

Exploring the Benefits of Using *CryoSat-2*'s Cross-Track Interferometry to Improve the Resolution of Multisatellite Mesoscale Fields

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Abstract:

Sea surface height (SSH) measurements provided by pulse-limited radar altimeters are one-dimensional profiles along the satellite's nadir track, with no information whatsoever in the cross-track direction. The anisotropy of resulting SSH profiles is the most limiting factor of mesoscale SSH maps that merge the 1D profiles.

This paper explores the potential of the cross-track slope derived from the *Cryosphere Satellite-2* (*CryoSat-2*)'s synthetic aperture radar interferometry (SARin) mode to increase the resolution of mesoscale fields in the cross-track direction. Through idealized 1D simulations, this study shows that it is possible to exploit the dual SARin measurement (cross-track slope and SSH profile) in order to constrain mesoscale mapping in the cross-track direction.

An error-free SSH slope allows a single SARin instrument to recover almost as much SSH variance as two coordinated altimeters. Noise-corrupted slopes can also be exploited to improve the mapping, and a breakthrough is observed for SARin errors ranging from 1 to 5 μrad for 150-km-radius features in strong currents, and 0.1–0.5 μrad for global mesoscale.

Although only limited experiments might be possible with the error level of current *CryoSat-2* data, this paper shows the potential of the SAR interferometry technology to reduce the anisotropy of altimeter measurements if the SARin error is significantly reduced in the future, and in particular in the context of a prospective SARin demonstrator optimized for oceanography.

Keywords: Altimetry, Remote sensing, Sampling, Satellite observations, Interpolation schemes, Inverse methods

1. Introduction and context

In contrast with wide-swath imagers (e.g., sea surface temperature or ocean color), the data record of radar altimeters is exceedingly anisotropic. Sea surface height (SSH) measurements from pulse-limited radar altimeters are one-dimensional profiles along the satellite's nadir

41 track, with no SSH information whatsoever in the cross-track direction. Figure 1 shows that,
42 for a single altimeter flying on the TOPEX/Jason orbit, the along-track (white segment)
43 resolution can be as small as 7 km (level 2 product, 1 Hz rate), whereas in the cross-track
44 resolution (black segment) it can be as large as 300 km.

45 To reconstruct 2D gridded fields of SSH or sea level anomalies (SLA), it is therefore
46 necessary to interpolate 1D profiles (e.g. AVISO 2010; Dibarboure et al, 2011a, Leben et al
47 2002, 2011). Optimal interpolation (OI) exploits an a priori statistical knowledge of the SLA
48 field characteristics (e.g. Jacobs et al 2001; Le Traon et al, 2003) and measurement error (e.g.
49 Philipps et al, 2012) as an additional constraint to merge 1D profiles from multiple sensors in
50 an optimal way (e.g. Ducet, et al 2000).

51 The anisotropy of SSH profiles is by far the most limiting factor of gridded SSH mesoscale
52 fields (Le Traon & Dibarboure, 2002, 2004; Pascual et al 2006), and especially in near real
53 time where measurements “from the map’s future” are not yet available (Pascual et al, 2008).
54 There are two practical consequences to this limitation.

55 Firstly, even if the spatial and temporal scales used to constrain the OI are derived from SSH
56 measurement of 2 to 4 satellite constellations, the mapping is limited in the cross-track
57 direction. Because 1D profiles from multiple sensors are blended into one map, 2D mesoscale
58 mapping uses a compromise between actual mesoscale correlations and the sampling
59 limitations from such constellations (Ducet, et al, 2000).

60 The resolution of mesoscale fields is dominated by the number of altimeters in operation.
61 Chelton & Schlax (2003), Le Traon & Dibarboure (2002), and Chelton et al (2011) have
62 shown that mesoscale maps have a limited global resolution capability. Higher resolution can
63 still be achieved, but only locally, at certain times, when enough 1D profiles are available
64 (Dussurget et al, 2011).

65 In this context, a new technology used on CryoSat-2 has the unprecedented potential to add
66 actual measurements to constrain mesoscale mapping in the *cross-track* direction. Indeed, in
67 addition to a classical pulse-limited radar altimeter measurement (also known as low
68 resolution mode or LRM), CryoSat-2's altimeter SIRAL features a synthetic aperture radar
69 interferometry (SARin) mode able to measure the SSH slope in the cross-track direction
70 (Francis et al, 2007) as illustrated by Figure 2. In this paper, the cross-track slope (CTS) is
71 given in micro-radians: a 1 μrad slope is approximately equal to a SSH gradient of 1 cm for
72 10 km, or a geostrophic current of 10 cm/s at mid-latitudes.

73 In this paper we use idealized OI simulations to explore the potential of the CTS derived from
74 CryoSat-2's SARin mode to increase the resolution of mesoscale fields in the cross-track
75 direction (methodology introduced in section 3). Our approach is to look at SARin technology
76 in optimal conditions in section 4 and then to discuss what can be done in practice with
77 current and future datasets in section 5.

78 **3 Methodology**

79 **3.1 Overview**

80 Le Traon and Dibarboure (2002, 2004), Chelton et al (2003) and Dibarboure et al (2011a)
81 have shown that 2D SSH mapping is affected by many parameters (e.g. geometry, phasing or
82 coordination of the constellation's orbits, high frequency ocean dynamics). To measure the
83 potential of using a SARin slope to constrain mesoscale mapping we therefore use a simpler
84 idealized 1D configuration.

85 We specifically focus on the cross-track direction (black segment from Figure 1) where the
86 resolution is limited by the number of satellites in the constellation. In other words, this is a
87 configuration where SARin slopes are ideal to complement lacking SSH measurements.

88 While performed for a given cross-track resolution (i.e. latitude), general conclusions can be
89 derived from our analysis because correlation scales were shown to decrease with latitude as
90 well (e.g. Jacobs et al 2001; Le Traon et al, 2003).

91 We measure the performance of mesoscale mapping using the following protocol:

- 92 • we simulate a mesoscale SSH “reality” profile, and we consider that the reality
93 profile is in a frozen state, i.e. stationary over the 10-day period of a T/P or Jason
94 repeat cycle (this strong assumption is discussed in section 5.3)
- 95 • the reality SSH field is sampled on measurement points to create error-free
96 observations,
- 97 • errors are optionally added to the observations,
- 98 • observations are injected into a 1D optimal interpolator to create a “reconstructed”
99 mesoscale field at the original resolution.

100 In this process, the statistical variance and correlation scales of the reality field are known
101 analytically. Consequently the reconstruction is perfect if performed from enough error-free
102 observations. In other words, differences between the reality and the reconstructed fields are
103 the result of the omission or sampling error (not enough data to observe the signal) and
104 commission or measurement error.

105 Note that, there is an additional error source in the mapping of real data: the imperfect
106 modeling of signal and error covariances. This point is discussed in section 5.3.

107 **3.2 Methodology**

108 In this paper, we generate our reality H_{real} as a spatially correlated random Gaussian process
109 (Equation 1). The default decorrelation scale is 150 km, i.e. consistent with findings from Le
110 Traon et al (2003). In our first simulations (section 4.1), the oceanic variability used is 20 cm

111 RMS, i.e. we focus on zones of intense mesoscale activity (e.g. western boundary currents).
112 Then we expand to different signal to error ratios in section 4.2. In section 5.3, we discuss the
113 validity of the Gaussian methodology.

$$\text{Corr}(H,H) = \exp (-x^2/d^2) \quad \text{Equation 1}$$

114 Our observation field H_{obs} is then constructed (Equation 2) by interpolating H_{real} at the desired
115 resolution (30 km for an along-track simulation, 300 km for a cross-track simulation on a
116 Jason-like orbit, and 100 km for a cross-track simulation on the CryoSat-2 orbit) and adding a
117 white noise of 0.5 to 2 cm to the interpolated SSH values. This is arguably a pessimistic error
118 level at 100+ km if compared with results from Dorandeu et al 2004, or Ablain et al 2011 : the
119 noise they observe at a 7 km resolution would be reduced by along-track filtering of the SSH
120 (factor of 2 for 30 km super-observations).

$$H_{\text{obs}}(\mathbf{x}) = H_{\text{real}}(\mathbf{x}) + \epsilon(\mathbf{x}) \quad \text{Equation 2}$$

121 Simulations shown in this paper do not include any along-track bias or long wavelength
122 correlated errors as our sensitivity studies show no significant difference with noise-limited
123 simulations. Although not shown in this paper, our simulated 1D mapping is degraded by
124 correlated errors like operational mesoscale 2D mapping (e.g. Dibarboure et al, 2011c), but
125 the anisotropy effect presented in section 3.3 and the impact of using SARin presented in
126 section 3.5 are the same.

127 The reconstruction of the estimated mesoscale field H_{est} is performed with a 1D optimal
128 interpolation derived from Bretherton et al (1976). H_{est} is obtained from Equation 3 where A
129 is the matrix describing the covariance between H_{est} and H_{obs} of Equation 2, and C the matrix
130 describing the covariance between the SSH observations (covariances are derived from
131 Equation 1). The formal reconstruction covariance error matrix E is obtained from Equation

132 4, although in practice only its diagonal is used here (1- σ gray envelope around reconstructed
133 SSH profiles).

$$\mathbf{H}_{\text{est}} = \mathbf{A} \cdot \mathbf{C}^{-1} \cdot \mathbf{H}_{\text{obs}} \quad \text{Equation 3}$$

134

$$\mathbf{E} = \mathbf{I} - \mathbf{A} \cdot \mathbf{C}^{-1} \cdot \mathbf{A}^T \quad \text{Equation 4}$$

135 Many figures shown in this paper are limited to 3000 km segments for the sake of illustration
136 but simulations were performed on very long profiles to ensure that the examples in this paper
137 are representative of the statistical behavior of each configuration.

138 **3.3 Observation anisotropy**

139 Figure 3 shows one reality segment, sampled in the along-track direction (every 30 km) with
140 2 cm white noise added. The reconstructed field after optimal interpolation is almost identical
141 to the reality field. The reconstruction error is 1.2 cm RMS i.e. 0.4% of the reality signal
142 variance (18 cm RMS). Similarly, the along-track slope (bottom panel) is almost perfectly
143 observed in the along-track direction.

144 Figure 4 shows the same reality segment, but positioned as a transect in the cross-track
145 direction (black segment from Figure 1). In other words, each measurement (black dot) is the
146 crossover between the transect and a different satellite track. In this figure, the SSH reality is
147 sampled by a LRM altimeter every 300 km, i.e. the worst case configuration of a TP/Jason
148 orbit. Because the Nyquist criterion is not met with a single satellite, many features are
149 missed entirely in the reconstruction (e.g. at km #1000 or #1800 or #2200). The error
150 reconstruction RMS is 46% of the signal variance. This figure illustrates the inability of a
151 single satellite to observe large mesoscale, let alone features with radii smaller than 150 km.

152 Adding a second LRM altimeter (perfectly coordinated with the first one, i.e. like in the
153 TOPEX/Jason tandem) significantly improves the resolution of the mesoscale field as shown

154 by Figure 5. Although the Nyquist criterion is barely met, the reconstruction error is
155 significantly reduced with an error of 8% of the signal variance. The error is consistent and
156 slightly larger than the 5% obtained by Le Traon et Dibarboure (2002) because this segment
157 represents the widest gap between roughly parallel tracks. These scenarios give the 1-LRM
158 and 2-LRM reference configurations to which SARin experiments can be compared to infer
159 the cross-track slope contribution in an ideal case.

160 The formal mapping error from Equation 4 is visible in each simulation as a grey envelope of
161 vertical bars. This theoretical error is *–for these idealized simulations–* a very accurate
162 statistical estimate of the error which could be made during the reconstruction process: the
163 differences between the real (plain) and the reconstructed (dashed) SSH are consistent with
164 the 1- σ boundaries defined with the grey envelope from Figure 3 and Figure 4.

165 The formal error represents the sum of the measurement error and the sampling error: it is as
166 small as 2 cm near observation points (measurement noise) and as large as tens of centimeters
167 at the center of the 300 km window between satellite tracks (sampling error). The anisotropy
168 of the altimetry system is illustrated by the difference between the along-track and the cross-
169 track formal errors. In the along-track direction (Figure 3) the error is always very small and
170 dictated by the measurement error level whereas in the cross-track direction (Figure 4) the
171 sampling error largely dominates between satellite tracks.

172 **3.4 CryoSat-2's cross-track measurement**

173 CryoSat-2 is ESA's ice mission (Francis, 2007). Equipped with an innovative radar altimeter
174 (SIRAL – Synthetic Aperture Interferometric Radar Altimeter), and high-precision orbit
175 determination (POD), CryoSat-2's primary objective is to serve Cryospheric science
176 (Wingham et al, 2006). Cryosat's altimeter is operated almost continuously over ocean,
177 mainly in LRM (i.e. conventional altimetry) or in the delay doppler / synthetic aperture radar

178 (SAR) mode which provides higher along-track resolution and lower noise level (Raney,
179 1998).

180 Furthermore, SIRAL also features a third mode: the SAR-interferometry (SARin) mode,
181 which uses CryoSat-2's two antennas (Francis, 2007). The combination of SAR and
182 interferometry makes it possible to determine the cross-track slope of the surface from which
183 the echoes are arriving. This is achieved by comparing the phase of one receive channel with
184 respect to the other as first suggested by Jensen, 1999.

185 With the SARin mode, CryoSat-2 can provide one estimate of the local CTS every 0.05
186 seconds, in addition to the classical topography measurement (Figure 2). Moreover, the along-
187 track resolution and the precision of the SSH is the same as for a LRM sensor (e.g. Jason-2).
188 The resolution is 300 m in the along-track direction (synthetic footprint), and the slope is
189 estimated from a cross-track footprint of the order of 7 km.

190 This unprecedented measurement was initially designed to be used over the margins of the
191 Greenland and Antarctic ice sheets, where the surface slopes are steep. To that extent,
192 SIRAL's SARin mode was designed to have a cross-track slope accuracy of 200 μ rad
193 (Wingham et al 2006), but Galin et al (2012) reported a noise level of 20 μ rad at a 7 km
194 resolution and a bias of 8 μ rad for 1000 km segments, using both detailed modeling of the
195 finite radar resolution in range and angle, and the thermally driven behavior of the
196 interferometer bench.

197 This should be compared to the typical mesoscale slope distribution in zones of intense
198 mesoscale activity which ranges from 1 to 5 μ rad at 150 km with values higher than 10 μ rad
199 on the edges of the largest eddies (observed on multi-satellite SSH maps from AVISO, 2010).
200 Assuming that the long wavelength errors described by Galin et al (2012) are minimized with
201 empirical cross-calibration mechanisms (discussed in section 5.2.1), and that the noise level is

202 reduced by along-track filtering (discussed in section 5.2.2), it would become possible to use
203 the SARin slope as a constraint for mesoscale mapping in the cross-track direction where
204 LRM altimeters are blind.

205 Because the error level reported on CryoSat-2 is high with respect to the oceanic signal, our
206 rationale is the following: we first look at the benefits of using error-free SARin CTS (section
207 4.1), then we perform sensitivity studies with respect to the ocean variability and
208 measurement errors (section 4.2). From this background, we discuss the practical case of
209 CryoSat-2 in section 5.

210 **3.5 Improving the reconstruction with the cross-track slope**

211 Figure 2 gives a qualitative illustration of how mesoscale mapping can exploit the SARin
212 cross-track slope. Subplot (a) shows a 500 km along-track LRM profile with SSH only
213 (simulated, error-free), whereas subplot (b) shows the information given by a SARin profile
214 with SSH and cross-track slope. Both plots correspond to the reality from subplots (c) and (d).

215 From the SSH+CTS sample (subplot b), one can assume that the maximum value at -100 km
216 is located on the right-hand side of the nadir track, that the minimum value at +150 km is
217 probably near the nadir track, and that the maximum value at +400 km is located on the left-
218 hand side of the nadir track. Adding a statistical description of mesoscale variability and
219 slopes, it is possible to enhance the mapping in the cross-track resolution up to a distance
220 equal to the spatial correlation radius.

221 This is achieved using a method derived from Le Traon and Hernandez (1992): we replace the
222 SSH observation vector H_{obs} in Equation 3 by a vector composed of all observations (SSH and
223 CTS), and we update matrixes A and C from Equation 3 and Equation 4 accordingly (see
224 Appendix).

225 **4 Results**

226 **4.1 Error-free simulations**

227 In this section we infer what would be the optimal mesoscale improvement using SARin on a
228 300-km cross-track resolution (i.e. a Jason-like orbit). It is optimal for SARin in the sense that
229 Le Traon and Dibarboure (2002, 2004) have shown that the main weakness of this orbit is the
230 cross-track resolution, and it is the reason why TOPEX/Jason and Jason-1/Jason-2 were put in
231 a spatially interleaved configuration. Thus we use this “reference orbit” and “reference
232 tandem” to SARin-based simulations. We discuss the difference between this Jason-like
233 configuration the (suboptimal) Cryosat orbit in section 5.1.

234 Adding the SARin slope constraint (error free) significantly improves the OI reconstruction as
235 shown by Figure 6. This plot should be compared to Figure 4 where one LRM altimeter was
236 barely able to recover 50% of the signal variance in the cross-track direction (Nyquist
237 sampling not achieved). Thanks to local constraints on the SSH derivative, it is possible to
238 recover features that were previously missed entirely (e.g. at km #1800 and #2200).

239 Quantitatively, on this example, the reconstruction error is only 6.96 cm RMS, i.e. 15% of the
240 signal variance (vs. 50% for the LRM scenario on Figure 4). In other words, about 35% of the
241 signal variance was recovered with the error-free slope. The 15% residual error should also be
242 compared to the 8% of the configuration with two LRM altimeters (Figure 5): in this idealized
243 simulation, a single SARin altimeter performs almost like two LRM altimeters.

244 Similarly, Figure 7 shows that a perfectly coordinated constellation of 2 SARin altimeters
245 flying on a Jason-like orbit (150 km cross-track resolution) is able to properly reconstruct the
246 SSH and slope reality fields even though the Nyquist criterion is barely met with SSH alone.
247 Because slopes and covariance models add the constraint needed, the reconstruction error is

248 only 1.83 cm RMS (i.e. 1% of the signal variance) and largely due to the error outlier of the
249 first measurement and the 2 cm SSH measurement noise.

250 **4.2 Sensitivity to signal to noise ratio**

251 We performed a series of sensitivity tests on the slope error for 1 and 2 SAR-in altimeter
252 constellations using very long simulations (2000 times the correlation radius). Figure 8 shows
253 the RMS of the error reconstruction as a function of the standard deviation (STD) of the
254 simulated SARin slope error (plain line). The 1-LRM and 2-LRM references are also given by
255 the black dotted and dashed lines. Note that the observed error is consistent with the formal
256 error given by Equation 4.

257 As expected, the reconstruction error decreases as the CTS error does, and the sigmoid shape
258 on the logarithmic abscissa scale indicates that the largest gains are obtained between 1 and 5
259 μrad , i.e. near the peak of the cross-track slope probability density function.

260 The upper asymptotic value for slope errors higher than 20 μrad is 49%, i.e. the mapping error
261 observed for 1 LRM sensor (dotted line). In other words, if the SARin error is large, it does
262 not improve the reconstruction with the OI. Yet as expected from a theoretical point of view,
263 this figure shows that even if the error STD of the CTS is 25 times larger than the SSH slope
264 STD (i.e. factor of 600 in the covariance matrix), the OI never underperforms w.r.t to the 1-
265 LRM scenario because untrustworthy observations are automatically downweighted by
266 covariance matrix C .

267 If the OI covariance matrixes are properly set up, adding very noisy slope estimates (e.g. 10 to
268 20 μrad unmitigated error from Galin et al, 2012) can still improve the cross-track mapping,
269 albeit in a very limited way.

270 The lower asymptotic value is 13% of the signal variance, i.e. only slightly larger than the 9%
271 error observed with 2 LRM sensors (dashed line): using an error-free SARin instrument in an
272 ideal configuration (1D, cross-track, 150 km radius for a 300 km sampling resolution) does
273 not allow to fully reconstruct the signal, but a single SARin instrument yields results almost
274 as good as two LRM sensors as per the example from Figure 5 and Figure 6. The residuals
275 arise from sampling errors: although additional error-free parameters are used, there are still
276 not enough measurements points to correctly resolve all mesoscale structures.

277 Results are similar for 2 x SARin simulations in Figure 9, even though the gain is more
278 limited owing to the fact that 2 coordinated LRM altimeters already have a good sampling
279 capability for 150 km radius features (Le Traon et Dibarboue, 2004). In this figure, the lower
280 asymptotic value is 1.2%, i.e. very close to the 1% obtained with 4 coordinated LRM sensors:
281 sampling errors would become marginal in a coordinated 2 x SARin configuration.

282 Because the variability of the cross-track slope is proportional to the variability of the SSH,
283 we performed sensitivity studies to the latter (using constant correlation scales and SSH noise
284 levels) to see how results from section 4.1 could be extrapolated out of intense mesoscale
285 activity zones.

286 Figure 10 confirms that the reconstruction error is still sigmoid-shaped, and shifted along the
287 abscissa axis as a function of the SSH variability. The breakthrough in mapping improvement
288 is always achieved for slope error STD ranging from $0.5 \cdot \sigma_{\text{slope}}$ to $2 \cdot \sigma_{\text{slope}}$.

289 To be used globally in mesoscale mapping, SARin slopes would require an error level of the
290 order of 0.1 to 0.5 μrad for mesoscale wavelengths. This is largely beyond what can be
291 achieved with current data from Cryosat-2 (discussed in section 5.2).

292 **5 Discussion: from theory to practice**

293 **5.1 Sensitivity to the satellite track geometry**

294 The sampling pattern of the CryoSat-2 orbit (current SARin mission) and the Jason orbit
295 (simulations from section 4) are very different. The latter has a 10-day repeat cycle (300 km
296 cross-track resolution from Figure 1). In contrast, CryoSat-2 has a one-year repeat cycle with
297 3-day and 30-day sub-cycles, i.e. globally homogeneous sampling patterns with 1000 and 100
298 km cross-track resolutions respectively (Francis et al, 2007). CryoSat-2's orbit has no sub-
299 cycle in the 10 to 20 day range associated with mesoscale temporal decorrelation (Jacobs et
300 al, 2001).

301 As a result, for any 10 to 20 day period, CryoSat-2's measurements are aggregated in band-
302 shaped patterns (100 km wide per 3-day sub-cycle) which are interleaved with band-shaped
303 "blind spots" with no recent SSH observation (Figure 11). The impact on mesoscale
304 observation in LRM mode is discussed by Dibarboure et al (2011c). As far as SARin slopes
305 are concerned, there are two consequences of CryoSat-2's sampling pattern.

306 **5.1.1 Track aggregation and data gaps**

307 Firstly the SARin slopes located on the outer edges of the band-shaped aggregation of satellite
308 tracks provide a unique capability to reduce the extent of the band-shaped blind spots by up to
309 $2 * 150$ km (one slope constraint on each side of the diamond not covered by CryoSat-2
310 tracks in Figure 11). This is useful to balance CryoSat-2's main sampling weakness when it
311 comes to mesoscale observation.

312 Figure 12 illustrates this point: it shows the OI reconstruction for a 1500 km cross-track
313 segment where CryoSat-2 measurements are aggregated in 100 km resolution bands where
314 mesoscale features (150 km radius) are resolved, and interleaved with a 500 km wide blind

315 spot where no CryoSat-2 track is available in the 15 day window corresponding the frozen
316 field approximation.

317 The SARin-based reconstruction (subplot a) is slightly better because the outer edges are
318 constrained by error-free slope estimates whereas the LRM-based reconstruction (subplot b)
319 is not able to observe even a fraction of the large eddy at km #700 and the total reconstruction
320 error is much higher (12.1 cm RMS vs. 6.7 cm RMS for SARin). Note that the overall
321 improvement is limited to the outer edges of the large data gap (one decorrelation radius on
322 each side) because the OI cannot “guess” the existence mesoscale structures if they are not
323 remotely observed.

324 **5.1.2 Orbit sampling differences**

325 The second consequence of CryoSat-2’s sampling pattern is the cross-track resolution *within*
326 the track aggregations. CryoSat-2’s sampling “bands” have a cross-track resolution of 100
327 km, i.e. more favourable to the observation of 150 km radius mesoscale features, albeit in
328 limited areas. In this context, SARin data from Cryosat might be used to recover smaller
329 mesoscale features (only within the satellite track aggregation).

330 Table 1 shows the mapping improvement (i.e. the reduction of cross-track reconstruction
331 error) when the “reality” and OI correlation radiuses range from 50 to 150 km and for the
332 Jason and CryoSat-2 orbits. All simulations were performed with a slope measurement noise
333 of 1 μ rad. On the Jason orbit, the cross-track mapping is improved mainly for large mesoscale
334 (18%) but not for short mesoscale (5%) because the SARin slope cannot balance the limited
335 resolution of the Jason orbit. The opposite is observed for CryoSat-2 (in the aggregation
336 bands) owing to its cross-track 100 km cross-track resolution: the improvement is limited for
337 100 km or more and the highest improvement is observed for a 50 km radius.

338 In other words, with the CryoSat orbit, the SARin slope is an asset to improve the cross-track
339 observation of smaller mesoscale features (in the band-shaped aggregation of satellite tracks),
340 something that would not be possible on a Jason orbit.

341 Yet higher wavenumber (K) mesoscale eddies also have a smaller amplitude (the SSH power
342 spectrum decreases as a function of $K^{-11/3}$ in the SQG theory, as per Le Traon et al., 2008).
343 Thus changing the correlation radius also induces a reduction of the SSH STD and a reduction
344 of the CTS STD from 2 μ rad to 1.5 μ rad (Table 1). In other words, higher precision SARin
345 slopes would be needed in CryoSat-2's sampling bands because the smaller signal of interest
346 also has weaker slopes.

347 To that extent, and considering the error level discussed in section 5.2, the CryoSat-2 orbit is
348 less attractive than a Jason-like resolution would be, because the gain with SARin is
349 geographically limited and because the orbit is more demanding in terms of CTS error budget.

350 **5.2 Slope error**

351 The simulations from section 4.2 showed that the enhancement of cross-track mesoscale
352 mapping was possible in favorable signal to ratio conditions. The expected benefit from actual
353 Cryosat-2 data raises the question of the error level of current datasets. Yet the error spectrum
354 of SARin data in a mesoscale context is not known. Indeed, SARin acquisition zones on
355 ocean are small and/or limited in time. So it is not possible to get datasets that are large
356 enough to observe correlated errors in space or in time. The study from Galin et al (2012) is
357 the first to provide a CTS error estimate as a bias and noise error on ocean through a
358 comparison with a geoid model.

359 **5.2.1 Biases and long wavelength errors**

360 Galin et al (2012) report biases of the order of 8 μ rad on their 1000 km segments. It is not so
361 much a true bias, as a long wavelength correlated error (e.g. orbital revolution) since they also
362 observe a correlation with thermal conditions on the orbit (i.e. linear on 1000 km segments).
363 Yet, in this paper, we are ignoring biases and long wavelength errors because we assume that
364 they can be accounted for by multi-satellite cross-calibration.

365 Indeed, at the intersection of satellite tracks (e.g. CryoSat-2 x CryoSat-2 or CryoSat-2 x
366 Jason-2) crossovers points provide a double measurement where the actual SSH anomaly
367 signal is partially cancelled if the temporal distance between both measurements is short
368 enough. It is thus possible to use this observation to detect and to mitigate spatially and
369 temporally correlated signals.

370 Tai et al (1988) have used this approach to empirically reduce orbit errors on the SSH and
371 Dibarboure et al (2011b) have demonstrated the feasibility of reducing the cross-track slope
372 error for the wide-swath altimetry mission SWOT. So, in theory, the same method could be
373 used to reduce CryoSat-2's SARin slope biases. The method would exploit crossover
374 observations using a combination of the along-track and cross-track slope for SARin / SARin
375 crossovers, and a projection of the along-track slope into the opposite along-track plane for
376 SARin / LRM crossovers.

377 Alternatively, long wavelength errors (500 km or more) can be accounted for in the mapping
378 process itself, with an approach derived from Ducet et al (2000). These techniques are used
379 operationally to remove SSH biases and 1000 km errors before mesoscale mapping
380 (Dibarboure et al, 2011a), including for datasets with limited coverage (e.g. ERS-2 after the
381 loss of its on board recorders). The same method could be used in the geographically-limited
382 SARin acquisition zones to cross-calibrate long-wavelength errors in the cross-track slope.

383 **5.2.2 Noise and short scale errors**

384 In the recovery of the cross-track slope, Galin et al (2012) also observe on average 20 μ rad of
385 speckle-related noise at 1 Hz or 7 km resolution. The slope is computed from a distance
386 ranging from 1 km to 8 km depending on the retrieval algorithm (phase-difference at the first
387 point of arrival VS. model fit) and significant wave height (SWH) conditions.

388 The spatial correlation of mesoscale slope makes along-track filtering possible (including
389 with non linear filters to remove spurious slopes) if the error is speckle-related (i.e. no along-
390 track correlation of the CTS error). If we assume that a simple running average is used to get
391 one super-observation for a 150 km radius (admittedly a crude filtering), the resulting
392 mesoscale slope precision would be less than 4 μ rad with current slope retrieval algorithms.

393 Moreover, Galin et al investigate the origin of residual SARin slope outliers such as the
394 influence of wind and so-called sigma0 blooms. Yet sigma0 blooms can be detected and
395 edited out in pulse-limited LRM altimetry (Thibaut et al, 2010). We can therefore assume that
396 the largest SARin slope outliers can be detected as well, thus decreasing the overall slope
397 error RMS of a non-Gaussian slope error distribution.

398 With Cryosat-2 we can probably observe only large eddies ($2\text{-}\sigma$) in zones of intense
399 mesoscale variability. Elsewhere, SARin slopes from Cryosat-2 can probably barely improve
400 cross-track mesoscale mapping because the instrument was not designed for this application
401 (insufficient signal to error ratio).

402 **5.2.3 MSS and geoid errors**

403 In this section, we discuss MSS model errors and their influence on SARin slope anomalies in
404 the context of mesoscale mapping. Indeed, mesoscale mapping is based on sea level
405 anomalies (SLA), not sea surface heights (Dibarboure et al, 2012) and the SLA is created as

406 the difference of the measured SSH and a temporal reference or $\langle \text{SSH} \rangle$. The orbit used by
407 CryoSat-2 is geodetic (one-year repeat cycle, described in Francis et al, 2007) so gridded
408 MSS or geoid models are used as a $\langle \text{SSH} \rangle$ reference. The same stands for CTS anomalies
409 which are the difference of the CTS measurement and the cross-track gradient of the MSS
410 model. Consequently any error in the models generates a CTS anomaly error (i.e. an
411 additional CTS error in Figure 10).

412 Pavlis et al (2008) show that in favourable conditions along the well-known TOPEX/Poseidon
413 tracks, they observe an error of $2 \mu\text{rad}$ at 1 Hz for EGM08. In a different context, Sandwell
414 and Smith (2009) have shown through comparisons with shipboard gravity that the accuracy
415 of altimetry-derived gridded gradients was of the order of a few μrad in zones of rugged
416 seafloor topography. More recently, Schaeffer et al (2012) have shown that the gradient error
417 of their MSS model (CNES/CLS2011) ranged from $1 \mu\text{rad}$ along charted tracks of repetitive
418 altimetry mission to $5 \mu\text{rad}$ in areas covered only by geodetic altimetry missions. Moreover,
419 Andersen & Rio (2011) and Dibarboure et al (2012) highlighted differences between
420 independent MSS models that range from 1 to 3 cm with wavelengths ranging from 3 to
421 hundreds of kilometers (a few μrad after along-track smoothing).

422 The MSS/geoid error is therefore quite significant in the error budget of a SARin CTS
423 anomaly, since it would add up to noise estimates from section 5.2.2. That error alone would
424 make error-free CTS measurements difficult to use except in zones of strong mesoscale
425 activity.

426 **5.2.4 Expected and possible improvements**

427 Comparing the figures of merit from section 5.2 to the sensitivity studies from section 4.2
428 shows that the precision needed to improve cross-track mesoscale mapping in strong currents
429 is at the limit of CryoSat-2's current observation capability.

430 However, one might expect some improvements in the future. The primary error sources were
431 shown to be speckle-related measurement noise and the MSS reference models used to
432 generate the slope anomaly.

433 Concerning the former, it might be technically possible to update onboard software to get
434 SAR data from both receive chains on ocean, and to change acquisition rates in SARin mode,
435 essentially yielding 4 times as many independent looks, and reducing the noise level.
436 Moreover, the SAR and SARin retrieval algorithms are relatively young, especially in an
437 oceanography context (CryoSat-2 is an ice mission), and Galin et al give some interesting
438 outlook that might result in a better precision: filtering, and weighting of beams...

439 And concerning the latter, our error estimate are derived from *2008-2001* generation MSS
440 models, which are not yet integrating new geodetic data from CryoSat-2, Jason-1 GM
441 (geodetic phase), let alone from new and upcoming missions flying on uncharted tracks (e.g.
442 Sentinel-3A and 3B, HY-2). It is likely that the current and future altimeter datasets will
443 decrease the error level of the future reference models, and especially at short wavelengths.

444 Beyond CryoSat-2, our findings raise the question of a prospective SARin demonstrator
445 optimized for oceanography (with synergies with other applications). In this context, the
446 outlook is even more promising because additional changes could be considered: on the orbit,
447 on the hardware, and reference surface models.

448 CryoSat-2's orbit was shown to be suboptimal for SARin usage in section 5.1 and a dedicated
449 mission could use a different orbit such as the ones analyzed by Dibarboure et al (2012) for
450 the geodetic phase of Jason-1.

451 Moreover, if a new instrument derived from SIRAL were used on a dedicated SARin
452 demonstrator, various upgrades could be considered to increase the number of statistically
453 independent looks and to decrease the speckle-related noise: antenna design and beam width,

454 baseline length, pulse timing (e.g. continuous or interleaved mode VS SIRAL's burst
455 mode)... However it is possible that the global mesoscale requirement from section 4.2
456 (precision of the order of 0.1 to 0.5 μ rad) might remain challenging.

457 Lastly, in the context of *global* SARin acquisition with a sufficient precision, such a
458 prospective mission would acquire east/west gradients which would help resolve the shortest
459 wavelengths in MSS, geoid or bathymetry models since they are difficult to resolve with the
460 current anisotropy of altimeter data (Sandwell et al., 2011). In turn, this would further
461 mitigate the errors from the \langle SSH \rangle reference discussed in section 5.2.3.

462 **5.3 Validity and limitations of this work**

463 In this section we discuss some approximations made in this paper, and the validity and
464 limitations of these factors as an outlook for future work: the Gaussian properties of our
465 “reality”, the perfect a priori knowledge used in the mapping process, the simple 1D mapping
466 methodology used, and the lack of temporal variability.

- 467 • In section 3.2, our reality is a random Gaussian process with a decorrelation function
468 consistent with scales reported by Le Traon et al (2003). In practice, our reality has a
469 flat power spectrum density for long wavelengths and a cut-off for shorter
470 wavelengths. In other words, we do not use the covariance model from operational
471 mesoscale mapping (e.g. Ducet et al, 2010), but our covariance model and the
472 associated variance-preserved power spectra are representative of a diversity of
473 wavelengths, much like along-track filtered altimeter measurements.
- 474 • In the OI, we use a priori knowledge of the covariance of the signal (H_{real}) and the
475 covariance of the error (ϵ) in matrices A and C from Equation 3. In this paper, we use
476 the true analytical covariance model used to simulate our dataset (i.e. the covariance

477 model of our Gaussian reality), resulting in a non-existent mapping error for error-free
478 measurements. However, in practice, we only have an approximate knowledge of the
479 true ocean decorrelation model (e.g. Jacobs et al, 2001 or Le Traon et al 2003) and of
480 the altimetry error, so the OI process is not perfect. The same stands for the CTS
481 parameter, and the a priori knowledge of the SARin data error. This can be a
482 significant implementation problem so our findings should be revisited with real data.
483 More importantly, this point highlights that one must acquire a better understanding of
484 the SARin error spectrum before such data can be used in an OI.

485 • Lastly one should note that the frozen field assumption and the 1D analysis (cross-
486 track direction) represent a best case configuration for SARin. In reality, mesoscale
487 signals temporally decorrelate over ± 15 to 20 days. Thus our results are optimistic
488 because they do not take into account the high frequency dynamics that Le Traon and
489 Dibarboue (2002) showed to be difficult to resolve with constellations with less than
490 4 altimeter missions. There is also a large panel of complex geometric configurations
491 that vary with latitude. Consequently, because 1D results are encouraging, the findings
492 of this paper should be extended to much more sophisticated 3D simulations (OI or
493 ocean model assimilation), taking into account orbit sampling dynamics
494 (measurements are not ubiquitous, nor regularly spaced out) and the temporal
495 variability of the ocean (reality is not frozen).

496 **6 Conclusion**

497 CryoSat-2's SAR interferometry (SARin) mode has the unprecedented capability to measure
498 the sea surface height slope in the cross-track direction. It is possible to use this parameter to
499 constrain mesoscale mapping, and to improve the resolution in the cross-track direction where
500 the traditional (LRM) radar altimetry is limited by the number of satellites in operations.

501 Idealized mapping simulations show that a single error-free SARin sensor on a Jason-like
502 orbit has the potential to perform almost like two coordinated LRM instruments. Sensitivity
503 studies show that the breakthrough in mapping improvement is achieved for slope errors
504 between 1 and 5 μrad for 150 km macro-observations, in zones of intense mesoscale activity.
505 A better slope precision of the order of 0.1 μrad would be needed for global usage and/or to
506 resolve smaller features (radius < 100 km).

507 The precision needed to improve cross-track mesoscale mapping is probably at the limit of
508 current SARin products from CryoSat-2 (and only after multi-satellite cross-calibration and
509 along-track filtering) which might observe only the strongest slopes ($2\text{-}\sigma$) in very energetic
510 areas. The proof of concept is more attractive if we extrapolate to future improvements of
511 SARin processors and ancillary datasets (e.g. MSS) and to a prospective mission improving
512 upon SIRAL hardware and CryoSat-2 processors.

513 While encouraging, these results are optimistic, because all simulations were performed on a
514 frozen SSH field (ocean dynamics and high frequencies are not taken into account), and only
515 in the cross-track direction (i.e. optimal for the SARin slope) and they should be extended to
516 much more complex 3D studies, or with real data from CryoSat-2.

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521 **8 Appendix: CTS methodology**

522 To use the cross-track slope in the OI process, we use the following covariance models to
 523 describe the relationship between the topography H and the slope S:

$$cov_{hh}(\tau) = K \cdot e^{-\left(\frac{\tau}{d}\right)^2} \quad \text{Equation 5}$$

$$cov_{hs}(\tau) = -\frac{2\tau}{d^2} \cdot K \cdot e^{-\left(\frac{\tau}{d}\right)^2} \quad \text{Equation 6}$$

$$cov_{ss}(\tau) = K \cdot \left(\frac{4\tau^2}{d^4} - \frac{2\tau}{d^2} \right) \cdot e^{-\left(\frac{\tau}{d}\right)^2} \quad \text{Equation 7}$$

524 To inverse the problem, we replace A and C and H_{obs} from Equation 3 (and Equation 4), by
 525 A', C' and H' _{obs}, where H' _{obs} is the new observation vector including topography and slope
 526 measurements as the sum of the true signal H_{real} or S_{real} and a random error ε_H and ε_S
 527 estimated on the across-track position vector x(i):

$$H'_{obs} = \begin{bmatrix} H_{real}(x(1)) + \varepsilon_H(x(1)) \\ \vdots \\ H_{real}(x(N)) + \varepsilon_H(x(N)) \\ \vdots \\ S_{real}(x(1)) + \varepsilon_S(x(1)) \\ \vdots \\ S_{real}(x(N)) + \varepsilon_S(x(N)) \end{bmatrix} \quad \text{Equation 8}$$

528 Matrix C' is the new covariance matrix taking into account both topography and slope
 529 estimates

$$C' = \begin{bmatrix} C_{hh} & C_{hs} \\ C_{hs} & C_{ss} \end{bmatrix} \quad \text{Equation 9}$$

530 where C_{hh} , C_{hs} , and C_{ss} are the three covariance matrixes for each couple of observation type,
 531 built as a function of the distance $d_{ij} = |x(i) - x(j)|$ separating measurements points #i and #j.

$$C_{hh} = \begin{bmatrix} cov_{hh}(d_{1,1}) & \dots & cov_{hh}(d_{1,N}) \\ \vdots & cov_{hh}(d_{i,j}) & \vdots \\ cov_{hh}(d_{N,1}) & \dots & cov_{hh}(d_{N,N}) \end{bmatrix} \quad \text{Equation 10}$$

$$C_{hs} = \begin{bmatrix} cov_{hs}(d_{1,1}) & \dots & cov_{hs}(d_{1,N}) \\ \vdots & cov_{hs}(d_{i,j}) & \vdots \\ cov_{hs}(d_{N,1}) & \dots & cov_{hs}(d_{N,N}) \end{bmatrix} \quad \text{Equation 11}$$

$$C_{ss} = \begin{bmatrix} cov_{ss}(d_{1,1}) & \dots & cov_{ss}(d_{1,N}) \\ \vdots & cov_{ss}(d_{i,j}) & \vdots \\ cov_{ss}(d_{N,1}) & \dots & cov_{ss}(d_{N,N}) \end{bmatrix} \quad \text{Equation 12}$$

532 When the inversion is optimal we also account for the uncorrelated error ε_H and ε_S in the
 533 diagonal of C_{hh} , C_{hs} , and C_{ss} (not shown).

534 Matrix A' describing the covariance between the topography we want to reconstruct H_{est} and
 535 the new observation vector H'_{obs} is created with the method used for C' , but using the distance
 536 d'_{ij} between the position $x(i)$ of our observation points and the position $x'(i)$ our unknown
 537 grid points.

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674 **10 Figure captions**

675 Figure 1: 10-day sampling from an altimeter on the TOPEX/Jason orbit. The white segment
676 highlights the along-track direction with one measurement every 7 km, and the black segment
677 highlights the worst case configuration in the cross-track direction with one measurement
678 every $315 \text{ km} \times \cos(\text{latitude})$.

679 Figure 2: Difference between LRM (a) and SARin (b) measurements in the optimal
680 interpolation for common profiles of cross-track slope (c) and SSH (d). The SARin
681 measurement allows to observe the cross-track slope in addition to the SSH profile given by
682 the LRM mode.

683 Figure 3: Simulated Gaussian field (plain line), observation by a LRM altimeter in the along-
684 track direction (dots, 30 km resolution), and reconstruction at each time step with an optimal
685 interpolation (dashed line) with formal error estimates (grey bars). The upper figure shows the
686 SSH (in cm), and the bottom panel the SSH slope (in μrad).

687 Figure 4: Simulated Gaussian field (plain line), observation by a LRM altimeter in the cross-
688 track direction (dots, 300 km resolution), and reconstruction at each time step with an optimal
689 interpolation (dashed line) with formal error estimates (grey bars). The upper figure shows the
690 SSH (in cm), and the bottom panel the SSH slope (in μrad).

691 Figure 5: Simulated Gaussian field (plain line), observation by two LRM altimeters in the
692 cross-track direction (dots, 150 km resolution), and reconstruction at each time step with an
693 optimal interpolation (dashed line) with formal error estimates (grey bars). The upper figure
694 shows the SSH (in cm), and the bottom panel the SSH slope (in μrad).

695

696 Figure 6: Simulated Gaussian field (plain line), observation by one SARin altimeter in the
697 cross-track direction (dots couples, 300 km resolution), and reconstruction at each time step
698 with an optimal interpolation (dashed line) with formal error estimates (grey bars). The upper
699 figure shows the SSH (in cm), and the bottom panel the SSH slope (in μrad).

700 Figure 7: Simulated Gaussian field (plain line), observation by two SARin altimeters in the
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705 a function of the cross-track slope observation error (standard deviation in μrad). The black

706 dotted lines show the reconstruction error for one LRM altimeter and the black dashed line
707 the reconstruction error for two LRM altimeters. The grey dashed line highlights the curve's
708 point of inflection.

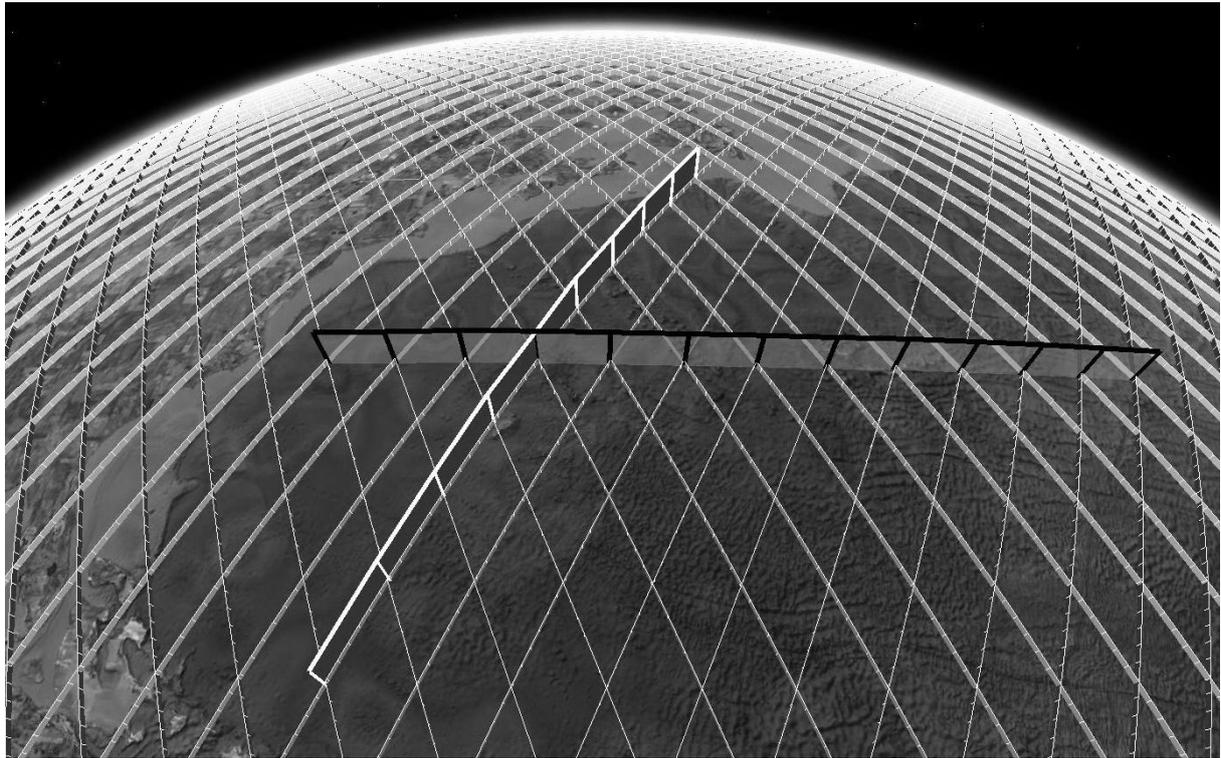
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710 as a function of the cross-track slope observation error (in μrad). The black dashed lines show
711 the reconstruction error for two LRM altimeters and the black line the reconstruction error for
712 four LRM altimeters. The grey dashed line highlights the curve's point of inflection.

713 Figure 10: Cross-track reconstruction error (in % of signal variance) for one SARin altimeter
714 as a function of the cross-track slope observation error (in μrad) and for 3 levels of SSH
715 variability.

716 Figure 11: CryoSat-2's sampling for 15 consecutive days. Satellite tracks (white lines) are
717 aggregated in 500 km wide bands thanks to the 3 and 30 day sub-cycles, and interleaved with
718 500 km bands with few/no satellite tracks.

719 Figure 12: Simulated Gaussian field (plain line), observation in the cross-track direction
720 (dots) by one SARin altimeter on the CryoSat-2 orbit (100 km resolution, packet-aggregated
721 tracks), and reconstruction at each time step with an optimal interpolation (dashed line).
722 Difference between SARin observation (top) and LRM observation (bottom) to constrain 1D
723 OI reconstruction in the 500 km wide blind spot (grey rectangles).

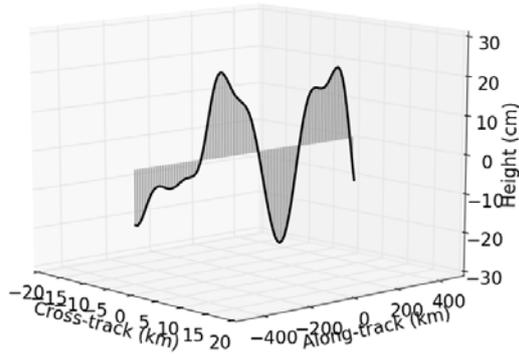
724 Table 1: Reduction of the mapping error from LRM to SARin as a function of the simulation
725 correlation radius. Unit: % of the signal variance. Right-hand side columns show the
726 decreasing amplitude and slope of the eddy as a function of its radius (approximation of SQG
727 theory).



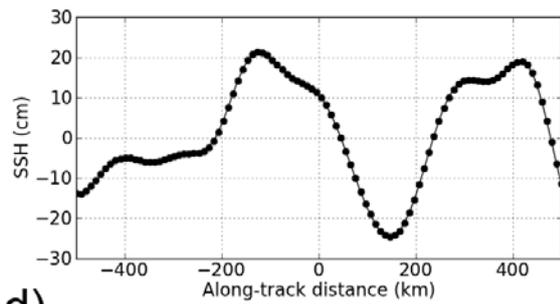
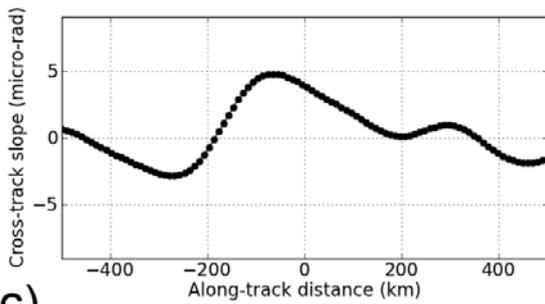
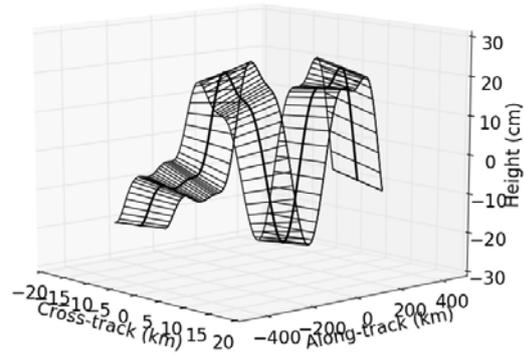
729

730 **Figure 1 : 10-day sampling from an altimeter on the TOPEX/Jason orbit. The white segment**
731 **highlights the along-track direction with one measurement every 7 km, and the black**
732 **segment highlights the worst case configuration in the cross-track direction with one**
733 **measurement every $315 \text{ km} \times \cos(\text{latitude})$.**

a) LRM sea surface height profile



b) SARin height & slope profile



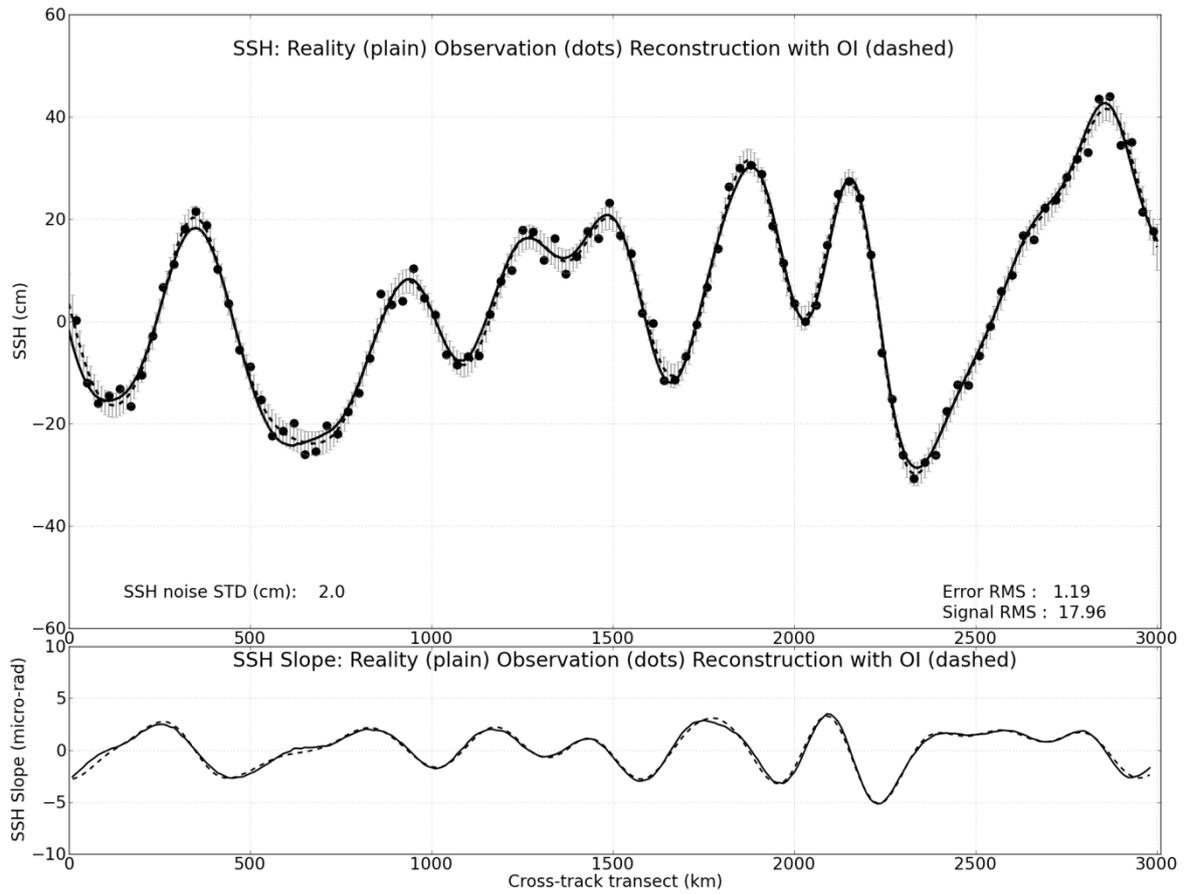
c)

d)

734

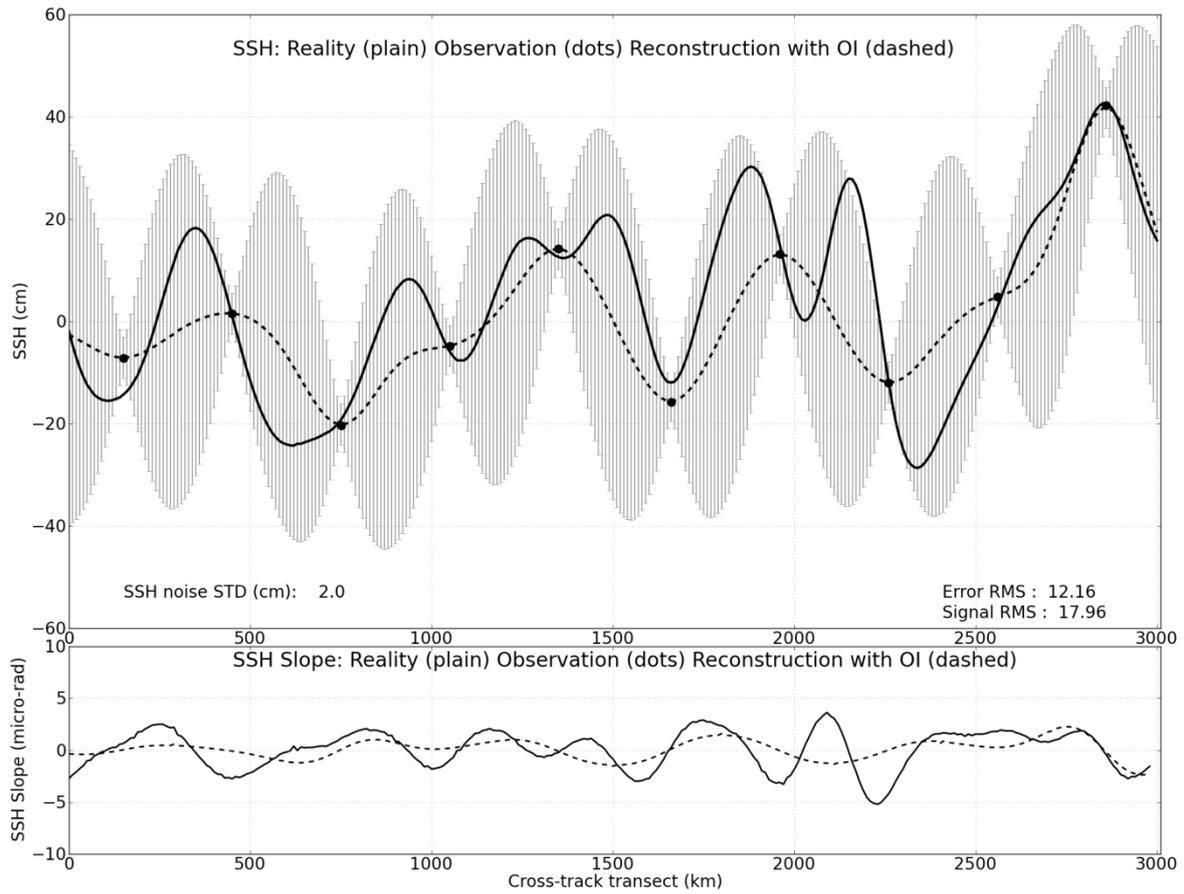
735 **Figure 2 : Difference between LRM (a) and SARin (b) measurements in the optimal**
 736 **interpolation for common profiles of cross-track slope (c) and SSH (d). The SARin**
 737 **measurement allows to observe the cross-track slope in addition to the SSH profile given**
 738 **by the LRM mode.**

739



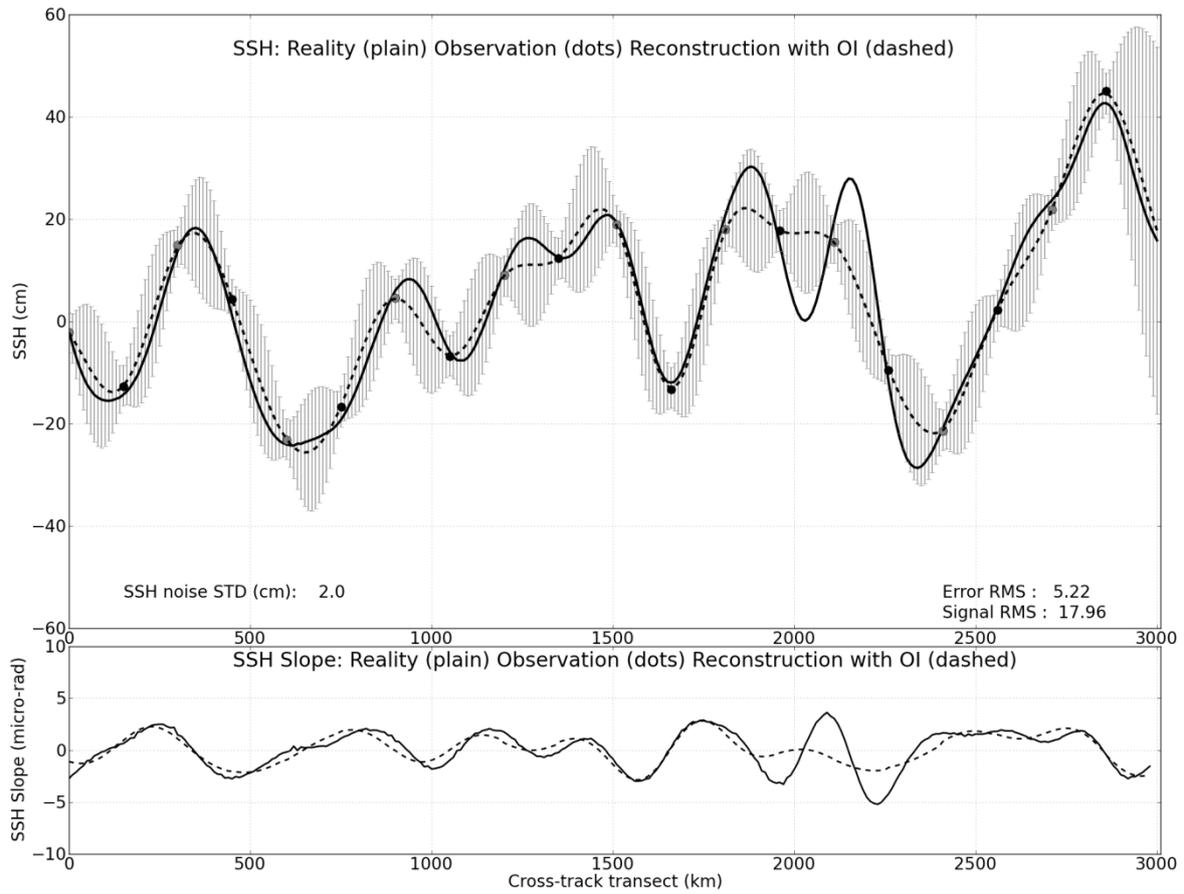
740

741 **Figure 3 : Simulated Gaussian field (plain line), observation by a LRM altimeter in the**
 742 **along-track direction (dots, 30 km resolution), and reconstruction at each time step with**
 743 **an optimal interpolation (dashed line) with formal error estimates (grey bars). The upper**
 744 **figure shows the SSH (in cm), and the bottom panel the SSH slope (in μrad).**



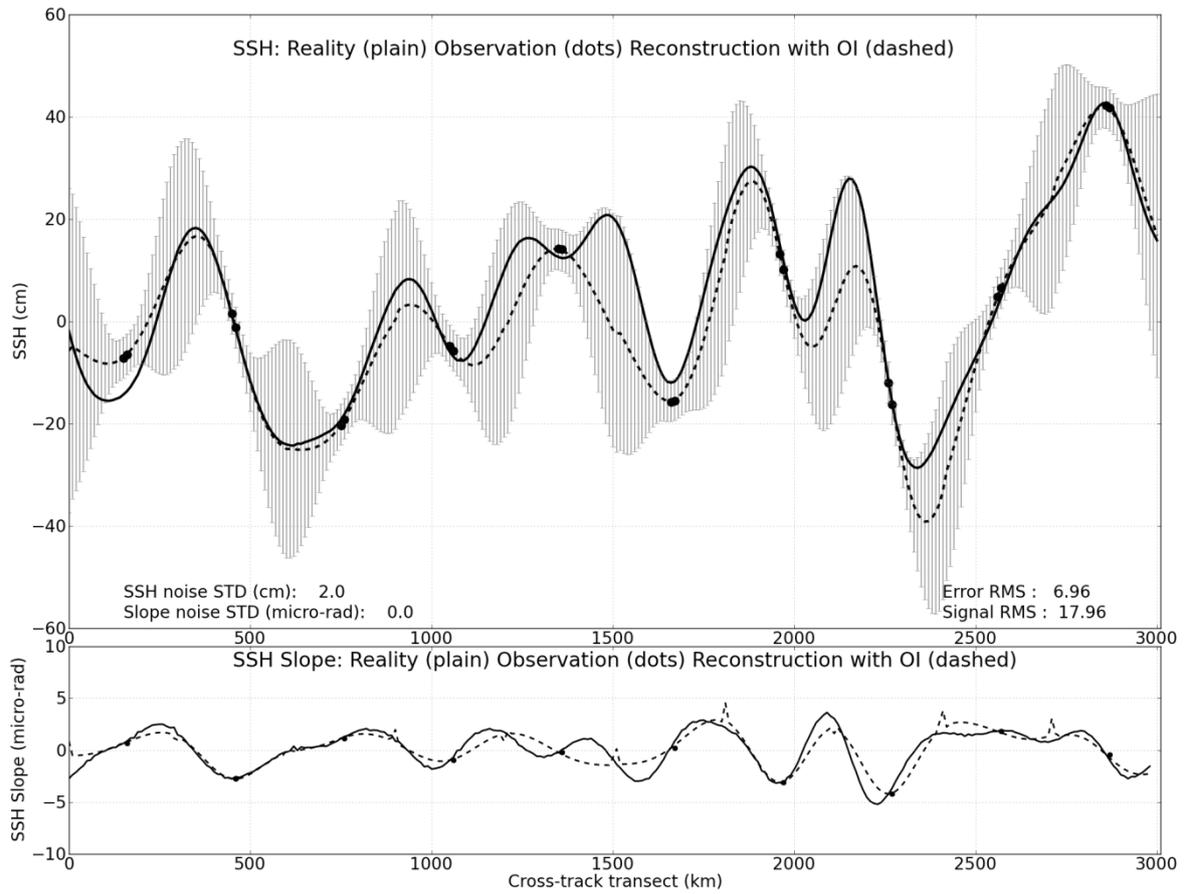
745

746 **Figure 4 : Simulated Gaussian field (plain line), observation by a LRM altimeter in the cross-**
 747 **track direction (dots, 300 km resolution), and reconstruction at each time step with an**
 748 **optimal interpolation (dashed line) with formal error estimates (grey bars). The upper**
 749 **figure shows the SSH (in cm), and the bottom panel the SSH slope (in μ rad).**



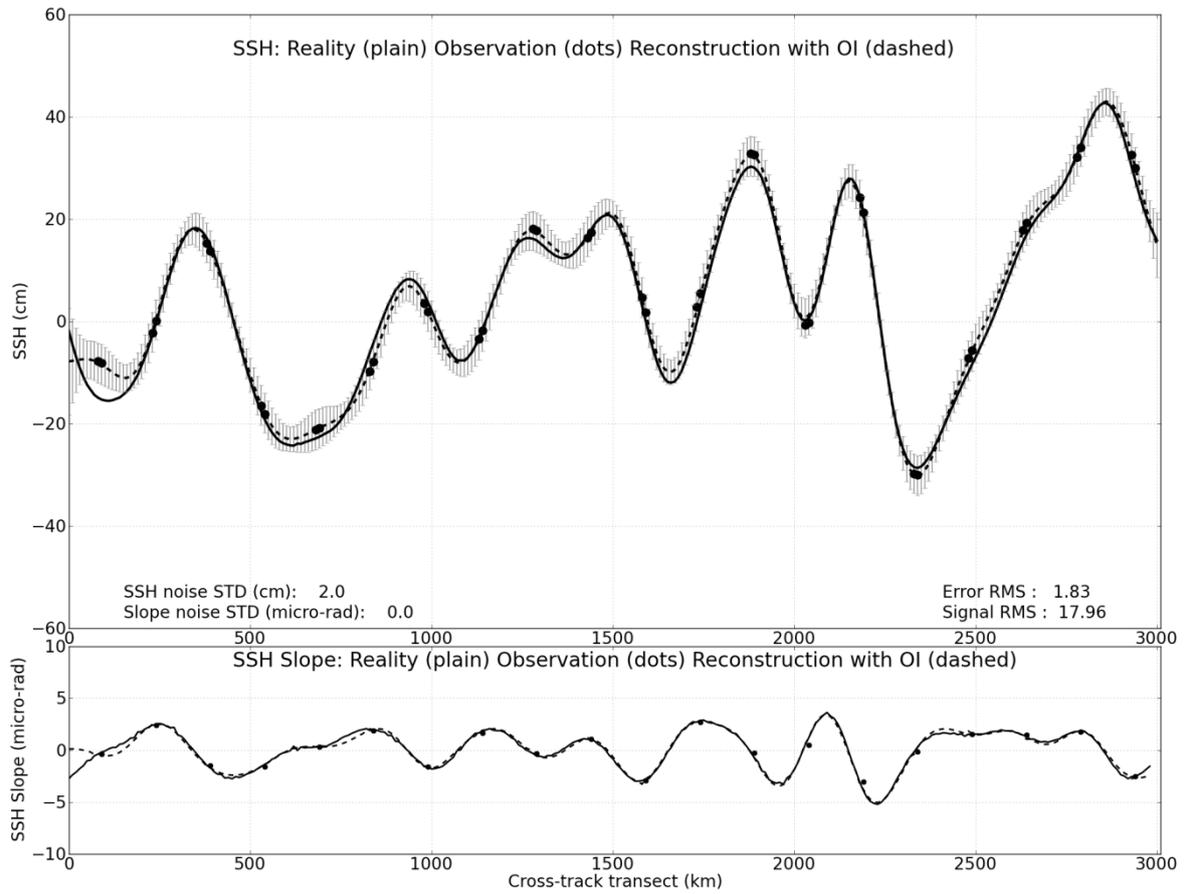
750

751 **Figure 5 : Simulated Gaussian field (plain line), observation by two LRM altimeters in the**
 752 **cross-track direction (dots, 150 km resolution), and reconstruction at each time step with**
 753 **an optimal interpolation (dashed line) with formal error estimates (grey bars). The upper**
 754 **figure shows the SSH (in cm), and the bottom panel the SSH slope (in μrad).**



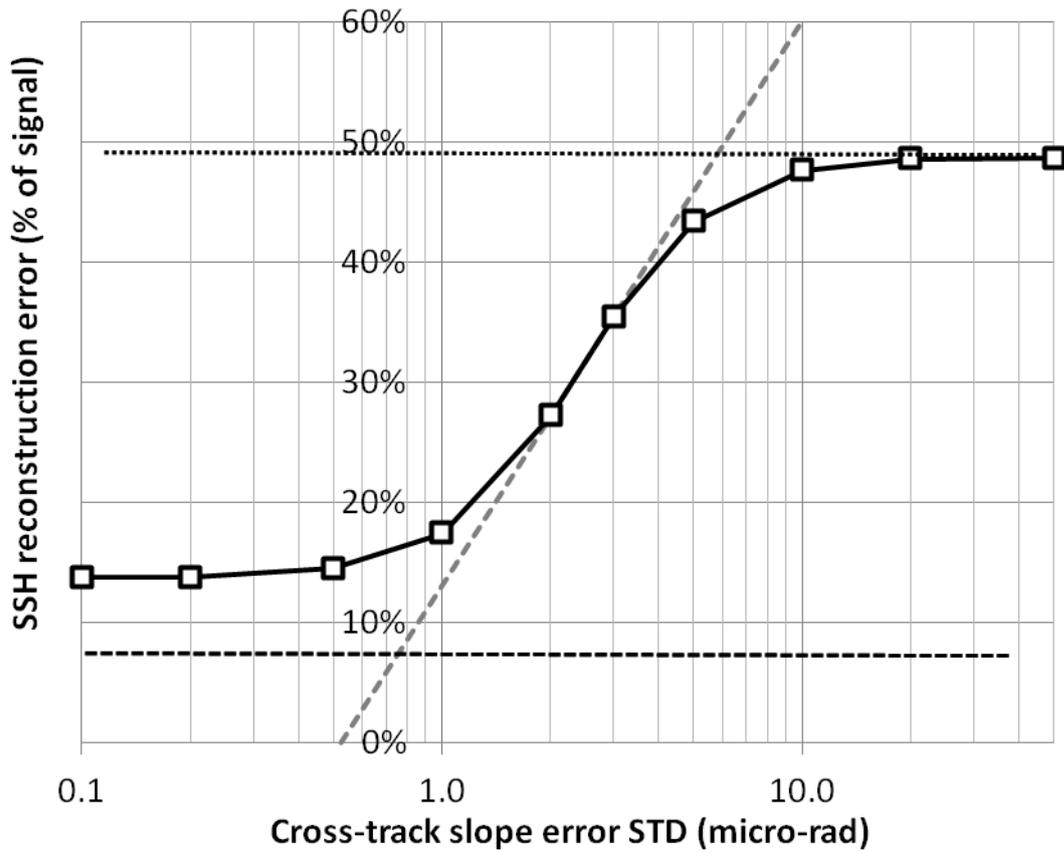
755

756 **Figure 6 : Simulated Gaussian field (plain line), observation by one SARin altimeter in the**
 757 **cross-track direction (dots couples, 300 km resolution), and reconstruction at each time**
 758 **step with an optimal interpolation (dashed line) with formal error estimates (grey bars).**
 759 **The upper figure shows the SSH (in cm), and the bottom panel the SSH slope (in μrad).**



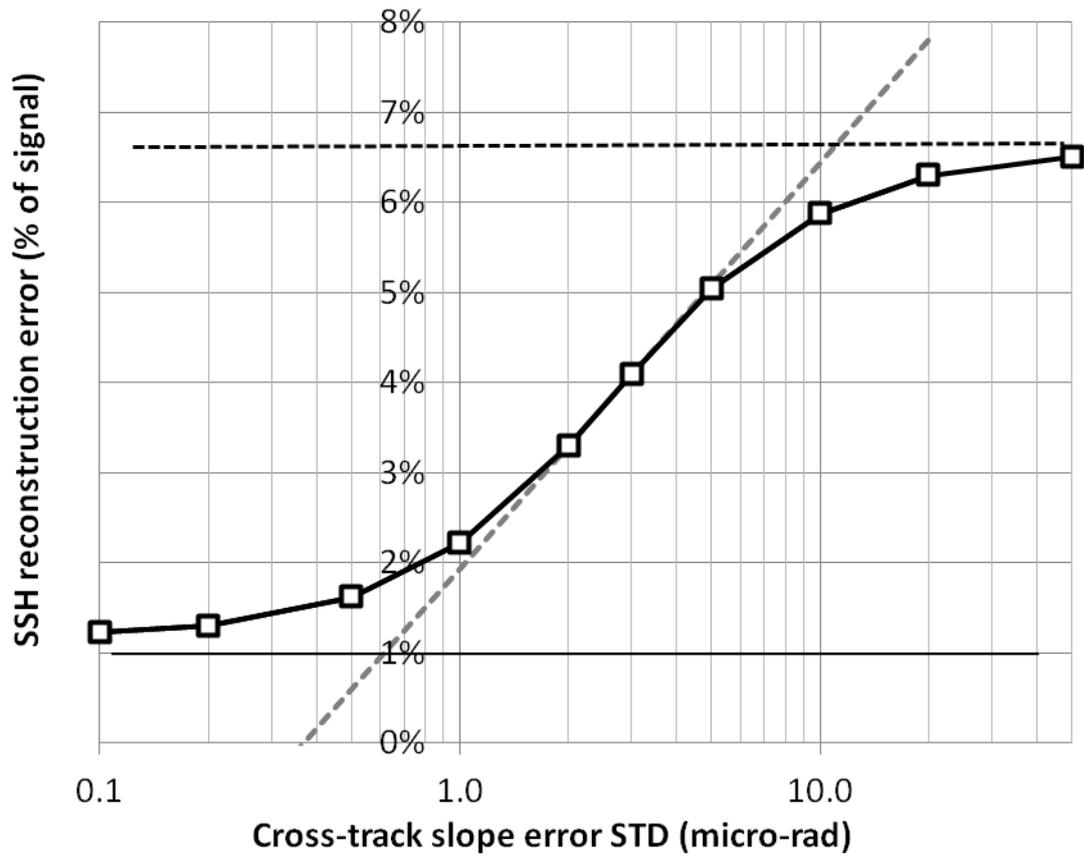
760

761 **Figure 7 : Simulated Gaussian field (plain line), observation by two SARin altimeters in the**
 762 **cross-track direction (dots couples, 300 km resolution), and reconstruction at each time**
 763 **step with an optimal interpolation (dashed line) with formal error estimates (grey bars).**
 764 **The upper figure shows the SSH (in cm), and the bottom panel the SSH slope (in μ rad).**



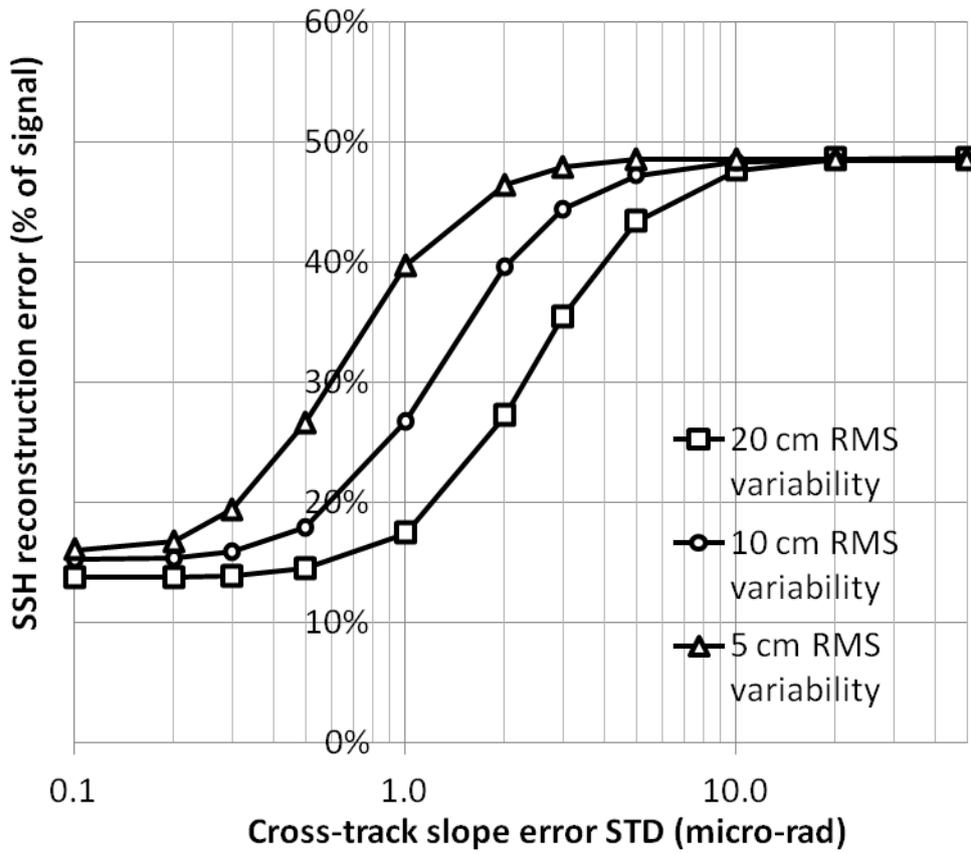
765

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 767 **as a function of the cross-track slope observation error (standard deviation in μrad). The**
 768 **black dotted lines show the reconstruction error for one LRM altimeter and the black**
 769 **dashed line the reconstruction error for two LRM altimeters. The grey dashed line**
 770 **highlights the curve's point of inflection.**



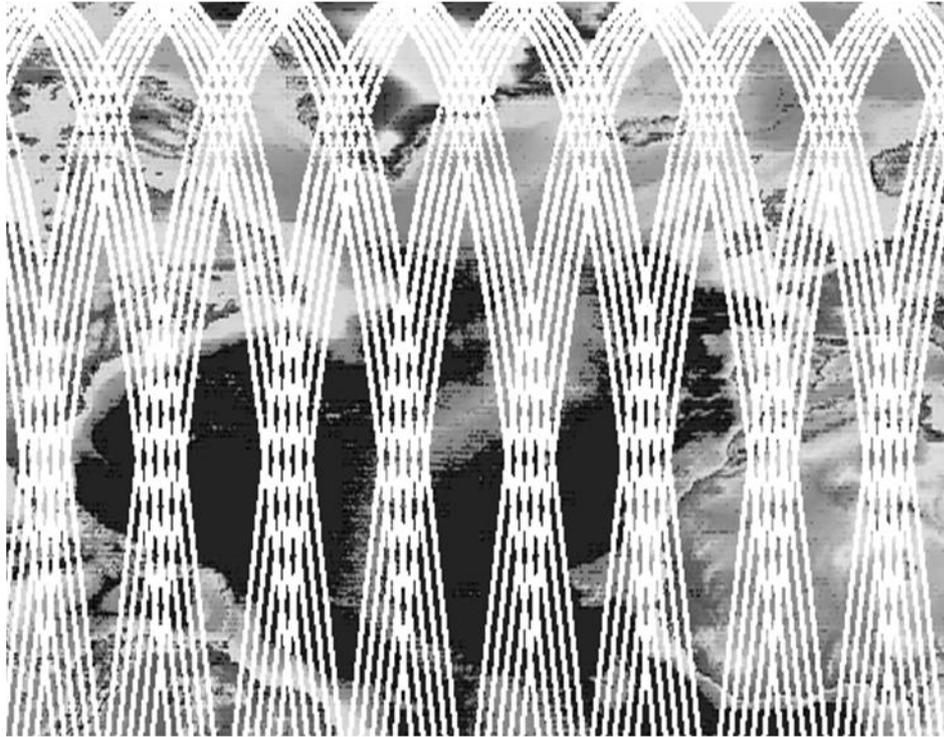
771

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 774 **show the reconstruction error for two LRM altimeters and the black line the**
 775 **reconstruction error for four LRM altimeters. The grey dashed line highlights the curve's**
 776 **point of inflection.**



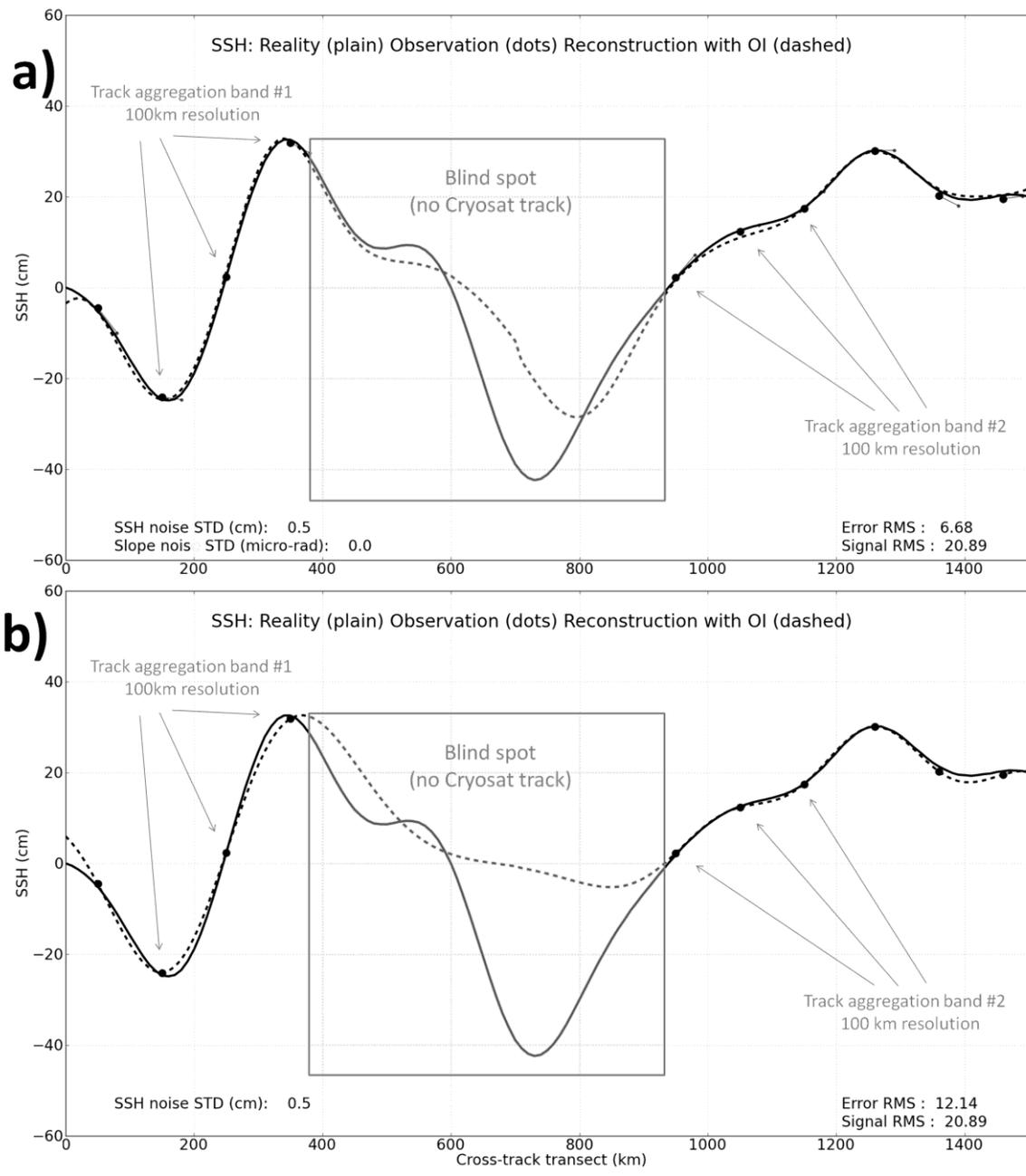
777

778 **Figure 10 : Cross-track reconstruction error (in % of signal variance) for one SARin**
 779 **altimeter as a function of the cross-track slope observation error (in μ rad) and for 3 levels**
 780 **of SSH variability.**



781

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783 **aggregated in 500 km wide bands thanks to the 3 and 30 day sub-cycles, and interleaved**
784 **with 500 km bands with few/no satellite tracks.**



785

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 788 **aggregated tracks), and reconstruction at each time step with an optimal interpolation**
 789 **(dashed line). Difference between SARin observation (top) and LRM observation (bottom)**
 790 **to constrain 1D OI reconstruction in the 500 km wide blind spot (grey rectangles).**

Radius	Mapping Improvement		Reality signal properties	
	Cryosat-2 Orbit	TP/Jason Orbit	SSH STD (cm)	SSH Slope STD (μ rad)
150	0%	18%	20.0	1.9
125	1%	9%	15.8	1.8
100	2%	7%	11.8	1.7
75	8%	5%	8.1	1.5
50	14%	5%	4.8	1.4

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Table 1: Reduction of the mapping error from LRM to SARin as a function of the simulation correlation radius. Unit: % of the signal variance. Right-hand side columns show the decreasing amplitude and slope of the eddy as a function of its radius (approximation of SQG theory).