

## Exploring the Benefits of Using *CryoSat-2*'s Cross-Track Interferometry to Improve the Resolution of Multisatellite Mesoscale Fields

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### Abstract:

Sea surface height (SSH) measurements provided by pulse-limited radar altimeters are one-dimensional profiles along the satellite's nadir track, with no information whatsoever in the cross-track direction. The anisotropy of resulting SSH profiles is the most limiting factor of mesoscale SSH maps that merge the 1D profiles.

This paper explores the potential of the cross-track slope derived from the *Cryosphere Satellite-2* (*CryoSat-2*)'s synthetic aperture radar interferometry (SARin) mode to increase the resolution of mesoscale fields in the cross-track direction. Through idealized 1D simulations, this study shows that it is possible to exploit the dual SARin measurement (cross-track slope and SSH profile) in order to constrain mesoscale mapping in the cross-track direction.

An error-free SSH slope allows a single SARin instrument to recover almost as much SSH variance as two coordinated altimeters. Noise-corrupted slopes can also be exploited to improve the mapping, and a breakthrough is observed for SARin errors ranging from 1 to 5  $\mu\text{rad}$  for 150-km-radius features in strong currents, and 0.1–0.5  $\mu\text{rad}$  for global mesoscale.

Although only limited experiments might be possible with the error level of current *CryoSat-2* data, this paper shows the potential of the SAR interferometry technology to reduce the anisotropy of altimeter measurements if the SARin error is significantly reduced in the future, and in particular in the context of a prospective SARin demonstrator optimized for oceanography.

**Keywords:** Altimetry, Remote sensing, Sampling, Satellite observations, Interpolation schemes, Inverse methods

### 1. Introduction and context

In contrast with wide-swath imagers (e.g., sea surface temperature or ocean color), the data record of radar altimeters is exceedingly anisotropic. Sea surface height (SSH) measurements from pulse-limited radar altimeters are one-dimensional profiles along the satellite's nadir

41 track, with no SSH information whatsoever in the cross-track direction. Figure 1 shows that,  
42 for a single altimeter flying on the TOPEX/Jason orbit, the along-track (white segment)  
43 resolution can be as small as 7 km (level 2 product, 1 Hz rate), whereas in the cross-track  
44 resolution (black segment) it can be as large as 300 km.

45 To reconstruct 2D gridded fields of SSH or sea level anomalies (SLA), it is therefore  
46 necessary to interpolate 1D profiles (e.g. AVISO 2010; Dibarboure et al, 2011a, Leben et al  
47 2002, 2011). Optimal interpolation (OI) exploits an a priori statistical knowledge of the SLA  
48 field characteristics (e.g. Jacobs et al 2001; Le Traon et al, 2003) and measurement error (e.g.  
49 Philipps et al, 2012) as an additional constraint to merge 1D profiles from multiple sensors in  
50 an optimal way (e.g. Ducet, et al 2000).

51 The anisotropy of SSH profiles is by far the most limiting factor of gridded SSH mesoscale  
52 fields (Le Traon & Dibarboure, 2002, 2004; Pascual et al 2006), and especially in near real  
53 time where measurements “from the map’s future” are not yet available (Pascual et al, 2008).  
54 There are two practical consequences to this limitation.

55 Firstly, even if the spatial and temporal scales used to constrain the OI are derived from SSH  
56 measurement of 2 to 4 satellite constellations, the mapping is limited in the cross-track  
57 direction. Because 1D profiles from multiple sensors are blended into one map, 2D mesoscale  
58 mapping uses a compromise between actual mesoscale correlations and the sampling  
59 limitations from such constellations (Ducet, et al, 2000).

60 The resolution of mesoscale fields is dominated by the number of altimeters in operation.  
61 Chelton & Schlax (2003), Le Traon & Dibarboure (2002), and Chelton et al (2011) have  
62 shown that mesoscale maps have a limited global resolution capability. Higher resolution can  
63 still be achieved, but only locally, at certain times, when enough 1D profiles are available  
64 (Dussurget et al, 2011).

65 In this context, a new technology used on CryoSat-2 has the unprecedented potential to add  
66 actual measurements to constrain mesoscale mapping in the *cross-track* direction. Indeed, in  
67 addition to a classical pulse-limited radar altimeter measurement (also known as low  
68 resolution mode or LRM), CryoSat-2's altimeter SIRAL features a synthetic aperture radar  
69 interferometry (SARin) mode able to measure the SSH slope in the cross-track direction  
70 (Francis et al, 2007) as illustrated by Figure 2. In this paper, the cross-track slope (CTS) is  
71 given in micro-radians: a 1  $\mu\text{rad}$  slope is approximately equal to a SSH gradient of 1 cm for  
72 10 km, or a geostrophic current of 10 cm/s at mid-latitudes.

73 In this paper we use idealized OI simulations to explore the potential of the CTS derived from  
74 CryoSat-2's SARin mode to increase the resolution of mesoscale fields in the cross-track  
75 direction (methodology introduced in section 3). Our approach is to look at SARin technology  
76 in optimal conditions in section 4 and then to discuss what can be done in practice with  
77 current and future datasets in section 5.

## 78 **3 Methodology**

### 79 **3.1 Overview**

80 Le Traon and Dibarboure (2002, 2004), Chelton et al (2003) and Dibarboure et al (2011a)  
81 have shown that 2D SSH mapping is affected by many parameters (e.g. geometry, phasing or  
82 coordination of the constellation's orbits, high frequency ocean dynamics). To measure the  
83 potential of using a SARin slope to constrain mesoscale mapping we therefore use a simpler  
84 idealized 1D configuration.

85 We specifically focus on the cross-track direction (black segment from Figure 1) where the  
86 resolution is limited by the number of satellites in the constellation. In other words, this is a  
87 configuration where SARin slopes are ideal to complement lacking SSH measurements.

88 While performed for a given cross-track resolution (i.e. latitude), general conclusions can be  
89 derived from our analysis because correlation scales were shown to decrease with latitude as  
90 well (e.g. Jacobs et al 2001; Le Traon et al, 2003).

91 We measure the performance of mesoscale mapping using the following protocol:

- 92 • we simulate a mesoscale SSH “reality” profile, and we consider that the reality  
93 profile is in a frozen state, i.e. stationary over the 10-day period of a T/P or Jason  
94 repeat cycle (this strong assumption is discussed in section 5.3)
- 95 • the reality SSH field is sampled on measurement points to create error-free  
96 observations,
- 97 • errors are optionally added to the observations,
- 98 • observations are injected into a 1D optimal interpolator to create a “reconstructed”  
99 mesoscale field at the original resolution.

100 In this process, the statistical variance and correlation scales of the reality field are known  
101 analytically. Consequently the reconstruction is perfect if performed from enough error-free  
102 observations. In other words, differences between the reality and the reconstructed fields are  
103 the result of the omission or sampling error (not enough data to observe the signal) and  
104 commission or measurement error.

105 Note that, there is an additional error source in the mapping of real data: the imperfect  
106 modeling of signal and error covariances. This point is discussed in section 5.3.

### 107 **3.2 Methodology**

108 In this paper, we generate our reality  $H_{\text{real}}$  as a spatially correlated random Gaussian process  
109 (Equation 1). The default decorrelation scale is 150 km, i.e. consistent with findings from Le  
110 Traon et al (2003). In our first simulations (section 4.1), the oceanic variability used is 20 cm

111 RMS, i.e. we focus on zones of intense mesoscale activity (e.g. western boundary currents).  
112 Then we expand to different signal to error ratios in section 4.2. In section 5.3, we discuss the  
113 validity of the Gaussian methodology.

$$\text{Corr}(H,H) = \exp ( -x^2/d^2 ) \quad \text{Equation 1}$$

114 Our observation field  $H_{\text{obs}}$  is then constructed (Equation 2) by interpolating  $H_{\text{real}}$  at the desired  
115 resolution (30 km for an along-track simulation, 300 km for a cross-track simulation on a  
116 Jason-like orbit, and 100 km for a cross-track simulation on the CryoSat-2 orbit) and adding a  
117 white noise of 0.5 to 2 cm to the interpolated SSH values. This is arguably a pessimistic error  
118 level at 100+ km if compared with results from Dorandeu et al 2004, or Ablain et al 2011 : the  
119 noise they observe at a 7 km resolution would be reduced by along-track filtering of the SSH  
120 (factor of 2 for 30 km super-observations).

$$H_{\text{obs}}(\mathbf{x}) = H_{\text{real}}(\mathbf{x}) + \epsilon(\mathbf{x}) \quad \text{Equation 2}$$

121 Simulations shown in this paper do not include any along-track bias or long wavelength  
122 correlated errors as our sensitivity studies show no significant difference with noise-limited  
123 simulations. Although not shown in this paper, our simulated 1D mapping is degraded by  
124 correlated errors like operational mesoscale 2D mapping (e.g. Dibarboure et al, 2011c), but  
125 the anisotropy effect presented in section 3.3 and the impact of using SARin presented in  
126 section 3.5 are the same.

127 The reconstruction of the estimated mesoscale field  $H_{\text{est}}$  is performed with a 1D optimal  
128 interpolation derived from Bretherton et al (1976).  $H_{\text{est}}$  is obtained from Equation 3 where  $A$   
129 is the matrix describing the covariance between  $H_{\text{est}}$  and  $H_{\text{obs}}$  of Equation 2, and  $C$  the matrix  
130 describing the covariance between the SSH observations (covariances are derived from  
131 Equation 1). The formal reconstruction covariance error matrix  $E$  is obtained from Equation

132 4, although in practice only its diagonal is used here (1- $\sigma$  gray envelope around reconstructed  
133 SSH profiles).

$$\mathbf{H}_{\text{est}} = \mathbf{A} \cdot \mathbf{C}^{-1} \cdot \mathbf{H}_{\text{obs}} \quad \text{Equation 3}$$

134

$$\mathbf{E} = \mathbf{I} - \mathbf{A} \cdot \mathbf{C}^{-1} \cdot \mathbf{A}^T \quad \text{Equation 4}$$

135 Many figures shown in this paper are limited to 3000 km segments for the sake of illustration  
136 but simulations were performed on very long profiles to ensure that the examples in this paper  
137 are representative of the statistical behavior of each configuration.

### 138 **3.3 Observation anisotropy**

139 Figure 3 shows one reality segment, sampled in the along-track direction (every 30 km) with  
140 2 cm white noise added. The reconstructed field after optimal interpolation is almost identical  
141 to the reality field. The reconstruction error is 1.2 cm RMS i.e. 0.4% of the reality signal  
142 variance (18 cm RMS). Similarly, the along-track slope (bottom panel) is almost perfectly  
143 observed in the along-track direction.

144 Figure 4 shows the same reality segment, but positioned as a transect in the cross-track  
145 direction (black segment from Figure 1). In other words, each measurement (black dot) is the  
146 crossover between the transect and a different satellite track. In this figure, the SSH reality is  
147 sampled by a LRM altimeter every 300 km, i.e. the worst case configuration of a TP/Jason  
148 orbit. Because the Nyquist criterion is not met with a single satellite, many features are  
149 missed entirely in the reconstruction (e.g. at km #1000 or #1800 or #2200). The error  
150 reconstruction RMS is 46% of the signal variance. This figure illustrates the inability of a  
151 single satellite to observe large mesoscale, let alone features with radii smaller than 150 km.

152 Adding a second LRM altimeter (perfectly coordinated with the first one, i.e. like in the  
153 TOPEX/Jason tandem) significantly improves the resolution of the mesoscale field as shown

154 by Figure 5. Although the Nyquist criterion is barely met, the reconstruction error is  
155 significantly reduced with an error of 8% of the signal variance. The error is consistent and  
156 slightly larger than the 5% obtained by Le Traon et Dibarboure (2002) because this segment  
157 represents the widest gap between roughly parallel tracks. These scenarios give the 1-LRM  
158 and 2-LRM reference configurations to which SARin experiments can be compared to infer  
159 the cross-track slope contribution in an ideal case.

160 The formal mapping error from Equation 4 is visible in each simulation as a grey envelope of  
161 vertical bars. This theoretical error is *–for these idealized simulations–* a very accurate  
162 statistical estimate of the error which could be made during the reconstruction process: the  
163 differences between the real (plain) and the reconstructed (dashed) SSH are consistent with  
164 the 1- $\sigma$  boundaries defined with the grey envelope from Figure 3 and Figure 4.

165 The formal error represents the sum of the measurement error and the sampling error: it is as  
166 small as 2 cm near observation points (measurement noise) and as large as tens of centimeters  
167 at the center of the 300 km window between satellite tracks (sampling error). The anisotropy  
168 of the altimetry system is illustrated by the difference between the along-track and the cross-  
169 track formal errors. In the along-track direction (Figure 3) the error is always very small and  
170 dictated by the measurement error level whereas in the cross-track direction (Figure 4) the  
171 sampling error largely dominates between satellite tracks.

### 172 **3.4 CryoSat-2's cross-track measurement**

173 CryoSat-2 is ESA's ice mission (Francis, 2007). Equipped with an innovative radar altimeter  
174 (SIRAL – Synthetic Aperture Interferometric Radar Altimeter), and high-precision orbit  
175 determination (POD), CryoSat-2's primary objective is to serve Cryospheric science  
176 (Wingham et al, 2006). Cryosat's altimeter is operated almost continuously over ocean,  
177 mainly in LRM (i.e. conventional altimetry) or in the delay doppler / synthetic aperture radar

178 (SAR) mode which provides higher along-track resolution and lower noise level (Raney,  
179 1998).

180 Furthermore, SIRAL also features a third mode: the SAR-interferometry (SARin) mode,  
181 which uses CryoSat-2's two antennas (Francis, 2007). The combination of SAR and  
182 interferometry makes it possible to determine the cross-track slope of the surface from which  
183 the echoes are arriving. This is achieved by comparing the phase of one receive channel with  
184 respect to the other as first suggested by Jensen, 1999.

185 With the SARin mode, CryoSat-2 can provide one estimate of the local CTS every 0.05  
186 seconds, in addition to the classical topography measurement (Figure 2). Moreover, the along-  
187 track resolution and the precision of the SSH is the same as for a LRM sensor (e.g. Jason-2).  
188 The resolution is 300 m in the along-track direction (synthetic footprint), and the slope is  
189 estimated from a cross-track footprint of the order of 7 km.

190 This unprecedented measurement was initially designed to be used over the margins of the  
191 Greenland and Antarctic ice sheets, where the surface slopes are steep. To that extent,  
192 SIRAL's SARin mode was designed to have a cross-track slope accuracy of 200  $\mu$ rad  
193 (Wingham et al 2006), but Galin et al (2012) reported a noise level of 20  $\mu$ rad at a 7 km  
194 resolution and a bias of 8  $\mu$ rad for 1000 km segments, using both detailed modeling of the  
195 finite radar resolution in range and angle, and the thermally driven behavior of the  
196 interferometer bench.

197 This should be compared to the typical mesoscale slope distribution in zones of intense  
198 mesoscale activity which ranges from 1 to 5  $\mu$ rad at 150 km with values higher than 10  $\mu$ rad  
199 on the edges of the largest eddies (observed on multi-satellite SSH maps from AVISO, 2010).  
200 Assuming that the long wavelength errors described by Galin et al (2012) are minimized with  
201 empirical cross-calibration mechanisms (discussed in section 5.2.1), and that the noise level is

202 reduced by along-track filtering (discussed in section 5.2.2), it would become possible to use  
203 the SARin slope as a constraint for mesoscale mapping in the cross-track direction where  
204 LRM altimeters are blind.

205 Because the error level reported on CryoSat-2 is high with respect to the oceanic signal, our  
206 rationale is the following: we first look at the benefits of using error-free SARin CTS (section  
207 4.1), then we perform sensitivity studies with respect to the ocean variability and  
208 measurement errors (section 4.2). From this background, we discuss the practical case of  
209 CryoSat-2 in section 5.

### 210 **3.5 Improving the reconstruction with the cross-track slope**

211 Figure 2 gives a qualitative illustration of how mesoscale mapping can exploit the SARin  
212 cross-track slope. Subplot (a) shows a 500 km along-track LRM profile with SSH only  
213 (simulated, error-free), whereas subplot (b) shows the information given by a SARin profile  
214 with SSH and cross-track slope. Both plots correspond to the reality from subplots (c) and (d).

215 From the SSH+CTS sample (subplot b), one can assume that the maximum value at -100 km  
216 is located on the right-hand side of the nadir track, that the minimum value at +150 km is  
217 probably near the nadir track, and that the maximum value at +400 km is located on the left-  
218 hand side of the nadir track. Adding a statistical description of mesoscale variability and  
219 slopes, it is possible to enhance the mapping in the cross-track resolution up to a distance  
220 equal to the spatial correlation radius.

221 This is achieved using a method derived from Le Traon and Hernandez (1992): we replace the  
222 SSH observation vector  $H_{\text{obs}}$  in Equation 3 by a vector composed of all observations (SSH and  
223 CTS), and we update matrixes A and C from Equation 3 and Equation 4 accordingly (see  
224 Appendix).

## 225 **4 Results**

### 226 **4.1 Error-free simulations**

227 In this section we infer what would be the optimal mesoscale improvement using SARin on a  
228 300-km cross-track resolution (i.e. a Jason-like orbit). It is optimal for SARin in the sense that  
229 Le Traon and Dibarboure (2002, 2004) have shown that the main weakness of this orbit is the  
230 cross-track resolution, and it is the reason why TOPEX/Jason and Jason-1/Jason-2 were put in  
231 a spatially interleaved configuration. Thus we use this “reference orbit” and “reference  
232 tandem” to SARin-based simulations. We discuss the difference between this Jason-like  
233 configuration the (suboptimal) Cryosat orbit in section 5.1.

234 Adding the SARin slope constraint (error free) significantly improves the OI reconstruction as  
235 shown by Figure 6. This plot should be compared to Figure 4 where one LRM altimeter was  
236 barely able to recover 50% of the signal variance in the cross-track direction (Nyquist  
237 sampling not achieved). Thanks to local constraints on the SSH derivative, it is possible to  
238 recover features that were previously missed entirely (e.g. at km #1800 and #2200).

239 Quantitatively, on this example, the reconstruction error is only 6.96 cm RMS, i.e. 15% of the  
240 signal variance (vs. 50% for the LRM scenario on Figure 4). In other words, about 35% of the  
241 signal variance was recovered with the error-free slope. The 15% residual error should also be  
242 compared to the 8% of the configuration with two LRM altimeters (Figure 5): in this idealized  
243 simulation, a single SARin altimeter performs almost like two LRM altimeters.

244 Similarly, Figure 7 shows that a perfectly coordinated constellation of 2 SARin altimeters  
245 flying on a Jason-like orbit (150 km cross-track resolution) is able to properly reconstruct the  
246 SSH and slope reality fields even though the Nyquist criterion is barely met with SSH alone.  
247 Because slopes and covariance models add the constraint needed, the reconstruction error is

248 only 1.83 cm RMS (i.e. 1% of the signal variance) and largely due to the error outlier of the  
249 first measurement and the 2 cm SSH measurement noise.

## 250 **4.2 Sensitivity to signal to noise ratio**

251 We performed a series of sensitivity tests on the slope error for 1 and 2 SAR-in altimeter  
252 constellations using very long simulations (2000 times the correlation radius). Figure 8 shows  
253 the RMS of the error reconstruction as a function of the standard deviation (STD) of the  
254 simulated SARin slope error (plain line). The 1-LRM and 2-LRM references are also given by  
255 the black dotted and dashed lines. Note that the observed error is consistent with the formal  
256 error given by Equation 4.

257 As expected, the reconstruction error decreases as the CTS error does, and the sigmoid shape  
258 on the logarithmic abscissa scale indicates that the largest gains are obtained between 1 and 5  
259  $\mu\text{rad}$ , i.e. near the peak of the cross-track slope probability density function.

260 The upper asymptotic value for slope errors higher than 20  $\mu\text{rad}$  is 49%, i.e. the mapping error  
261 observed for 1 LRM sensor (dotted line). In other words, if the SARin error is large, it does  
262 not improve the reconstruction with the OI. Yet as expected from a theoretical point of view,  
263 this figure shows that even if the error STD of the CTS is 25 times larger than the SSH slope  
264 STD (i.e. factor of 600 in the covariance matrix), the OI never underperforms w.r.t to the 1-  
265 LRM scenario because untrustworthy observations are automatically downweighted by  
266 covariance matrix  $C$ .

267 If the OI covariance matrixes are properly set up, adding very noisy slope estimates (e.g. 10 to  
268 20  $\mu\text{rad}$  unmitigated error from Galin et al, 2012) can still improve the cross-track mapping,  
269 albeit in a very limited way.

270 The lower asymptotic value is 13% of the signal variance, i.e. only slightly larger than the 9%  
271 error observed with 2 LRM sensors (dashed line): using an error-free SARin instrument in an  
272 ideal configuration (1D, cross-track, 150 km radius for a 300 km sampling resolution) does  
273 not allow to fully reconstruct the signal, but a single SARin instrument yields results almost  
274 as good as two LRM sensors as per the example from Figure 5 and Figure 6. The residuals  
275 arise from sampling errors: although additional error-free parameters are used, there are still  
276 not enough measurements points to correctly resolve all mesoscale structures.

277 Results are similar for 2 x SARin simulations in Figure 9, even though the gain is more  
278 limited owing to the fact that 2 coordinated LRM altimeters already have a good sampling  
279 capability for 150 km radius features (Le Traon et Dibarboure, 2004). In this figure, the lower  
280 asymptotic value is 1.2%, i.e. very close to the 1% obtained with 4 coordinated LRM sensors:  
281 sampling errors would become marginal in a coordinated 2 x SARin configuration.

282 Because the variability of the cross-track slope is proportional to the variability of the SSH,  
283 we performed sensitivity studies to the latter (using constant correlation scales and SSH noise  
284 levels) to see how results from section 4.1 could be extrapolated out of intense mesoscale  
285 activity zones.

286 Figure 10 confirms that the reconstruction error is still sigmoid-shaped, and shifted along the  
287 abscissa axis as a function of the SSH variability. The breakthrough in mapping improvement  
288 is always achieved for slope error STD ranging from  $0.5 \cdot \sigma_{\text{slope}}$  to  $2 \cdot \sigma_{\text{slope}}$ .

289 To be used globally in mesoscale mapping, SARin slopes would require an error level of the  
290 order of 0.1 to 0.5  $\mu\text{rad}$  for mesoscale wavelengths. This is largely beyond what can be  
291 achieved with current data from Cryosat-2 (discussed in section 5.2).

## 292 **5 Discussion: from theory to practice**

### 293 **5.1 Sensitivity to the satellite track geometry**

294 The sampling pattern of the CryoSat-2 orbit (current SARin mission) and the Jason orbit  
295 (simulations from section 4) are very different. The latter has a 10-day repeat cycle (300 km  
296 cross-track resolution from Figure 1). In contrast, CryoSat-2 has a one-year repeat cycle with  
297 3-day and 30-day sub-cycles, i.e. globally homogeneous sampling patterns with 1000 and 100  
298 km cross-track resolutions respectively (Francis et al, 2007). CryoSat-2's orbit has no sub-  
299 cycle in the 10 to 20 day range associated with mesoscale temporal decorrelation (Jacobs et  
300 al, 2001).

301 As a result, for any 10 to 20 day period, CryoSat-2's measurements are aggregated in band-  
302 shaped patterns (100 km wide per 3-day sub-cycle) which are interleaved with band-shaped  
303 "blind spots" with no recent SSH observation (Figure 11). The impact on mesoscale  
304 observation in LRM mode is discussed by Dibarboure et al (2011c). As far as SARin slopes  
305 are concerned, there are two consequences of CryoSat-2's sampling pattern.

#### 306 **5.1.1 Track aggregation and data gaps**

307 Firstly the SARin slopes located on the outer edges of the band-shaped aggregation of satellite  
308 tracks provide a unique capability to reduce the extent of the band-shaped blind spots by up to  
309  $2 * 150$  km (one slope constraint on each side of the diamond not covered by CryoSat-2  
310 tracks in Figure 11). This is useful to balance CryoSat-2's main sampling weakness when it  
311 comes to mesoscale observation.

312 Figure 12 illustrates this point: it shows the OI reconstruction for a 1500 km cross-track  
313 segment where CryoSat-2 measurements are aggregated in 100 km resolution bands where  
314 mesoscale features (150 km radius) are resolved, and interleaved with a 500 km wide blind

315 spot where no CryoSat-2 track is available in the 15 day window corresponding the frozen  
316 field approximation.

317 The SARin-based reconstruction (subplot a) is slightly better because the outer edges are  
318 constrained by error-free slope estimates whereas the LRM-based reconstruction (subplot b)  
319 is not able to observe even a fraction of the large eddy at km #700 and the total reconstruction  
320 error is much higher (12.1 cm RMS vs. 6.7 cm RMS for SARin). Note that the overall  
321 improvement is limited to the outer edges of the large data gap (one decorrelation radius on  
322 each side) because the OI cannot “guess” the existence mesoscale structures if they are not  
323 remotely observed.

### 324 **5.1.2 Orbit sampling differences**

325 The second consequence of CryoSat-2’s sampling pattern is the cross-track resolution *within*  
326 the track aggregations. CryoSat-2’s sampling “bands” have a cross-track resolution of 100  
327 km, i.e. more favourable to the observation of 150 km radius mesoscale features, albeit in  
328 limited areas. In this context, SARin data from Cryosat might be used to recover smaller  
329 mesoscale features (only within the satellite track aggregation).

330 Table 1 shows the mapping improvement (i.e. the reduction of cross-track reconstruction  
331 error) when the “reality” and OI correlation radiuses range from 50 to 150 km and for the  
332 Jason and CryoSat-2 orbits. All simulations were performed with a slope measurement noise  
333 of 1  $\mu$ rad. On the Jason orbit, the cross-track mapping is improved mainly for large mesoscale  
334 (18%) but not for short mesoscale (5%) because the SARin slope cannot balance the limited  
335 resolution of the Jason orbit. The opposite is observed for CryoSat-2 (in the aggregation  
336 bands) owing to its cross-track 100 km cross-track resolution: the improvement is limited for  
337 100 km or more and the highest improvement is observed for a 50 km radius.

338 In other words, with the CryoSat orbit, the SARin slope is an asset to improve the cross-track  
339 observation of smaller mesoscale features (in the band-shaped aggregation of satellite tracks),  
340 something that would not be possible on a Jason orbit.

341 Yet higher wavenumber (K) mesoscale eddies also have a smaller amplitude (the SSH power  
342 spectrum decreases as a function of  $K^{-11/3}$  in the SQG theory, as per Le Traon et al., 2008).  
343 Thus changing the correlation radius also induces a reduction of the SSH STD and a reduction  
344 of the CTS STD from 2  $\mu$ rad to 1.5  $\mu$ rad (Table 1). In other words, higher precision SARin  
345 slopes would be needed in CryoSat-2's sampling bands because the smaller signal of interest  
346 also has weaker slopes.

347 To that extent, and considering the error level discussed in section 5.2, the CryoSat-2 orbit is  
348 less attractive than a Jason-like resolution would be, because the gain with SARin is  
349 geographically limited and because the orbit is more demanding in terms of CTS error budget.

## 350 **5.2 Slope error**

351 The simulations from section 4.2 showed that the enhancement of cross-track mesoscale  
352 mapping was possible in favorable signal to ratio conditions. The expected benefit from actual  
353 Cryosat-2 data raises the question of the error level of current datasets. Yet the error spectrum  
354 of SARin data in a mesoscale context is not known. Indeed, SARin acquisition zones on  
355 ocean are small and/or limited in time. So it is not possible to get datasets that are large  
356 enough to observe correlated errors in space or in time. The study from Galin et al (2012) is  
357 the first to provide a CTS error estimate as a bias and noise error on ocean through a  
358 comparison with a geoid model.

359           **5.2.1 Biases and long wavelength errors**

360 Galin et al (2012) report biases of the order of 8  $\mu$ rad on their 1000 km segments. It is not so  
361 much a true bias, as a long wavelength correlated error (e.g. orbital revolution) since they also  
362 observe a correlation with thermal conditions on the orbit (i.e. linear on 1000 km segments).  
363 Yet, in this paper, we are ignoring biases and long wavelength errors because we assume that  
364 they can be accounted for by multi-satellite cross-calibration.

365 Indeed, at the intersection of satellite tracks (e.g. CryoSat-2 x CryoSat-2 or CryoSat-2 x  
366 Jason-2) crossovers points provide a double measurement where the actual SSH anomaly  
367 signal is partially cancelled if the temporal distance between both measurements is short  
368 enough. It is thus possible to use this observation to detect and to mitigate spatially and  
369 temporally correlated signals.

370 Tai et al (1988) have used this approach to empirically reduce orbit errors on the SSH and  
371 Dibarboure et al (2011b) have demonstrated the feasibility of reducing the cross-track slope  
372 error for the wide-swath altimetry mission SWOT. So, in theory, the same method could be  
373 used to reduce CryoSat-2's SARin slope biases. The method would exploit crossover  
374 observations using a combination of the along-track and cross-track slope for SARin / SARin  
375 crossovers, and a projection of the along-track slope into the opposite along-track plane for  
376 SARin / LRM crossovers.

377 Alternatively, long wavelength errors (500 km or more) can be accounted for in the mapping  
378 process itself, with an approach derived from Ducet et al (2000). These techniques are used  
379 operationally to remove SSH biases and 1000 km errors before mesoscale mapping  
380 (Dibarboure et al, 2011a), including for datasets with limited coverage (e.g. ERS-2 after the  
381 loss of its on board recorders). The same method could be used in the geographically-limited  
382 SARin acquisition zones to cross-calibrate long-wavelength errors in the cross-track slope.

### 383           **5.2.2 Noise and short scale errors**

384   In the recovery of the cross-track slope, Galin et al (2012) also observe on average 20  $\mu$ rad of  
385   speckle-related noise at 1 Hz or 7 km resolution. The slope is computed from a distance  
386   ranging from 1 km to 8 km depending on the retrieval algorithm (phase-difference at the first  
387   point of arrival VS. model fit) and significant wave height (SWH) conditions.

388   The spatial correlation of mesoscale slope makes along-track filtering possible (including  
389   with non linear filters to remove spurious slopes) if the error is speckle-related (i.e. no along-  
390   track correlation of the CTS error). If we assume that a simple running average is used to get  
391   one super-observation for a 150 km radius (admittedly a crude filtering), the resulting  
392   mesoscale slope precision would be less than 4  $\mu$ rad with current slope retrieval algorithms.

393   Moreover, Galin et al investigate the origin of residual SARin slope outliers such as the  
394   influence of wind and so-called sigma0 blooms. Yet sigma0 blooms can be detected and  
395   edited out in pulse-limited LRM altimetry (Thibaut et al, 2010). We can therefore assume that  
396   the largest SARin slope outliers can be detected as well, thus decreasing the overall slope  
397   error RMS of a non-Gaussian slope error distribution.

398   With Cryosat-2 we can probably observe only large eddies ( $2\text{-}\sigma$ ) in zones of intense  
399   mesoscale variability. Elsewhere, SARin slopes from Cryosat-2 can probably barely improve  
400   cross-track mesoscale mapping because the instrument was not designed for this application  
401   (insufficient signal to error ratio).

### 402           **5.2.3 MSS and geoid errors**

403   In this section, we discuss MSS model errors and their influence on SARin slope anomalies in  
404   the context of mesoscale mapping. Indeed, mesoscale mapping is based on sea level  
405   anomalies (SLA), not sea surface heights (Dibarboure et al, 2012) and the SLA is created as

406 the difference of the measured SSH and a temporal reference or  $\langle \text{SSH} \rangle$ . The orbit used by  
407 CryoSat-2 is geodetic (one-year repeat cycle, described in Francis et al, 2007) so gridded  
408 MSS or geoid models are used as a  $\langle \text{SSH} \rangle$  reference. The same stands for CTS anomalies  
409 which are the difference of the CTS measurement and the cross-track gradient of the MSS  
410 model. Consequently any error in the models generates a CTS anomaly error (i.e. an  
411 additional CTS error in Figure 10).

412 Pavlis et al (2008) show that in favourable conditions along the well-known TOPEX/Poseidon  
413 tracks, they observe an error of  $2 \mu\text{rad}$  at 1 Hz for EGM08. In a different context, Sandwell  
414 and Smith (2009) have shown through comparisons with shipboard gravity that the accuracy  
415 of altimetry-derived gridded gradients was of the order of a few  $\mu\text{rad}$  in zones of rugged  
416 seafloor topography. More recently, Schaeffer et al (2012) have shown that the gradient error  
417 of their MSS model (CNES/CLS2011) ranged from  $1 \mu\text{rad}$  along charted tracks of repetitive  
418 altimetry mission to  $5 \mu\text{rad}$  in areas covered only by geodetic altimetry missions. Moreover,  
419 Andersen & Rio (2011) and Dibarboure et al (2012) highlighted differences between  
420 independent MSS models that range from 1 to 3 cm with wavelengths ranging from 3 to  
421 hundreds of kilometers (a few  $\mu\text{rad}$  after along-track smoothing).

422 The MSS/geoid error is therefore quite significant in the error budget of a SARin CTS  
423 anomaly, since it would add up to noise estimates from section 5.2.2. That error alone would  
424 make error-free CTS measurements difficult to use except in zones of strong mesoscale  
425 activity.

#### 426 **5.2.4 Expected and possible improvements**

427 Comparing the figures of merit from section 5.2 to the sensitivity studies from section 4.2  
428 shows that the precision needed to improve cross-track mesoscale mapping in strong currents  
429 is at the limit of CryoSat-2's current observation capability.

430 However, one might expect some improvements in the future. The primary error sources were  
431 shown to be speckle-related measurement noise and the MSS reference models used to  
432 generate the slope anomaly.

433 Concerning the former, it might be technically possible to update onboard software to get  
434 SAR data from both receive chains on ocean, and to change acquisition rates in SARin mode,  
435 essentially yielding 4 times as many independent looks, and reducing the noise level.  
436 Moreover, the SAR and SARin retrieval algorithms are relatively young, especially in an  
437 oceanography context (CryoSat-2 is an ice mission), and Galin et al give some interesting  
438 outlook that might result in a better precision: filtering, and weighting of beams...

439 And concerning the latter, our error estimate are derived from *2008-2001* generation MSS  
440 models, which are not yet integrating new geodetic data from CryoSat-2, Jason-1 GM  
441 (geodetic phase), let alone from new and upcoming missions flying on uncharted tracks (e.g.  
442 Sentinel-3A and 3B, HY-2). It is likely that the current and future altimeter datasets will  
443 decrease the error level of the future reference models, and especially at short wavelengths.

444 Beyond CryoSat-2, our findings raise the question of a prospective SARin demonstrator  
445 optimized for oceanography (with synergies with other applications). In this context, the  
446 outlook is even more promising because additional changes could be considered: on the orbit,  
447 on the hardware, and reference surface models.

448 CryoSat-2's orbit was shown to be suboptimal for SARin usage in section 5.1 and a dedicated  
449 mission could use a different orbit such as the ones analyzed by Dibarboure et al (2012) for  
450 the geodetic phase of Jason-1.

451 Moreover, if a new instrument derived from SIRAL were used on a dedicated SARin  
452 demonstrator, various upgrades could be considered to increase the number of statistically  
453 independent looks and to decrease the speckle-related noise: antenna design and beam width,

454 baseline length, pulse timing (e.g. continuous or interleaved mode VS SIRAL's burst  
455 mode)... However it is possible that the global mesoscale requirement from section 4.2  
456 (precision of the order of 0.1 to 0.5  $\mu$ rad) might remain challenging.

457 Lastly, in the context of *global* SARin acquisition with a sufficient precision, such a  
458 prospective mission would acquire east/west gradients which would help resolve the shortest  
459 wavelengths in MSS, geoid or bathymetry models since they are difficult to resolve with the  
460 current anisotropy of altimeter data (Sandwell et al., 2011). In turn, this would further  
461 mitigate the errors from the  $\langle$ SSH $\rangle$  reference discussed in section 5.2.3.

### 462 **5.3 Validity and limitations of this work**

463 In this section we discuss some approximations made in this paper, and the validity and  
464 limitations of these factors as an outlook for future work: the Gaussian properties of our  
465 “reality”, the perfect a priori knowledge used in the mapping process, the simple 1D mapping  
466 methodology used, and the lack of temporal variability.

- 467 • In section 3.2, our reality is a random Gaussian process with a decorrelation function  
468 consistent with scales reported by Le Traon et al (2003). In practice, our reality has a  
469 flat power spectrum density for long wavelengths and a cut-off for shorter  
470 wavelengths. In other words, we do not use the covariance model from operational  
471 mesoscale mapping (e.g. Ducet et al, 2010), but our covariance model and the  
472 associated variance-preserved power spectra are representative of a diversity of  
473 wavelengths, much like along-track filtered altimeter measurements.
- 474 • In the OI, we use a priori knowledge of the covariance of the signal ( $H_{\text{real}}$ ) and the  
475 covariance of the error ( $\epsilon$ ) in matrices A and C from Equation 3. In this paper, we use  
476 the true analytical covariance model used to simulate our dataset (i.e. the covariance

477 model of our Gaussian reality), resulting in a non-existent mapping error for error-free  
478 measurements. However, in practice, we only have an approximate knowledge of the  
479 true ocean decorrelation model (e.g. Jacobs et al, 2001 or Le Traon et al 2003) and of  
480 the altimetry error, so the OI process is not perfect. The same stands for the CTS  
481 parameter, and the a priori knowledge of the SARin data error. This can be a  
482 significant implementation problem so our findings should be revisited with real data.  
483 More importantly, this point highlights that one must acquire a better understanding of  
484 the SARin error spectrum before such data can be used in an OI.

485 • Lastly one should note that the frozen field assumption and the 1D analysis (cross-  
486 track direction) represent a best case configuration for SARin. In reality, mesoscale  
487 signals temporally decorrelate over  $\pm 15$  to 20 days. Thus our results are optimistic  
488 because they do not take into account the high frequency dynamics that Le Traon and  
489 Dibarboue (2002) showed to be difficult to resolve with constellations with less than  
490 4 altimeter missions. There is also a large panel of complex geometric configurations  
491 that vary with latitude. Consequently, because 1D results are encouraging, the findings  
492 of this paper should be extended to much more sophisticated 3D simulations (OI or  
493 ocean model assimilation), taking into account orbit sampling dynamics  
494 (measurements are not ubiquitous, nor regularly spaced out) and the temporal  
495 variability of the ocean (reality is not frozen).

## 496 **6 Conclusion**

497 CryoSat-2's SAR interferometry (SARin) mode has the unprecedented capability to measure  
498 the sea surface height slope in the cross-track direction. It is possible to use this parameter to  
499 constrain mesoscale mapping, and to improve the resolution in the cross-track direction where  
500 the traditional (LRM) radar altimetry is limited by the number of satellites in operations.

501 Idealized mapping simulations show that a single error-free SARin sensor on a Jason-like  
502 orbit has the potential to perform almost like two coordinated LRM instruments. Sensitivity  
503 studies show that the breakthrough in mapping improvement is achieved for slope errors  
504 between 1 and 5  $\mu\text{rad}$  for 150 km macro-observations, in zones of intense mesoscale activity.  
505 A better slope precision of the order of 0.1  $\mu\text{rad}$  would be needed for global usage and/or to  
506 resolve smaller features (radius < 100 km).

507 The precision needed to improve cross-track mesoscale mapping is probably at the limit of  
508 current SARin products from CryoSat-2 (and only after multi-satellite cross-calibration and  
509 along-track filtering) which might observe only the strongest slopes ( $2\text{-}\sigma$ ) in very energetic  
510 areas. The proof of concept is more attractive if we extrapolate to future improvements of  
511 SARin processors and ancillary datasets (e.g. MSS) and to a prospective mission improving  
512 upon SIRAL hardware and CryoSat-2 processors.

513 While encouraging, these results are optimistic, because all simulations were performed on a  
514 frozen SSH field (ocean dynamics and high frequencies are not taken into account), and only  
515 in the cross-track direction (i.e. optimal for the SARin slope) and they should be extended to  
516 much more complex 3D studies, or with real data from CryoSat-2.

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521 **8 Appendix: CTS methodology**

522 To use the cross-track slope in the OI process, we use the following covariance models to  
 523 describe the relationship between the topography H and the slope S:

$$cov_{hh}(\tau) = K \cdot e^{-\left(\frac{\tau}{d}\right)^2} \quad \text{Equation 5}$$

$$cov_{hs}(\tau) = -\frac{2\tau}{d^2} \cdot K \cdot e^{-\left(\frac{\tau}{d}\right)^2} \quad \text{Equation 6}$$

$$cov_{ss}(\tau) = K \cdot \left( \frac{4\tau^2}{d^4} - \frac{2\tau}{d^2} \right) \cdot e^{-\left(\frac{\tau}{d}\right)^2} \quad \text{Equation 7}$$

524 To inverse the problem, we replace A and C and H<sub>obs</sub> from Equation 3 (and Equation 4), by  
 525 A', C' and H' <sub>obs</sub>, where H' <sub>obs</sub> is the new observation vector including topography and slope  
 526 measurements as the sum of the true signal H<sub>real</sub> or S<sub>real</sub> and a random error ε<sub>H</sub> and ε<sub>S</sub>  
 527 estimated on the across-track position vector x(i):

$$H'_{obs} = \begin{bmatrix} H_{real}(x(1)) + \varepsilon_H(x(1)) \\ \vdots \\ H_{real}(x(N)) + \varepsilon_H(x(N)) \\ \vdots \\ S_{real}(x(1)) + \varepsilon_S(x(1)) \\ \vdots \\ S_{real}(x(N)) + \varepsilon_S(x(N)) \end{bmatrix} \quad \text{Equation 8}$$

528 Matrix C' is the new covariance matrix taking into account both topography and slope  
 529 estimates

$$C' = \begin{bmatrix} C_{hh} & C_{hs} \\ C_{hs} & C_{ss} \end{bmatrix} \quad \text{Equation 9}$$

530 where  $C_{hh}$ ,  $C_{hs}$ , and  $C_{ss}$  are the three covariance matrixes for each couple of observation type,  
 531 built as a function of the distance  $d_{ij} = |x(i) - x(j)|$  separating measurements points #i and #j.

$$C_{hh} = \begin{bmatrix} cov_{hh}(d_{1,1}) & \dots & cov_{hh}(d_{1,N}) \\ \vdots & cov_{hh}(d_{i,j}) & \vdots \\ cov_{hh}(d_{N,1}) & \dots & cov_{hh}(d_{N,N}) \end{bmatrix} \quad \text{Equation 10}$$

$$C_{hs} = \begin{bmatrix} cov_{hs}(d_{1,1}) & \dots & cov_{hs}(d_{1,N}) \\ \vdots & cov_{hs}(d_{i,j}) & \vdots \\ cov_{hs}(d_{N,1}) & \dots & cov_{hs}(d_{N,N}) \end{bmatrix} \quad \text{Equation 11}$$

$$C_{ss} = \begin{bmatrix} cov_{ss}(d_{1,1}) & \dots & cov_{ss}(d_{1,N}) \\ \vdots & cov_{ss}(d_{i,j}) & \vdots \\ cov_{ss}(d_{N,1}) & \dots & cov_{ss}(d_{N,N}) \end{bmatrix} \quad \text{Equation 12}$$

532 When the inversion is optimal we also account for the uncorrelated error  $\varepsilon_H$  and  $\varepsilon_S$  in the  
 533 diagonal of  $C_{hh}$ ,  $C_{hs}$ , and  $C_{ss}$  (not shown).

534 Matrix A' describing the covariance between the topography we want to reconstruct  $H_{est}$  and  
 535 the new observation vector  $H'_{obs}$  is created with the method used for C', but using the distance  
 536  $d'_{ij}$  between the position  $x(i)$  of our observation points and the position  $x'(i)$  our unknown  
 537 grid points.

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## 674 **10 Figure captions**

675 Figure 1: 10-day sampling from an altimeter on the TOPEX/Jason orbit. The white segment  
676 highlights the along-track direction with one measurement every 7 km, and the black segment  
677 highlights the worst case configuration in the cross-track direction with one measurement  
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679 Figure 2: Difference between LRM (a) and SARin (b) measurements in the optimal  
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691 Figure 5: Simulated Gaussian field (plain line), observation by two LRM altimeters in the  
692 cross-track direction (dots, 150 km resolution), and reconstruction at each time step with an  
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696 Figure 6: Simulated Gaussian field (plain line), observation by one SARin altimeter in the  
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708 point of inflection.

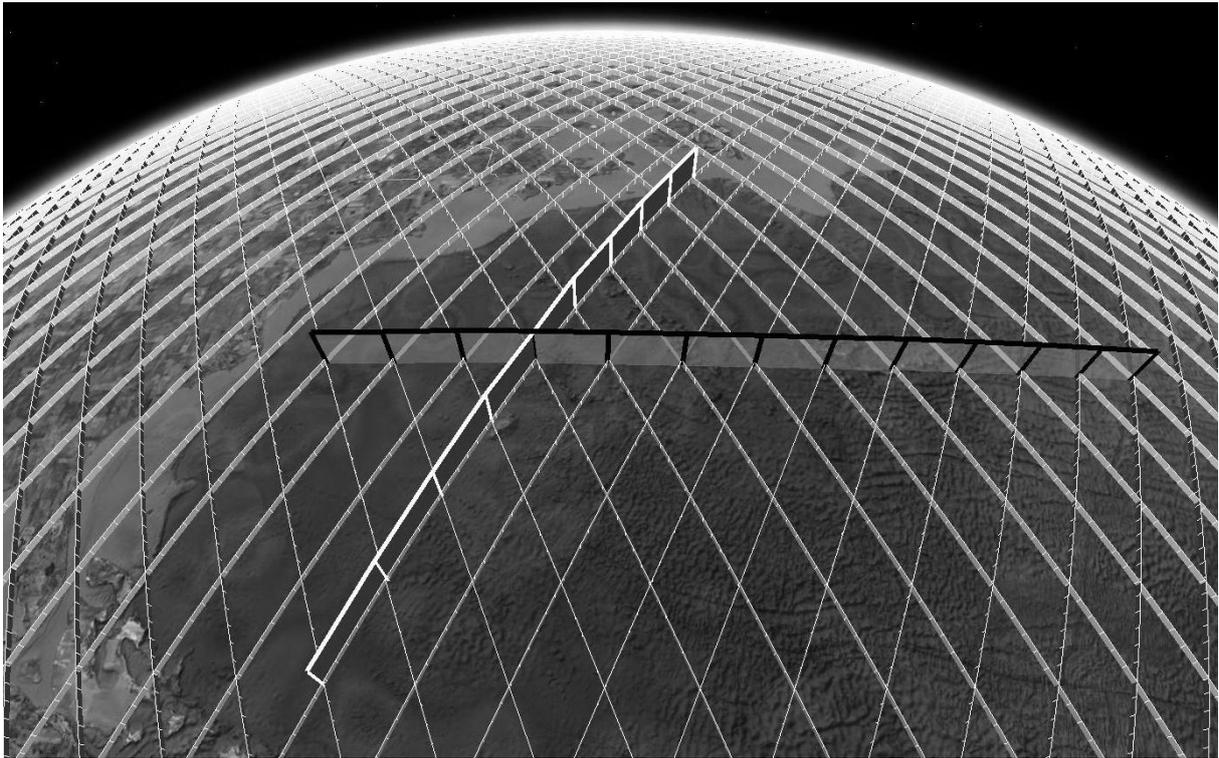
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713 Figure 10: Cross-track reconstruction error (in % of signal variance) for one SARin altimeter  
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717 aggregated in 500 km wide bands thanks to the 3 and 30 day sub-cycles, and interleaved with  
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719 Figure 12: Simulated Gaussian field (plain line), observation in the cross-track direction  
720 (dots) by one SARin altimeter on the CryoSat-2 orbit (100 km resolution, packet-aggregated  
721 tracks), and reconstruction at each time step with an optimal interpolation (dashed line).  
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723 OI reconstruction in the 500 km wide blind spot (grey rectangles).

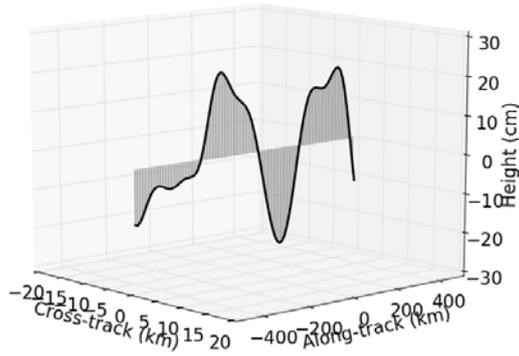
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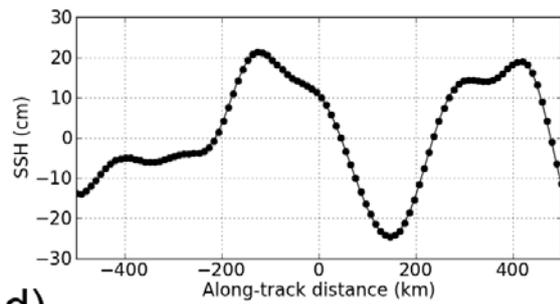
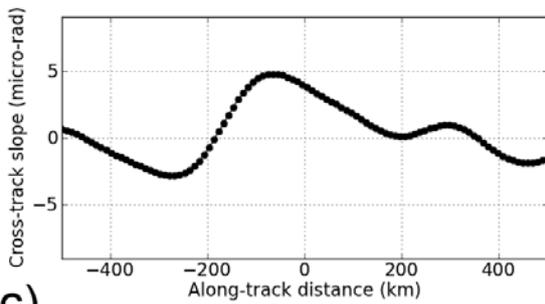
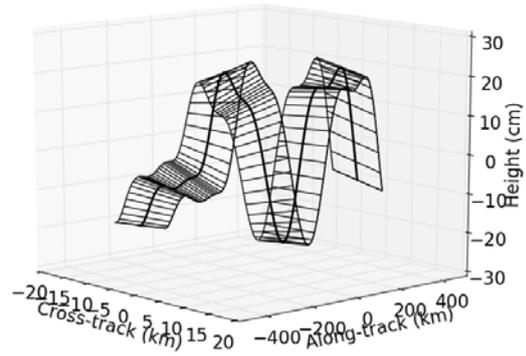
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a) LRM sea surface height profile



b) SARin height & slope profile



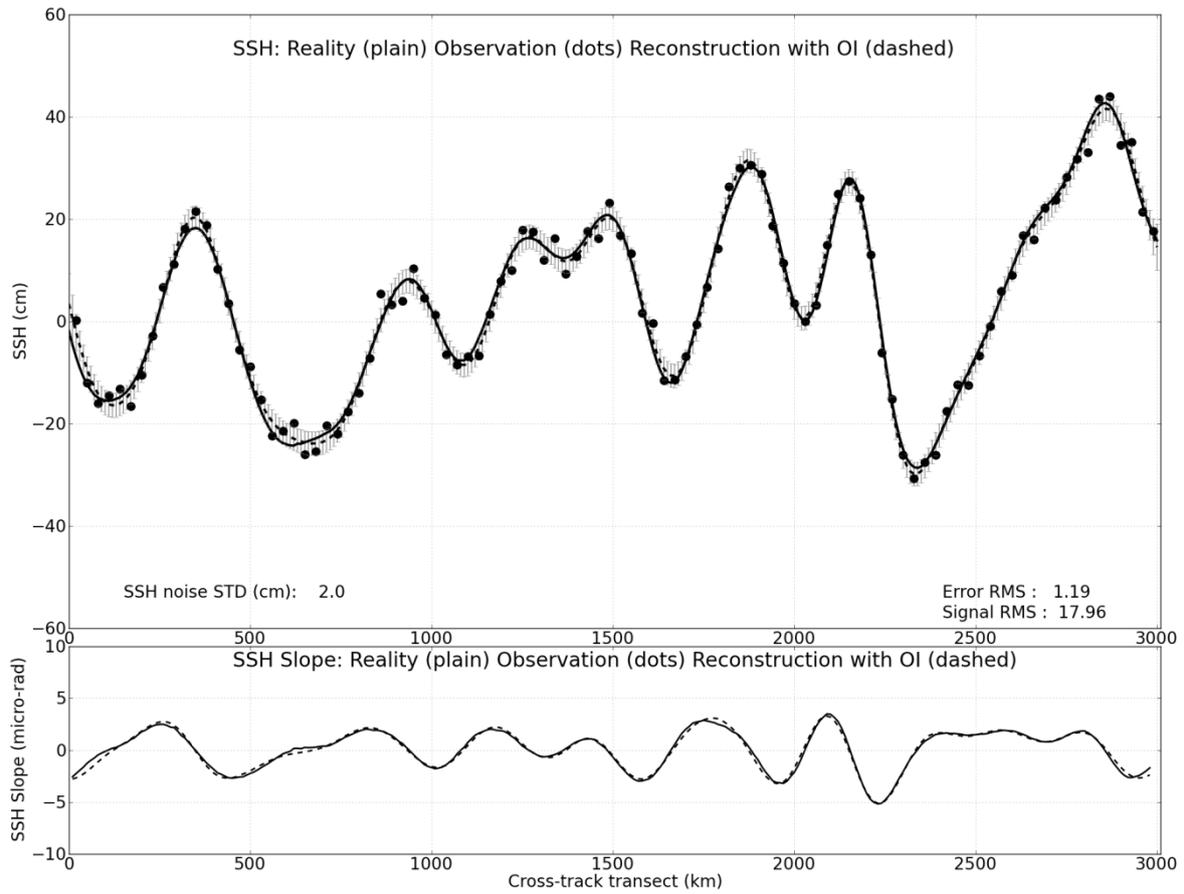
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c)

d)

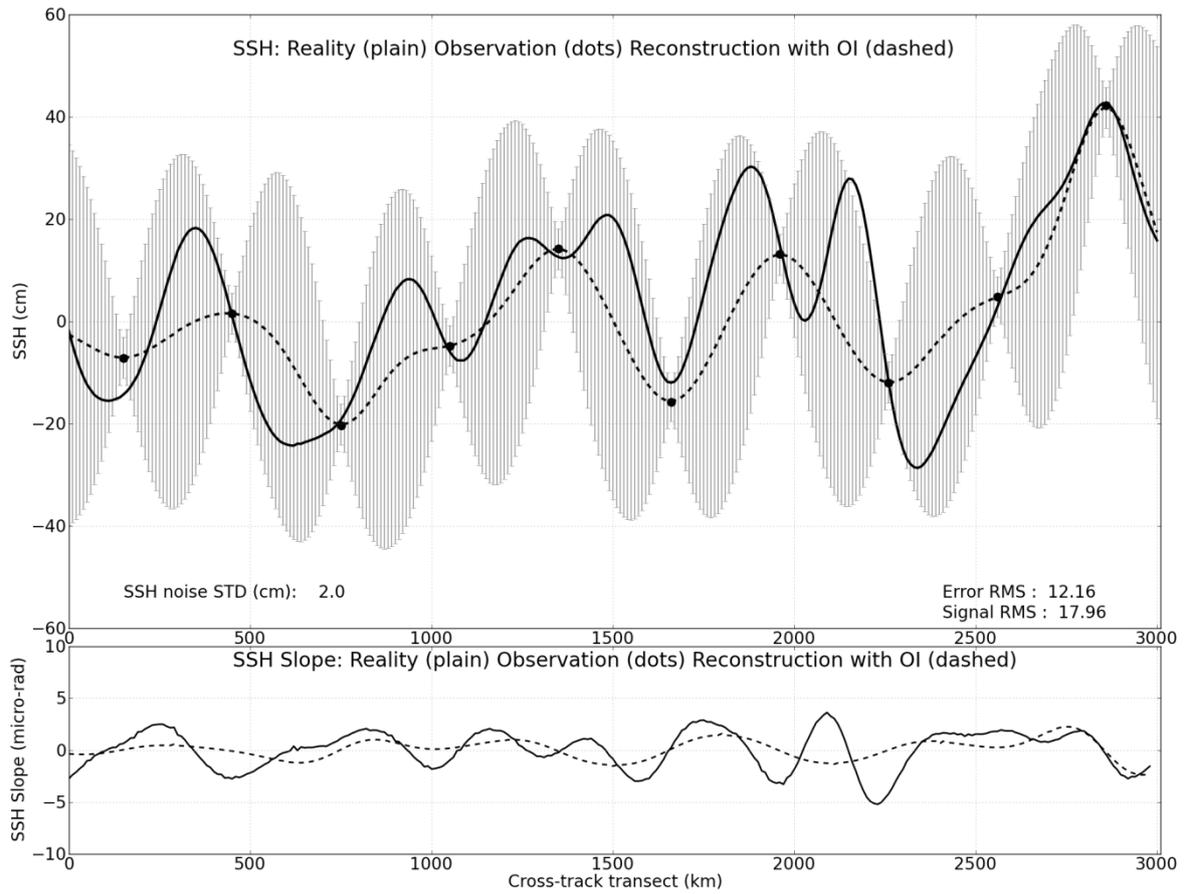
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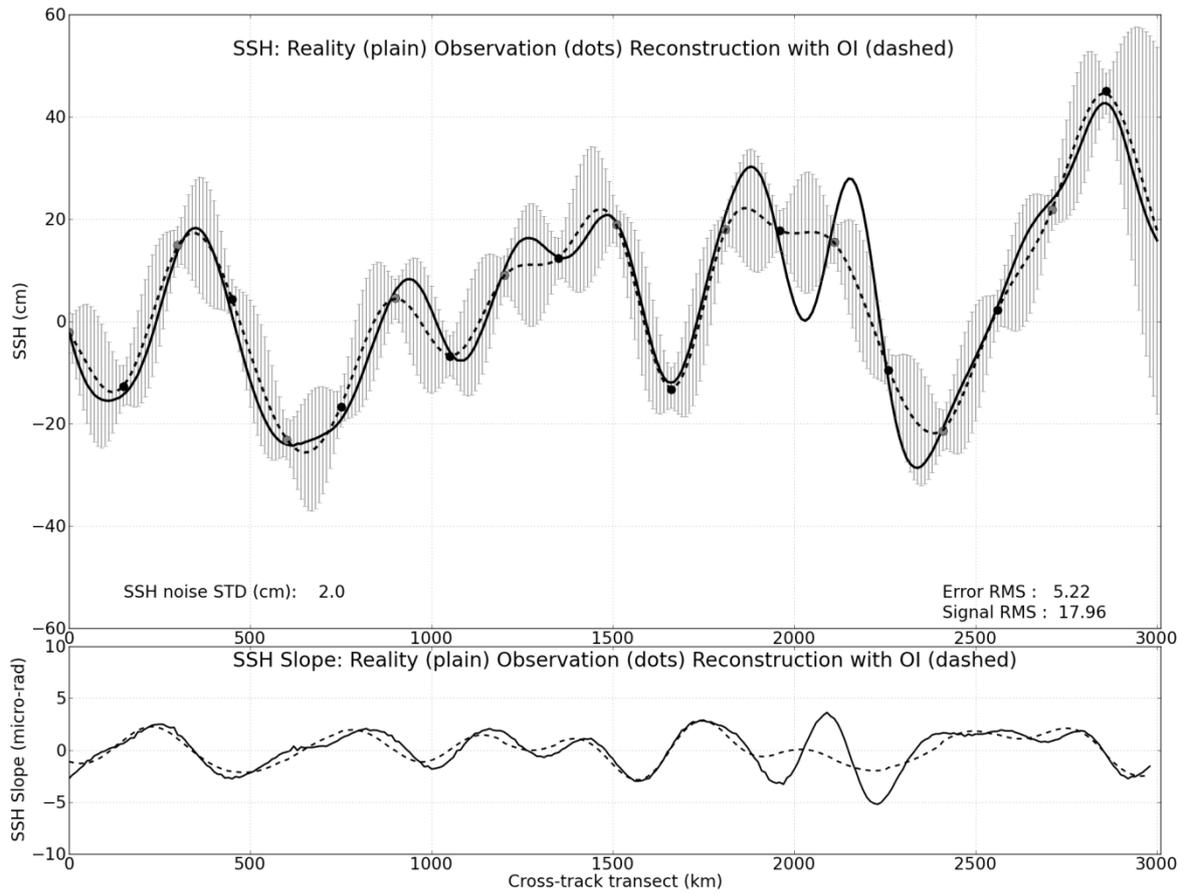
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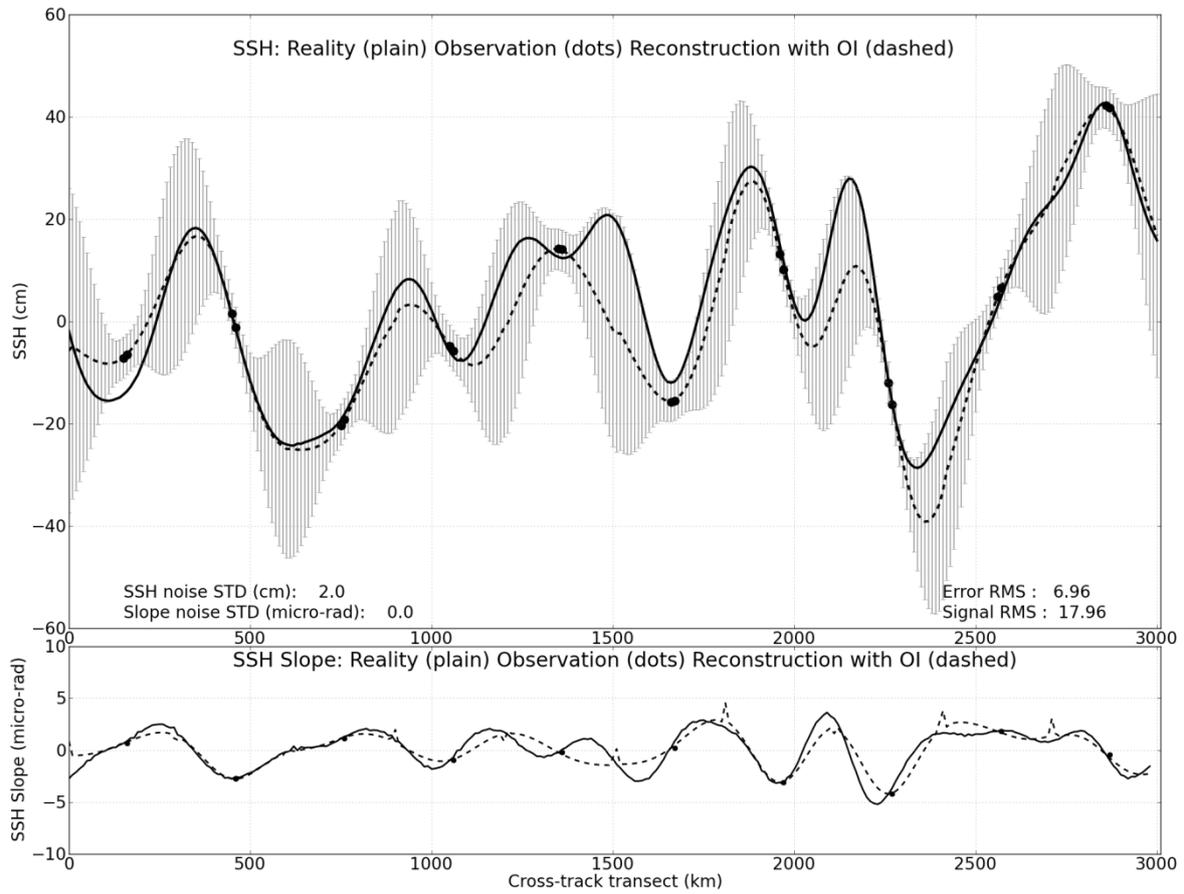
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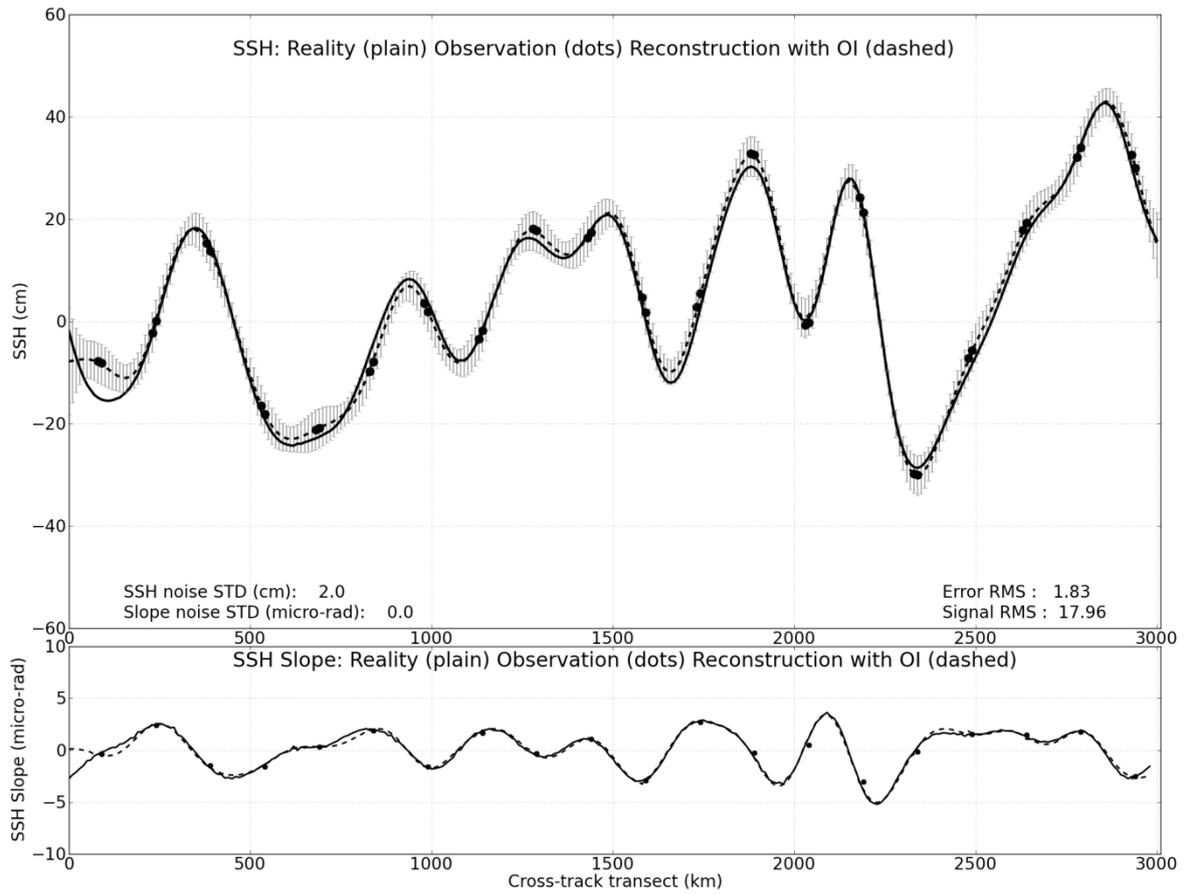
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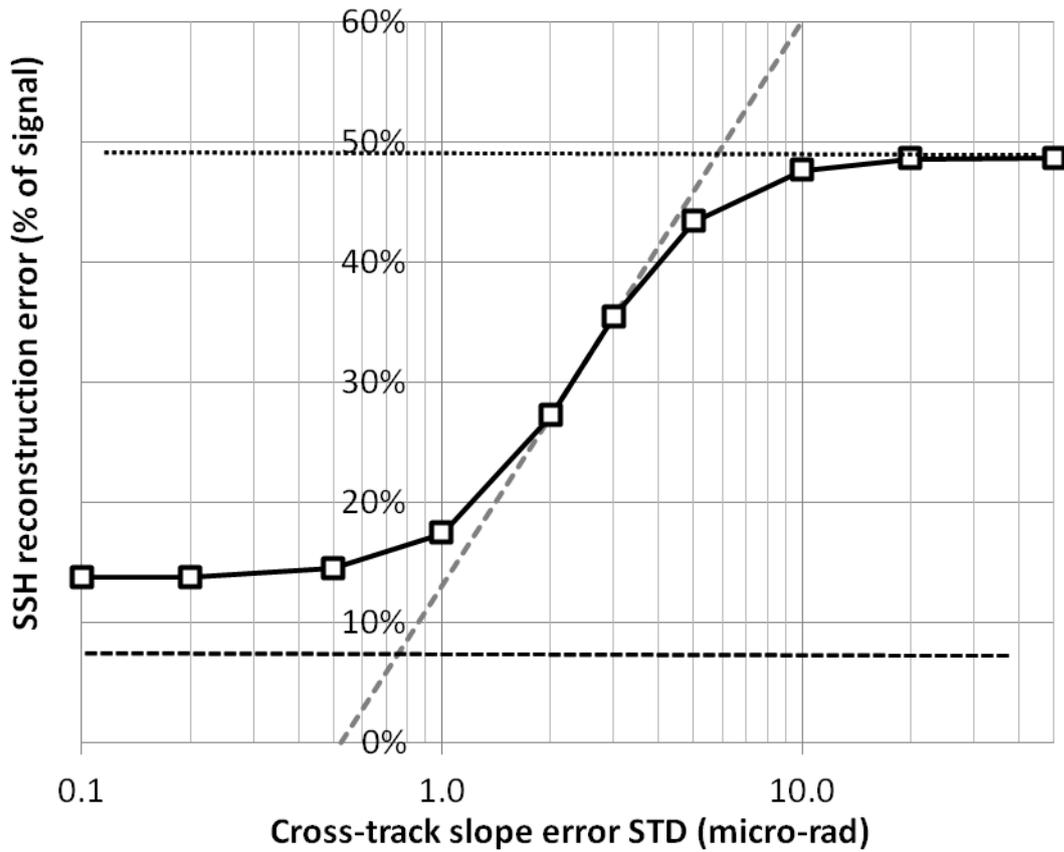
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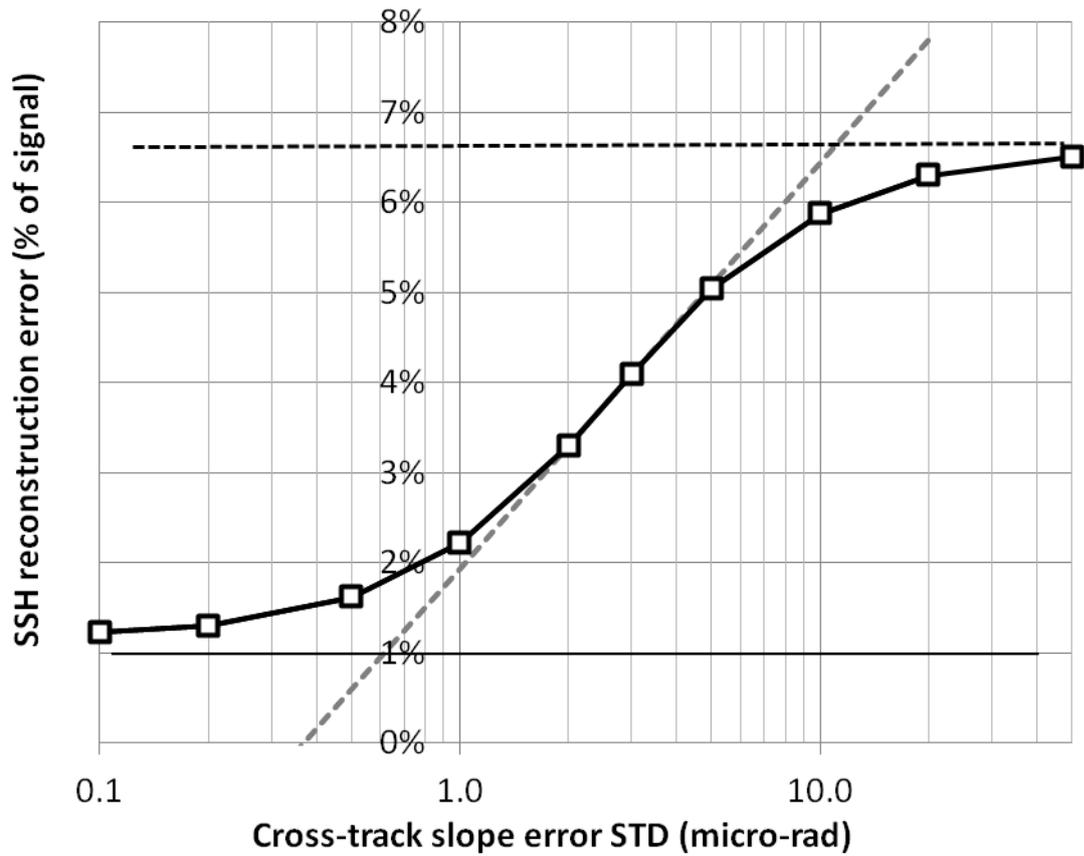
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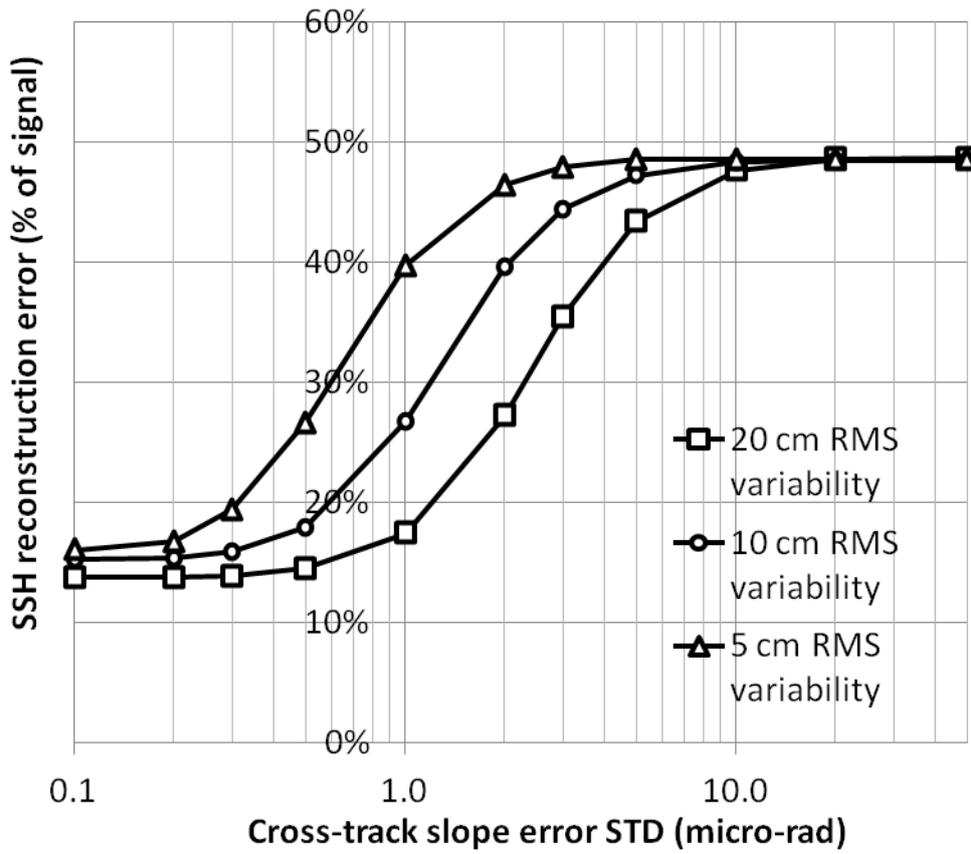
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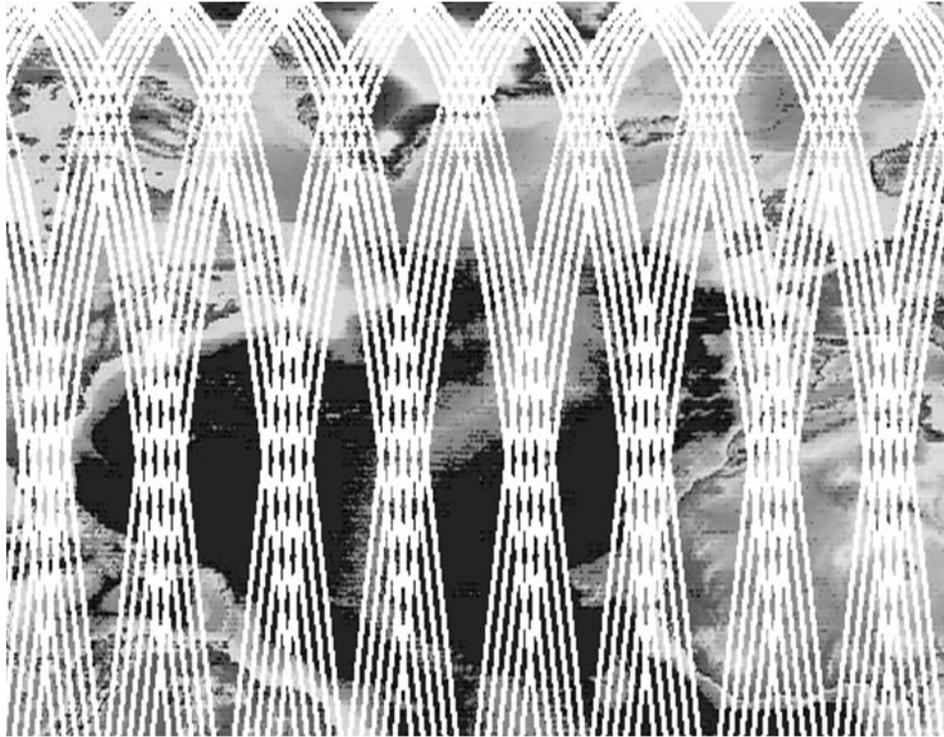
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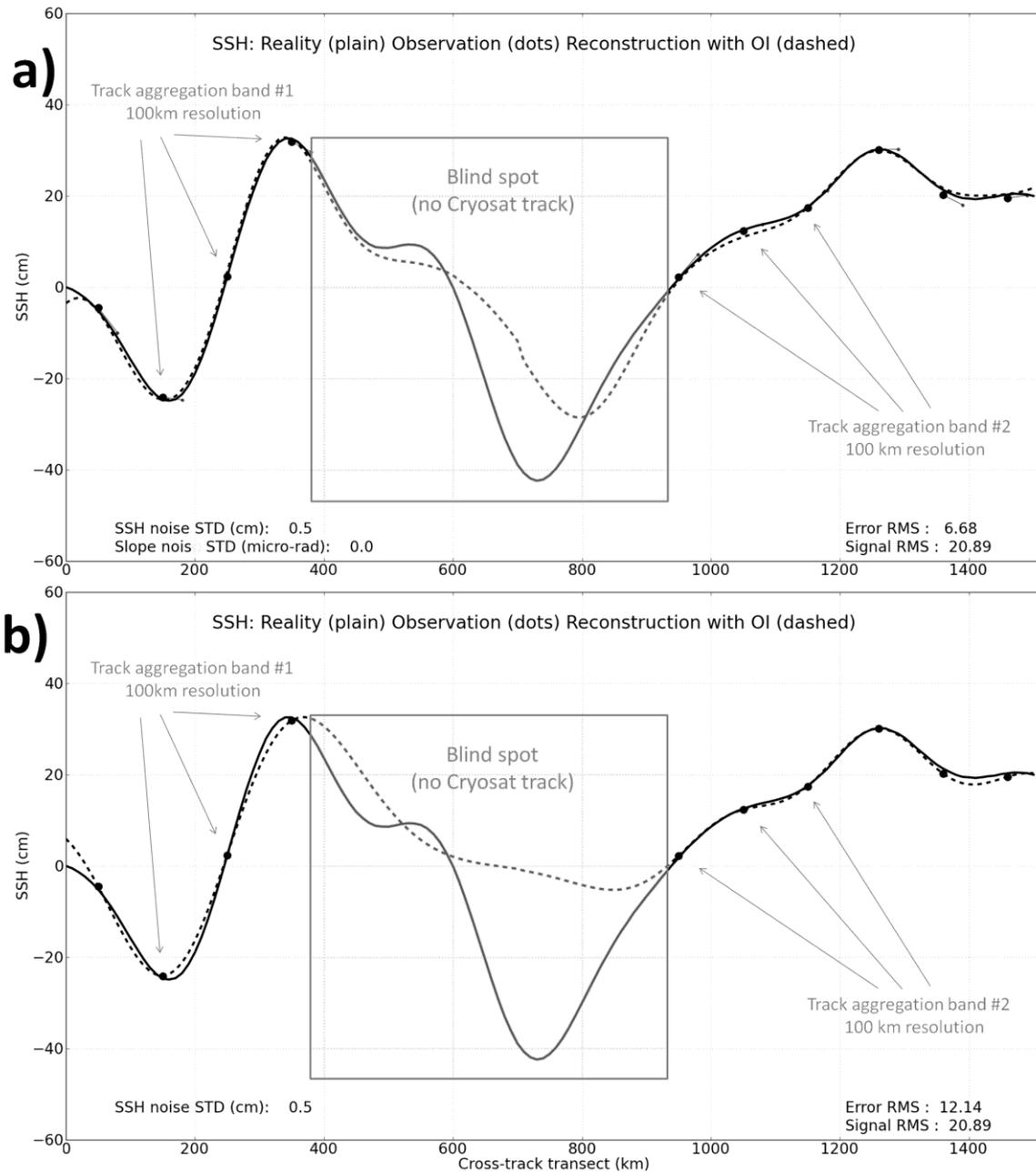
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Radius	Mapping Improvement		Reality signal properties	
	Cryosat-2 Orbit	TP/Jason Orbit	SSH STD (cm)	SSH Slope STD ( $\mu$ rad)
150	0%	18%	20.0	1.9
125	1%	9%	15.8	1.8
100	2%	7%	11.8	1.7
75	8%	5%	8.1	1.5
50	14%	5%	4.8	1.4

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