

Effect of Long Waves on Ku-Band Ocean Radar Backscatter at Low Incidence Angles Using TRMM and Altimeter Data

Ngan Tran, B. Chapron, and D. Vandemark

Abstract—This letter uses a large ocean satellite data set to document relationships between Ku-band radar backscatter (σ_0) of the sea surface, near-surface wind speed (U), and ocean wave height (SWH). The observations come from satellite crossovers of the Tropical Rainfall Mapping Mission (TRMM) precipitation radar (PR) and two satellite altimeters, namely: 1) Jason-1 and 2) Environmental Satellite. At these nodes, we obtain TRMM clear-air normalized radar cross-section data along with coincident altimeter-derived significant wave height. Wind speed estimates come from the European Centre for Medium-Range Weather Forecast. TRMM PR is the first satellite to measure low incidence Ku-band ocean backscatter at a continuum of incidence angles from 0° to 18° . This letter utilizes these global ocean data to assess hypotheses developed in past theoretical and field studies—namely that variations in ocean sea state are measurably and systematically related to Ku-band σ_0 , that the impact changes with incidence angle, and that it will affect the retrieval of wind speed from σ_0 . Results have bearing on near-nadir ocean radar missions such as Surface Waves Investigation and Monitoring from Satellite, Advanced Scatterometer, TRMM, and the wide-swath altimeter.

Index Terms—Author, please supply your own keywords or send a blank e-mail to keywords@ieee.org to receive a list of suggested keywords.

I. INTRODUCTION

SATELLITE radars are used to infer the wind speed just above the sea surface through their measurement of backscattered signal power, which is a signal that changes with the amount and steepness of ocean waves. This normalized radar backscatter cross-section (σ_0) term also depends on the frequency, polarization, and incidence angle (θ) of the incident radiation. Two now-standard satellite systems for ocean wind estimation are the altimeter and the scatterometer. The former views the sea from a downlooking ($\theta = 0^\circ$) incidence angle, whereas the latter uses side-looking angles from 20° to 60° . It is widely held that centimeter-scale ocean gravity-capillary waves and their growth or decay with wind forcing are the dominant controls of σ_0 variation for both sensors, but the ocean reflection is distinctly different for these two systems that is consistent with the optical expectation; increased wave roughness

decreases altimeter σ_0 but increases it for the scatterometer. This is because the incidence angle leads to differing scattering processes: dominant quasi-specular scattering for the former and Bragg resonance diffraction for the latter. Regardless of such differences, the linkage between σ_0 and wind forcing is used for both sensors to empirically derive wind speed inversion algorithms that are well validated and widely used. However, long-wavelength tilting of short-scale waves is a known effect inducing fundamental perturbations in the precise relationship between local wind forcing and local radar backscatter variations. What is central for this letter is the acknowledgement that long-wave tilting of these short-scale waves is an additional second-order but fundamental perturbation that can impinge on any assumed direct relation between local wind forcing and radar backscatter variation, particularly at near-nadir incidence angles. A substantial fraction of the longer tilting gravity wave field is due to swell and wind seas generated by distant or turning winds, which are uncoupled and misaligned with the local wind.

Previous investigations have used bulk wave statistical parameters such as significant wave height (SWH) to demonstrate long-wave variability impacts upon σ_0 [15], [22], [23], [29] and more widely on retrieved winds [12], [13], [21], [27], [28]. Such observations have clearly shown long-wave effects on altimeter backscatter and have led to the development of an operational wind speed model for the satellite altimeter that utilizes both σ_0 and SWH [7], [13], where fortuitously both measurements are made from the same platform. SWH is not retrievable using scatterometry.

While altimeter ocean backscatter has been successfully modeled with quasi-specular scattering theory, off-nadir radar backscatter represents a mixture of specular and tilted Bragg resonance diffraction processes as the incidence angle extends away from 0° out toward 10° – 15° . The transition between the two scattering regimes depends upon the instrument wavelength and the wind speed and has been proposed to occur near an incidence angle of 10° . A notable observation is that close to this angle a lower sensitivity between σ_0 and wind speed is found [14], [19], [25]. This particular feature has been exploited over the ocean to calibrate airborne and spaceborne precipitation or cloud radars—the objective being to minimize uncertainty due to surface wind variations.

The low, or near-nadir, incidence angle range of $1^\circ \leq \theta \leq 8^\circ$, is currently covered by the precipitation radar (PR) on the Tropical Rainfall Mapping Mission (TRMM) [17], [18]. Though designed specifically for the measurement of precipitation profiles in the atmosphere over both land and ocean, 91

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92 the PR system also acquires sea surface σ_0 under rain-free
 93 conditions. Using wind speed estimates from the TRMM Mi-
 94 crowave Imager (TMI), a fully empirical model function was
 95 built to relate cross section to wind speed for incidence angles
 96 from 0° to 18° [11]. This is the first and only satellite system
 97 to provide such angle-resolved scattering near nadir, and the
 98 objective here is to further examine these data to help bridge
 99 what is known regarding the effects of waves on the altimeter
 100 and scatterometer. In this letter, we take the advantage of a
 101 large collocated database, which is compiled using PR and both
 102 Jason-1 and Environmental Satellite (ENVISAT) altimeters,
 103 to extend the description of PR σ_0 in terms of wind speed,
 104 significant wave height, and incidence angle through a tabulated
 105 model function $\sigma_0(\theta, U, SWH)$. This provides a compact and
 106 statistically accurate representation permitting the study of the
 107 expected wave tilting impacts on the sea surface scattering at
 108 these low incidence angles.

109 II. DATA SETS

110 A. TRMM PR Cross Section

111 The TRMM satellite was launched in November 1997 car-
 112 rying five instruments including the PR. Since the focus of
 113 TRMM is to measure rainfall in the tropics, a low inclination
 114 non-sun-synchronous orbit was selected to confine the satellite
 115 ground track between 35°S and 35°N . The PR is a Ku-band
 116 pulsed radar operating at 13.8 GHz and horizontal polarization.
 117 σ_0 measurements are collected through the atmospheric column
 118 and from the surface. The PR antenna is an electronically
 119 scanned phased array that scans a plane normal to the flight
 120 direction (cross-track) through the nadir with measurements at
 121 49 beam positions (e.g., the angle bins 1, 25, and 49 correspond
 122 to the incidence angles $+18^\circ$, 0.1° , and -18° , respectively) over
 123 a 215-km ground swath. The scan duration is equal to 0.6 s with
 124 a surface pixel provided every 4.3 km both along and cross-
 125 track [16], [17] for the original orbit height.

126 The TRMM orbit was raised from 350 to 403 km in Au-
 127 gust 2001 to increase the duration of the mission. The spatial
 128 resolution of the PR is thus degraded slightly, increasing to
 129 5.0 km by 5.0 km. Our data analysis covers the one-year
 130 period of 2003. The high quality of the PR surface ocean σ_0
 131 data for this period was confirmed in two recent studies [11],
 132 [31]. The data product used herein is TRMM PR standard
 133 product 2A21 (ver. 5) from the Goddard Distributed Active
 134 Archive Center. These data include normalized radar cross-
 135 section measurements, associated quality flags, and a rain/no-
 136 rain flag for each incidence angle bin or pixel [18]. Data over
 137 land, with any data quality issue, or with rain over the ocean
 138 target are all excluded from the composite data set. Further data
 139 processing and satellite-to-satellite crossover selection details
 140 follow [31] except that for this letter the search was performed
 141 over all incidence angles in the PR ground tracks. As shown
 142 in the previous study, the density of crossovers increases with
 143 latitude due to the combined altimeter-PR orbit characteristics.

144 B. Wind Speed and Significant Wave Height Data

145 We use surface wind speed estimates (U) from the surface
 146 model analysis provided by the European Centre for Medium-

Range Weather Forecast (ECMWF) as a common reference to 147
 quantify PR σ_0 wind dependence. SWH data for the study come 148
 from the Jason-1 and ENVISAT altimeters. These satellites 149
 also provide an estimate of wind speed using altimeter σ_0 150
 measurements, but ECMWF model winds are used for the 151
 model functions developed here explicitly because we do not 152
 wish to introduce the known sea state impacts that lie within 153
 altimeter wind speed data into the present PR results. The 154
 potential negative impact of using the model wind products is 155
 that these data are extracted and interpolated from six hourly 156
 1° grid data set and that model winds will always disagree 157
 with *in situ* measurements to a certain degree. Thus, the model 158
 functions to be developed will be slightly impacted, particularly 159
 at lightest wind speed, by this interpolation but previous studies 160
 (e.g., [13]) have shown that the systematic nature of wave 161
 height impacts should still be quite apparent and similar when 162
 using the ECMWF model winds and it is this impact that is 163
 the main focus of this letter. While one could go another step 164
 to gather TRMM/scatterometer/altimeter triplet crossovers to 165
 replace ECMWF winds with those from scatterometry, this 166
 step dramatically reduces the data set size without dramatically 167
 increasing the quality of the result as the agreement between 168
 the ECMWF and scatterometer product is high. 169

Time/space interpolated ECMWF wind speed and standard 170
 altimeter SWH estimates are both available in the Geophysical 171
 Data Records (GDR) for these two altimeters. The Jason-1 172
 altimetric mission was launched in December 2001 and placed 173
 in the same ground track as its predecessor TOPEX/Poseidon. 174
 It carries the Poseidon-2 altimeter that was derived from the 175
 experimental Poseidon-1 instrument aboard TOPEX/Poseidon. 176
 The satellite flies a nonsun-synchronous orbit at an altitude 177
 of 1336 km with an inclination of 66° . Detailed description 178
 of the mission and the Poseidon-2 instrument are provided, 179
 respectively, in [20] and [5]. The ENVISAT altimeter (called 180
 Radar Altimeter 2) was launched on March 2002 and is derived 181
 from the European Remote Sensing satellite (ERS)-1 and ERS- 182
 2 altimeters [2]. The satellite orbit is sun-synchronous at an 183
 altitude of 800 km with an inclination of 98.55° allowing 184
 measurement closer to the poles than Jason-1. More details can 185
 be found in the ENVISAT product handbook [3]. Parameters 186
 from both Jason-1 and ENVISAT GDRs over the one-year 187
 period of 2003 are used for this letter. Erroneous altimeter 188
 estimates are discarded using conventional data quality flag- 189
 ging [3], [24]. Further data filtering follows from the Cal/Val 190
 quality assessment that is routinely performed at the Collecte 191
 Localisation Satellites [8]. We use only rain-free data. Since 192
 the Jason-1 rain flag, which is currently available at the time of 193
 this analysis, uses a TOPEX-derived algorithm that was not yet 194
 fine-tuned on Jason-1 measurements, a Jason-1 rain flag was 195
 calculated using a more recent algorithm [30]. This algorithm 196
 shows higher sensitivity to low intensity rainfall as shown using 197
 TMI rain estimates [33]. 198

199 C. Crossover Selection

The criteria used for the collocation between PR and Jason-1 200
 or ENVISAT crossovers are given as follows: time separation 201
 within 1 h and spatial separation less than 100 km. The different 202
 collocation sets PR/altimeter/ECMWF are limited in latitude 203
 to the tropics within $\pm 35^\circ$ of the equator due to the TRMM 204

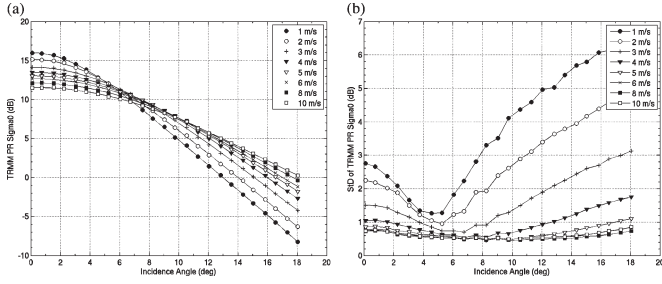


Fig. 1. (a) Mean values and (b) standard deviations of binned PR σ_0 as a function of incidence angle for different wind speeds (SWH between 0.5 and 6.5 m).

205 orbit. We merge the two data sets using, respectively, Jason-1
 206 and ENVISAT SWH estimates to obtain a unique data set
 207 over which the geophysical model function $\sigma_0(\theta, U, SWH)$ can
 208 be produced. To insure homogeneity and consistency between
 209 altimeter SWH estimates for the two missions, we applied small
 210 [O (cm)] SWH adjustments per the most recent correction
 211 model [26].

212 III. NEAR-NADIR SCATTERING MODEL

213 Following the standard quasi-specular backscattering ap-
 214 proach, near-nadir σ_0 can be written as

$$\sigma_0(\theta, U, SWH) = \frac{\rho(U)}{\text{mss}(U, SWH)} \sec^4(\theta) \exp\left[-\frac{\tan^2(\theta)}{\text{mss}(U, SWH)}\right] \quad (1)$$

215 where σ_0 is the normalized backscatter in natural units (not in
 216 decibels), and θ is the incidence angle as previously defined.
 217 ρ represents an effective nadir reflection coefficient, and mss
 218 is a measure of the effective mean square slope (see [4] for
 219 review). The model assumes that sea state dependence of ρ is
 220 unlikely or negligible, which is verified to a large extent using
 221 the dual frequency capabilities of the TOPEX altimeter [6], [9].
 222 The model also allows for the impact of sea state on the cross
 223 section. It is the overall degree of sea state development, which
 224 contributes to the mean squared tilting slopes. Yet, the analogy
 225 with optical scattering assumption implies that the incident
 226 radiation wavelength should be much shorter than all roughness
 227 lengths on the surface. For microwave probing of the ocean
 228 surface, this is untrue due to the presence of gravity-capillary
 229 waves. In particular, at first order and with a Gaussian statistical
 230 assumption, the slope variance in (1) corresponds to a filtered
 231 slope distribution [35].

232 As obtained in Fig. 1(a) (see also [11, Fig. 5]), the Gaussian
 233 assumption of (1) is qualitatively consistent with the PR data
 234 up to about 18° . Observed biases at nadir between altime-
 235 ter measurements and PR data can be attributed to absolute
 236 calibration issues [31]. According to (1), analysis of a single
 237 frequency radar altimeter with both wind speed and sea state
 238 proxy cannot, with certainty, separate the dependencies related
 239 to mss variations from those related to variations of ρ . This is
 240 because at nadir, (1) becomes

$$\sigma_0(\theta = 0, U, SWH) = \frac{\rho(U)}{\text{mss}(U, SWH)}. \quad (2)$$

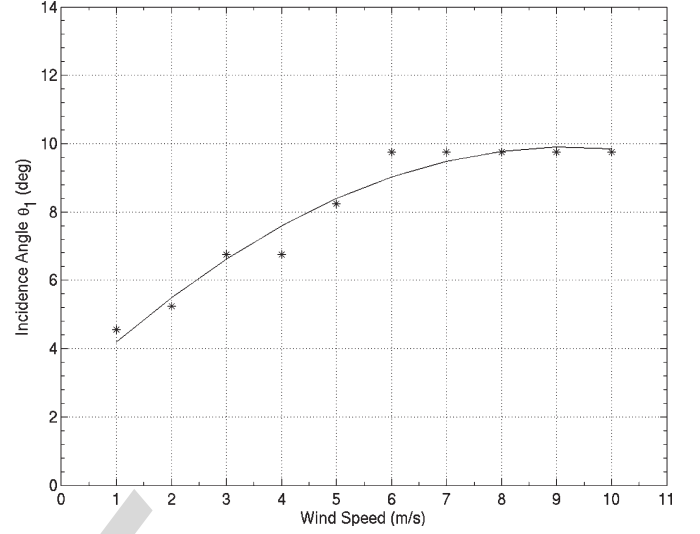


Fig. 2. Incidence angle θ_1 presenting the lowest standard deviation of binned PR σ_0 at a given wind speed as function of wind speed. Overlaid is a quadratic regression fit to better display the trend.

Perhaps more interestingly, the differentiation of (1) with
 respect to mss yields

$$\frac{\partial \sigma_0}{\partial \text{mss}} = \frac{\tan^2(\theta) - \text{mss}}{\text{mss}^2} \sigma_0. \quad (3)$$

The form of the fractional cross-section variation ($\Delta\sigma_0/\sigma_0$)
 in natural units (not in decibels), due to fractional change of
 mss, will be incidence angle dependent, i.e., 1) when $\tan^2(\theta) <$
 mss, $\Delta\sigma_0/\sigma_0 \propto (-\Delta\text{mss}/\text{mss})$ and the nadir viewing altimeter
 falls in this category, and 2) at higher incidence angles, when
 $\tan^2(\theta) >$ mss, $\Delta\sigma_0/\sigma_0 \propto (+\Delta\text{mss}/\text{mss})$ and the off-nadir
 viewing scatterometer falls into this category. The following
 analysis of the PR σ_0 documents this fractional change of σ_0
 with incidence angle.

IV. ANALYSIS OF PR BACKSCATTER

A. Geophysical Model Function for Ku-Band Ocean σ_0 at Low Incidence Angles

We restrict this letter to light-to-moderate wind speed condi-
 tions up to 11 m/s. At wind speeds above this range, complex
 nonlinear surface wave structure and foam involved with large-
 scale wave breaking become critical to the surface description
 and the radar scattering from it. While these higher winds are
 important, the extensive amount of data that fall at or below
 11 m/s and the physics associated with these conditions are the
 focus of this letter. Two empirical tabular model functions are
 developed. The first model is based on the analysis of measured
 σ_0 at each PR incidence angle within specified wind speed
 intervals and is denoted $\sigma_0(\theta, U)$. The 25 different incidence
 angles are given as follows: 0.1° (nadir), 0.75° , 1.55° , 2.25° ,
 3.05° , 3.75° , 4.55° , 5.25° , 6.05° , 6.75° , 7.55° , 8.25° ,
 9.05° , 9.75° , 10.55° , 11.30° , 12.05° , 12.85° , 13.55° ,
 14.35° , 15.05° , 15.85° , 16.55° , 17.35° , and 18.05° , and the bin width is about
 0.1° . The model is formed from the sample mean σ_0 in each 1-
 270 m/s wind speed and incidence angle 2-D bin. A 3σ filter is then
 applied to eliminate outlier measurements, giving Fig. 1(a). The
 second model function takes into account both wind speed and

274 significant wave height dependence at each incidence angle. It
 275 is denoted as $\sigma_0(\theta, U, SWH)$. The wind speed bin width is still
 276 1 m/s, and the SWH bin width is set to 1 m.

277 B. $\sigma_0(\theta, U)$

278 Fig. 1(a) shows that results from nadir to 5° in incidence
 279 angle are monotonically decreasing in σ_0 as wind speed in-
 280 creases. Above 10° , σ_0 becomes a monotonically increasing
 281 function of wind speed. In the range $5^\circ \leq \theta \leq 10^\circ$, σ_0 first
 282 increases, then decreases with increasing wind speed with a
 283 low sensitivity to wind speed. The standard deviations of the
 284 σ_0 measurements in each (θ, U) bin are shown in Fig. 1(b)
 285 with respect to incidence angle for different wind speeds. For
 286 all wind speeds, standard deviations reach a minimum value
 287 at an incidence θ_1 between 4° and 10° . Higher magnitudes
 288 of standard deviation are associated with light wind speeds,
 289 and these magnitudes decrease with increasing wind speed.
 290 Magnitudes are smaller at nadir (0.1°) than at 18° for light
 291 wind speeds up to 5 m/s. Above 5 m/s, results show similar
 292 values. In the range $4^\circ \leq \theta \leq 10^\circ$, σ_0 not only exhibits low
 293 sensitivity to wind speed but also an overall low variability. This
 294 lowered variability is related to (3). The angle θ_1 , in Fig. 2,
 295 roughly identifies the condition $\tan^2(\theta) = mss(U)$ for which
 296 the fractional cross-section variation is minimum, and the shift
 297 of θ_1 with wind speed corresponds to the anticipated increase of
 298 mss. As found, there is an increase of θ_1 with increasing wind
 299 speed up to 7 m/s followed by a saturation trend toward $\sim 10^\circ$
 300 for higher moderate winds.

301 C. $\sigma_0(\theta, U, SWH)$

302 The very large collocated data set compiled enables the
 303 analysis of the combined incidence angle and SWH depen-
 304 dencies on σ_0 using the narrow 1-m/s wind speed bin. Fig. 3
 305 displays a difference factor δ defined as $[\sigma_0(\theta, U, SWH) -$
 306 $\sigma_0(\theta, U)]$, in decibels, with respect to incidence angle at four
 307 selected wind speeds of 2, 5, 7, and 10 m/s. For all winds,
 308 behavior of δ as a function of SWH is clear. At low SWH
 309 (~ 1 m) representing young sea, δ decreases with increasing
 310 angle, whereas for higher SWH (~ 4 m associated mostly with
 311 mixed seas including swell) δ exhibits the opposite trend. The
 312 overall picture shows that at a given wind speed, all curves
 313 (linear least-squares fits) associated to the different 1-m SWH
 314 classes intersect at a particular value of incidence angle θ_2 that
 315 shifts with respect to wind speed value.

316 Very similar results are obtained when reducing the crossover
 317 collocation criteria. A subset is extracted for measurements
 318 collocated in time to within 1/2 h and 25 km in space. Fig. 4
 319 is the same as Fig. 3 but only for the case of two moderate
 320 wind speeds (7 and 10 m/s) for which there is still sufficient
 321 data (minimum of 100 samples per bin) to compute statistically
 322 stable indicators.

323 The relative magnitude of σ_0 for extreme conditions, i.e., low
 324 and high SWH (1 and 4 m, respectively), is shown in Fig. 5 as a
 325 function of incidence angle for light-to-moderate wind speeds.
 326 For all wind speeds, we observe a positive magnitude at low
 327 incidence angles that decreases to reach a negative value at
 328 higher incidence. At 2-m/s wind, the magnitude is, respectively,
 329 ~ 0.8 dB at nadir and -1.6 dB at $\sim 18^\circ$. At 10 m/s, we ob-
 330 serve almost similar absolute magnitude of variation (~ 0.8 dB)

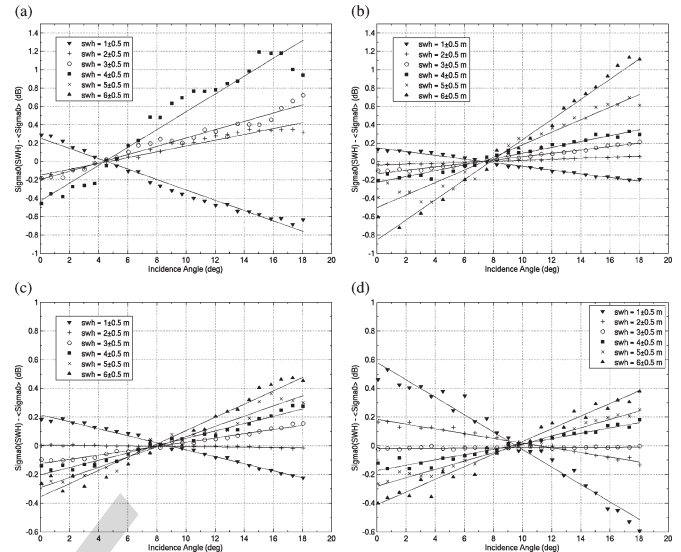


Fig. 3. Difference δ (between averaged PR σ_0 associated to a 1-m class of SWH and the averaged values estimated over all SWH) as function of incidence angles for various SWH classes at selected wind speeds: (a) 2 m/s, (b) 5 m/s, (c) 7 m/s, and (d) 10 m/s (1-m/s bin width). Overlaid are linear regression fits to better display the trends.

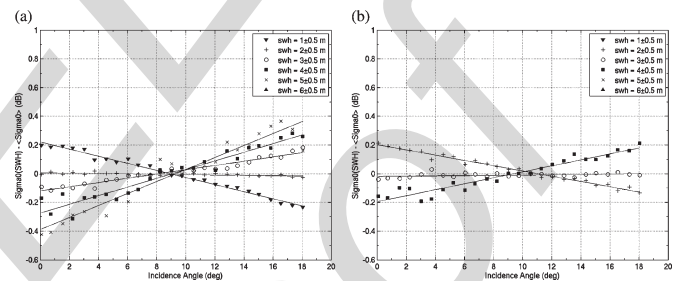


Fig. 4. Same as Fig. 3 but crossover data are selected with narrower selection criteria, 25 km and 30 min, for two selected wind speeds: (a) 7 m/s and (b) 10 m/s (1-m/s bin width).

between the two extreme incidence angles. These results are 331
 consistent with the previous analysis at higher incidence angles 332
 (20° , 30° , 40° , and 60°) in [23] using NUSCAT-SWADE data. 333
 However, in cases of moderate winds, these authors concluded 334
 that the existence of large waves with high SWH will not have 335
 significant impact on the radar backscatter since the observed 336
 differences were within the uncertainty of the radar (± 1 dB). 337
 The large amount of data available here helps to revise these 338
 conclusions. For these wind conditions, the presence of large 339
 waves significantly impact σ_0 from nadir to 18° except around 340
 a particular incidence angle, denoted θ_2 in Fig. 6, where σ_0 is 341
 insensitive to SWH at a given wind speed. As found, θ_1 and 342
 θ_2 angles are almost equal and correspond to the condition 343
 $\tan^2(\theta_1) = \tan^2(\theta_2) = mss(U)$. Around these critical angles, 344
 the backscatter cross section is insensitive to significant wave 345
 height variations at a given wind speed. 346

One point of note for these TRMM PR data is that they 347
 represent horizontally polarized returns. There is a recognized 348
 difference in the response of horizontal and vertical polarization 349
 returns from the sea surface (cf. [1]). The present results, in 350
 terms of overall features, can however be easily transposed to 351
 a vertically polarized result. Indeed, continuity between nadir 352
 viewing returns (no polarization) and scatterometer off-nadir 353
 returns in either one of the polarized states indicates that since 354

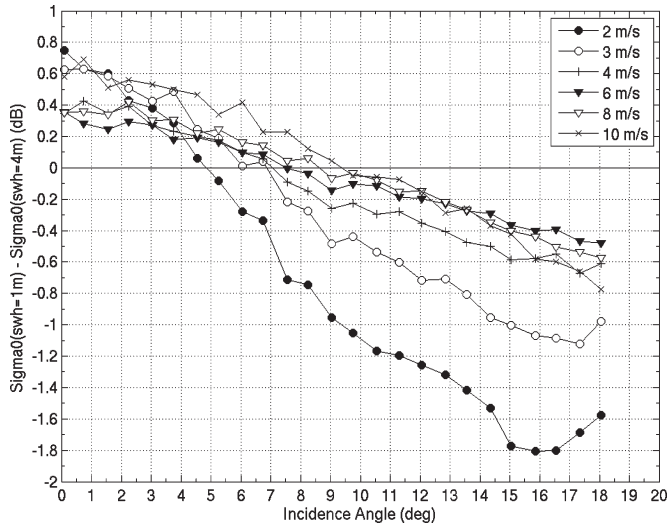


Fig. 5. Magnitude of the difference of σ_0 between low SWH (1 m) and high SWH (4 m) conditions as function of incidence angle for different wind speeds from light to moderate winds.

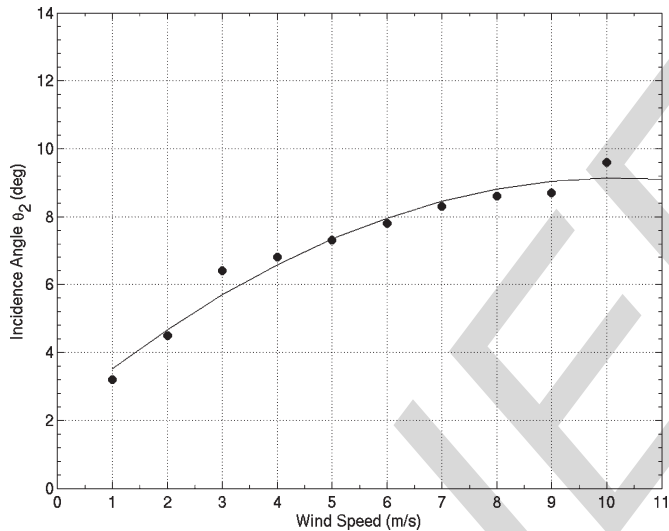


Fig. 6. Incidence angle θ_2 presenting a quasi-insensitivity of PR σ_0 to SWH at a given wind speed as function of wind speed. Overlaid is a quadratic regression fit to better display the trend.

355 sea state effects are observed in both polarizations when the
 356 backscatter is off nadir, all near-nadir measurements will dis-
 357 play the same trends (in HH and VV). Previous analysis of the
 358 National Aeronautics and Space Administration scatterometer
 359 (NSCAT) backscatter in each polarization state shows similar
 360 relative sea state impacts with respect to a global averaged
 361 backscatter that was derived by mixing all sea state conditions;
 362 they are slightly larger for NSCAT HH polarization measure-
 363 ment than on VV polarization data regardless of incidence
 364 angles between 16° and 50° [32].

365

V. CONCLUSION

366 New approaches for viewing the global ocean using satellites
 367 have become available in the last decade. This letter focuses on
 368 sea surface roughness remote sensing and what can be learned
 369 using a multiple satellite perspective with the specific goal
 370 being to provide new data to bridge the gap between what

is known about nadir and off-nadir microwave scattering and
 371 emission from the ocean. Near-surface wind speed is a first-
 372 order geophysical parameter to be derived from microwave
 373 ocean sensors (the scatterometer, radiometer, and altimeter), but
 374 it is well known that the transfer function between their raw
 375 measurements and wind speed must account for perturbation
 376 due to surface wave processes that often deviate from simple
 377 local wind forcing behavior. To reduce uncertainties in satellite
 378 wind speed retrieval from backscatter measured at different
 379 observation angles and to ensure proper assimilation of scat-
 380 terometer and altimeter data into numerical weather prediction
 381 models, it is increasingly apparent that a precise understanding
 382 of the relationship between surface roughness through radar
 383 backscatter measurement and both wind and wave conditions
 384 is very important [28], [29]. Gaining quantitative insight on
 385 these sea state perturbations using field studies is notoriously
 386 difficult due to the inability to gather the sufficient range of
 387 surface conditions and data population.

388
 This letter makes use of a multisatellite ocean observing
 389 opportunity, where a new type of ocean surface remote sensing
 390 data set, i.e., the TRMM cross-track scanning radar, is com-
 391 bined with coincident sea surface wave height information from
 392 crossing satellite altimeters to provide all-new data illustrat-
 393 ing wave impacts on radar backscatter at multiple incidence
 394 angles. The resulting TRMM PR model function provides
 395 results showing that long-wave tilting effects are quantitatively
 396 confirmed in line with recent airborne slope measurements [34].
 397 Accordingly, near-nadir cross-section measurements at a given
 398 fixed wind speed and ranging in incidence angles out to 20° are
 399 measurably related to the sea state dynamics. As a surrogate for
 400 the sea state's degree of development, the use of a collocated
 401 SWH parameter helps to document this impact and to clearly
 402 identify the off-nadir incidence angle that corresponds to the
 403 lowest fractional cross-section variation—a very useful angle to
 404 know in over-ocean radar calibration activities. For TRMM PR
 405 incidence angles that lie closer to scatterometer viewing angles
 406 (i.e., 16° – 20°), our results show that for light-to-moderate
 407 wind conditions the presence of large waves can affect the perfor-
 408 mance of surface wind retrieval algorithms, which is consis-
 409 tent with previous results (e.g., [27] and [29]). As expected,
 410 larger incidence angles are thus certainly to be recommended
 411 for surface wind scatterometry to minimize sea state impact.
 412 Combined use of both nadir and near-nadir single frequency
 413 measurements can also help to infer a sea surface slope variance
 414 that can potentially be related to surface wind stress [35] and
 415 assimilated into numerical wave models. As obtained, this sea
 416 surface slope variance will include both longer wave and shorter
 417 wave slope contributions. Dual frequency nadir measurements
 418 and/or use of the contemporaneous SWH measurements will
 419 then help to remove the longer wave contributions to leave the
 420 shorter ones [9]. At nadir and near-nadir configurations, dual
 421 frequency capability will thus improve short surface wave ob-
 422 servations and surface wind retrieval algorithm performances.

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424

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429

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