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## Tidal downscaling from the open ocean to the coast: a new approach applied to the Bay of Biscay

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### Abstract :

Downscaling physical processes from a large scale to a regional scale 3D model is a recurrent issue in coastal processes studies. The choice of boundary conditions will often greatly influence the solution within the 3D circulation model. In some regions, tides play a key role in coastal dynamics and must be accurately represented.

The Bay of Biscay is one of these regions, with highly energetic tides influencing coastal circulation and river plume dynamics. In this study, three strategies are tested to force with barotropic tides a 3D circulation model with a variable horizontal resolution. The tidal forcings, as well as the tidal elevations and currents resulting from the 3D simulations, are compared to tidal harmonics extracted from satellite altimetry and tidal gauges, and tidal currents harmonics obtained from ADCP data.

The results show a strong improvement of the M2 solution within the 3D model with a "tailored" tidal forcing generated on the same grid and bathymetry as the 3D configuration, compared to a global tidal atlas forcing. Tidal harmonics obtained from satellite altimetry data are particularly valuable to assess the performance of each simulation. Comparisons between sea surface height time series, a sea surface salinity database, and daily averaged 2D currents also show a better agreement with this tailored forcing.

### Highlights

► A new approach is proposed to tidal downscaling in 3D coastal models. ► The best tidal solution is not necessarily the best forcing for downscaling. ► The error on the M2 elevation in the Bay of Biscay is reduced by 75%. ► The success of the method is also assessed through its impact on circulation. ► Satellite altimetry is a powerful tool for tidal assessment.

**Keywords :** Tides, Downscaling, 3D coastal modelling, Boundary conditions, Satellite altimetry, Bay of Biscay



## 1 **1. Introduction**

2       Increasing efforts are made to improve the accuracy of global circulation models at regional scales, by improving  
3 the grid resolution, by taking into account more physical processes or through data assimilation techniques (e.g. Holt  
4 et al. (2017)). In spite of significant progresses in the recent years, the global or basin simulations performance  
5 generally remains insufficient to accurately study coastal phenomena, and regional models are still the best option,  
6 thanks to their higher resolution, tuned parameterizations or parameters, and to the consideration of comprehensive  
7 coastal processes such as tides, surface waves, estuarine processes, etc. Since regional physical processes are partly  
8 driven by large scale processes (Zheng and Weisberg, 2012), with this limited-area approach comes the issue of

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9 downscaling and managing open boundary conditions. As first stated by Oliger and Sundström (1978), open boundary  
10 conditions can never be considered as perfect. Several strategies have been developed to deal with this issue, as  
11 discussed for instance by Blayo and Debreu (2005) and Herzfeld (2009). More particularly, the nesting of several  
12 grids within each other are often used to gradually increase the resolution near the coast. However, the interpolation  
13 necessary due to resolution differences and bathymetry inconsistencies may induce errors at the open boundaries.

14 Modelling the 3D ocean circulation in coastal areas and shelf seas requires an accurate representation of the tidal  
15 dynamics, especially near the coast. The tidal solution in a regional circulation model results from the introduction  
16 of the astronomical tidal potential in the primitive equations, and from open boundary conditions in sea surface  
17 elevation (hereafter SSH) and currents. The accuracy of the tidal forcing at the open boundaries is critical for the  
18 representation of tides of course, but also for the simulation of mixing and circulation through different mechanisms:  
19 non-linear interactions between tidal currents and the general circulation, mixing induced by internal tides, bottom  
20 friction modulation by tidal currents, mixing enhancing by vertical tidal currents shear (Carter and Merrifield, 2007;  
21 Herzfeld, 2009; Guarnieri et al., 2013). Guarnieri et al. (2013) show the impact of tides on the Adriatic Sea circulation,  
22 with a 3D model. They find that tides influence the circulation by modifying the horizontal advection, especially  
23 during periods of weak wind stress. They also assess the impact of tides on mixing, this time for strong wind stress  
24 periods. Residual tidal flows due to non-linear interactions with the topography ('topographic rectification') can also  
25 be generated (González-Pola et al., 2012; Wang et al., 2013). Holt et al. (2017) show that the inclusion of tides in  
26 circulation models allows a better representation of seasonal stratification cycles than high resolution models without  
27 tides.

28 In principle, the tidal forcing at the open-boundaries is given either as a set of tidal constituents or as time-varying  
29 fields of sea surface elevation and horizontal currents. The latter option is for instance tempting when the coastal  
30 model is also forced at the open-boundaries by a large-scale circulation model that simulates both the tidal and non-  
31 tidal circulation. However, such an option requires the availability of the large-scale forcing at very high-frequency  
32 (a few minutes) which is, in practice, never (or very rarely) possible. That is why in most cases, the open-boundary  
33 conditions for tides and for the non-tidal circulation are prescribed as distinct sources. We have adopted such an  
34 approach in the present study.

35 Downscaling tides in a coastal (child) model is not a trivial issue: as for the general problem of open-boundary  
36 conditions, the difficulties come from the numerical scheme or from the prescribed fields (at last for incoming con-  
37 ditions or 'active boundaries') stemming from the parent model. Another difficulty is introduced if the model is also  
38 forced at the open boundaries (hereafter OB) by low frequency motions. For instance, Herzfeld and Gillibrand (2015)  
39 discuss the problem of dealing with multiple timescales in a scheme based on local adjustment of the flux at the OB;  
40 they propose an approach based on dual relaxation timescales for their scheme. In general, the use of prescribed tidal  
41 fields lead to inconsistencies with the interior solution, mainly due to differences in bathymetry between the forcing  
42 and forced models. As an example, Wang et al. (2013) note that an adjustment of the prescribed tidal barotropic  
43 velocity at the OB is necessary to ensure consistency of the depth integrated barotropic transport with the interior

44 solution. In the case of baroclinic tides, other complications come from possible inconsistencies between the child  
45 and parent stratification as well as from the non-stationary part of the internal wave fields that requires the availability  
46 of the parent outputs at very high-frequency. For these reasons, in cases of offline downscaling problems such as the  
47 one addressed in this paper, only the barotropic tides are taken into account.

48 To prescribe barotropic tides at the OB, two strategies are usually adopted. The most common one is based on the  
49 use of tidal atlases that provide tidal harmonics (amplitude and phase) of sea surface elevation and, in most case, of  
50 barotropic velocities for a given tidal spectrum. Several global atlases exist and are regularly updated (for a review see  
51 Stammer et al. (2014)): some of them are built from empirical adjustment mostly from satellite altimetry to a prior  
52 model, such as the GOT (Ray, 1999), or TPXO (Egbert and Erofeeva, 2002) models. Other atlases are solutions of  
53 barotropic hydrodynamical models constrained by satellite and/or in situ observations via data assimilation. Among  
54 the latter, FES2012 is the last distributed product from a long series of solutions obtained with the T-UGOm hydrody-  
55 namical model (Lyard et al., 2006) described in section 2.2. (At the time when we write this paper, the FES2014 atlas  
56 is under construction). Several examples of regional or coastal circulation models that prescribe tidal harmonics from  
57 global atlases at their open-boundaries are found in the recent literature: Dong et al. (2011) and Wang et al. (2013)  
58 use the TPXO.6 solution for their regional models in the Southern California Bight and Prince William Sound respec-  
59 tively; Katavouta and Thompson (2016) use FES2004 over the Nova Scotia Shelf. In coastal/estuarine applications,  
60 one or several levels of nesting are often necessary and the open-boundary conditions may be obtained from a larger  
61 scale model, as done in Touloukian et al. (2016).

62 Another strategy consists in running the regional or coastal model in a 2D mode without any other forcing than  
63 tidal harmonics in sea surface elevation at the OB (the latter provided by an atlas). The solution of this barotropic  
64 simulation gives tidal constituents that are then used to force the model in 3D mode. The tidal spectrum that can be  
65 estimated from the 2D run depends mainly on the length of the simulation. Such an approach has been used in the  
66 North-East Atlantic by Maraldi et al. (2013).

67 In this paper, we therefore address the issue of downscaling barotropic tides in a circulation model, where an ac-  
68 curate representation of tides is required either for the tidal signal itself (both barotropic tides and internal tides) or for  
69 its impact on the circulation and hydrology. There is a wide literature on open-boundary conditions (hereafter OBC)  
70 in regional models, and many variants of the Dirichlet, Flather, radiation and relaxation conditions are developed,  
71 based on different implementations on the model grid and different strategies regarding sponge layers. A thorough  
72 work with the SYMPHONIE model has been made to implement relevant OBC for coastal applications in presence  
73 of strong or weak tides and consideration to fundamental properties (such as conservation of mass, energy) has been  
74 given. This is summarized in the paper of Marsaleix et al. (2006). We have not found any drawbacks with this scheme.  
75 We do not claim its superiority to alternative schemes either.

76 The two main sources of errors arising with OBC are the errors linked to the equations and numerical implemen-  
77 tation of the OBC method and those due to the possible inconsistency between external forcing and interior dynamics;  
78 in this study we have made the choice to address the latter only. This is a choice motivated by the need to find a

79 relatively easy and fast-to-implement method, that can be applied in different configurations, as an alternative to the  
80 revisit or adjustment of the numerical scheme and equations.

81 Our objective in this paper is to propose a robust and simple approach that allows to improve the downscaling  
82 of barotropic tides for any given set of boundary equations and of external forcing, therefore being non-intrusive in  
83 the model equations. In other words, given a certain 3D circulation model, with a given grid and bathymetry, how  
84 can we improve the tidal forcing to reduce errors on the interior tidal solution? Our new approach is based on the  
85 additional use of a tidal model, here the T-UGOm model of Lyard et al. (2006). Our 3D coastal model is SYMPHONIE  
86 (Marsaleix et al., 2008, 2009). To avoid inconsistencies between the prescribed tides and the interior solution due to  
87 mesh resolution and bathymetry differences, tidal boundary conditions are generated on the same grid and bathymetry  
88 as the ones used by the 3D circulation model. The unstructured 2D spectral model T-UGOm was adapted to perform  
89 simulations on a structured, variable horizontal resolution grid, by introducing C-grid equivalent quadrangle elements.

90 This approach is applied to the Bay of Biscay, where tides are highly energetic, particularly over the western  
91 French shelf with tidal ranges reaching 6 m locally at the coast. Tides are dominated by M2 (Cavanie and Hyacinthe,  
92 1976; Cartwright et al., 1980; Le Cann, 1990), with amplitudes ranging between 1 to 2 m, against a few centimeters  
93 for K1. Non-linear interactions occurring between semi-diurnal constituents and the topography can result in the  
94 generation of overtides such as M4, which can reach amplitudes of 25 cm. Le Cann (1990) showed that the width of  
95 the Bay of Biscay is close to resonance for quarter-diurnal tides, leading to a strong amplification of these constituents.  
96 Figure 1 shows the distribution of the M2 tide (elevation and current) in the Bay of Biscay, taken from the FES2012  
97 tidal atlas. In addition, Table 1 gives the minimum, mean and maximum values for the tidal amplitude of M2, S2, M4  
98 and K1, in the Bay of Biscay.

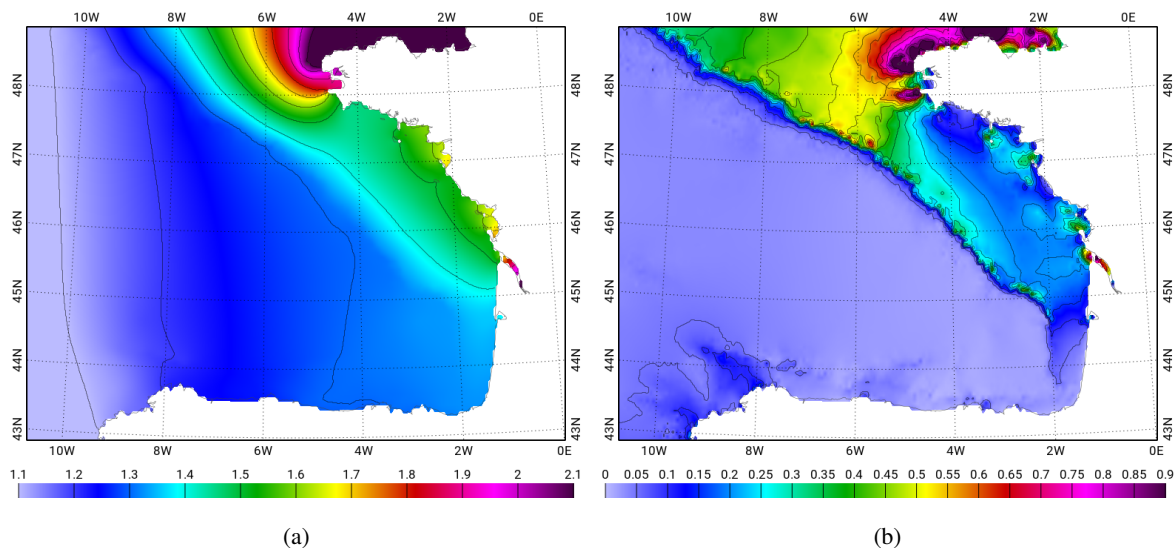


Figure 1: (a) M2 elevation (m), (b) M2 current (m/s), taken from the FES2012 tidal atlas

99 The work of Pairaud et al. (2008, 2010) has shown the ability of the SYMPHONIE model in a regional confi-

Table 1: Minimum, mean and maximum amplitudes (m) of the tidal elevation for M2, S2, M4 and K1, in the Bay of Biscay, from FES2012

	M2	S2	M4	K1
Mean	1.36	0.48	0.035	0.070
Min	1.20	0.42	0.0070	0.063
Max	2.1	0.77	0.25	0.076

100 uration (horizontal resolution of 1.5 km) to reproduce tides in the Bay of Biscay, the major sources of error being  
 101 the bathymetry and boundary conditions. The latter are prescribed from a regional atlas. Since these studies, the  
 102 bathymetry has been significantly improved by merging different datasets (Lyard, pers. comm., 2016). In this study,  
 103 our configuration of SYMPHONIE covers the bay from the deep plain to the shelf and coastal shallow waters.

104 In the first part of this paper, the configuration applied to the Bay of Biscay and the data used to assess the  
 105 solution are presented. The three strategies chosen to constrain the 3D circulation model boundaries with tides are  
 106 then detailed. The 3rd and 4th section are dedicated to the performance evaluation of the different tidal boundary  
 107 conditions. First, the forcing solutions are compared, followed by the 3D circulation simulations, which are evaluated  
 108 in two stages: the tidal solution, with respect to tidal elevations and currents; the "global" simulation, comparing more  
 109 integrating variables (SST, SSS, SSH and total currents). These results are then discussed before concluding.

## 110 2. Model and data

### 111 2.1. The SYMPHONIE model and the BOBSHELF configuration

112 In this section, the SYMPHONIE code used for this study is presented, as well as the BOBSHELF grid and  
 113 configuration, which is an application of SYMPHONIE to the Bay of Biscay.

#### 114 2.1.1. The SYMPHONIE circulation model

115 The SYMPHONIE model is based on the Boussinesq hydrostatic equations of momentum, temperature and salin-  
 116 ity. The primitive equations are discretized and solved on an Arakawa C-grid, using an energy conserving finite  
 117 difference method described in Marsaleix et al. (2008, 2009, 2012). Following Damien et al. (2017), horizontal ad-  
 118 vection and diffusion of momentum are respectively computed with a 4th order centered and a bi-harmonic scheme,  
 119 while vertical advection of momentum is given by a 2nd order centered scheme. Advection and diffusion of tracers  
 120 are computed using the QUICKEST scheme (Neumann et al., 2011). A wet and drying scheme is used with the same  
 121 vertical coordinates; when the water column thickness drops below 1m, the wetting and drying algorithm freezes the  
 122 tracers and cancels out the baroclinic velocities, which means that the model actually becomes locally a 2D barotropic  
 123 model. The k-epsilon turbulence closure scheme is implemented as in (Michaud et al., 2012).

124 Large scale forcing terms can be provided to the model. The barotropic tidal forcing consists of the harmonic  
 125 tidal components provided by an external tidal dataset, introduced through the open boundary conditions, and of

126 the astronomical tide potential, implemented in the momentum equations according to Pairaud et al. (2008). The  
 127 numerical scheme for the open boundary conditions is described in (Marsaleix et al., 2006); their implementation in  
 128 the present configuration is explained in the Appendix.

129 The air-sea fluxes are computed with the bulk formulae of Large and Yeager (2004) and variables from an atmo-  
 130 spheric model. The daily river discharges are prescribed at the mouth of the rivers or in estuaries and converted into  
 131 horizontal depth-averaged currents.

### 132 2.1.2. The BOBSHELF grid and configuration

133 The Bay of Biscay configuration used in this study is discretized on a curvilinear horizontal grid, implemented  
 134 using the equations and scaling factors described by Madec (2008) (Figure 2a). The resolution in the region of the  
 135 'Pertuis Charentais' and the Gironde estuary is of the order of 300m. Around Brittany, in the northern part of the  
 136 domain, the resolution is degraded. Therefore, this area will not be further discussed in this paper. The bathymetry  
 137 is obtained by merging a GEBCO dataset with several local databases (F. Lyard, pers. comm., 2016). The small  
 138 scales are removed by smoothing. The bathymetry is however not thresholded, in order to be able to represent the  
 139 intertidal zones using the wetting drying scheme of the model. This configuration allows the representation of physical  
 140 processes occurring at different spatial and temporal scales, from the deep plain to shallow areas. In this configuration,  
 141 large scale processes such as tides and shelf circulation can be studied, as well as fine scale processes like waves or  
 142 river plume dynamics.

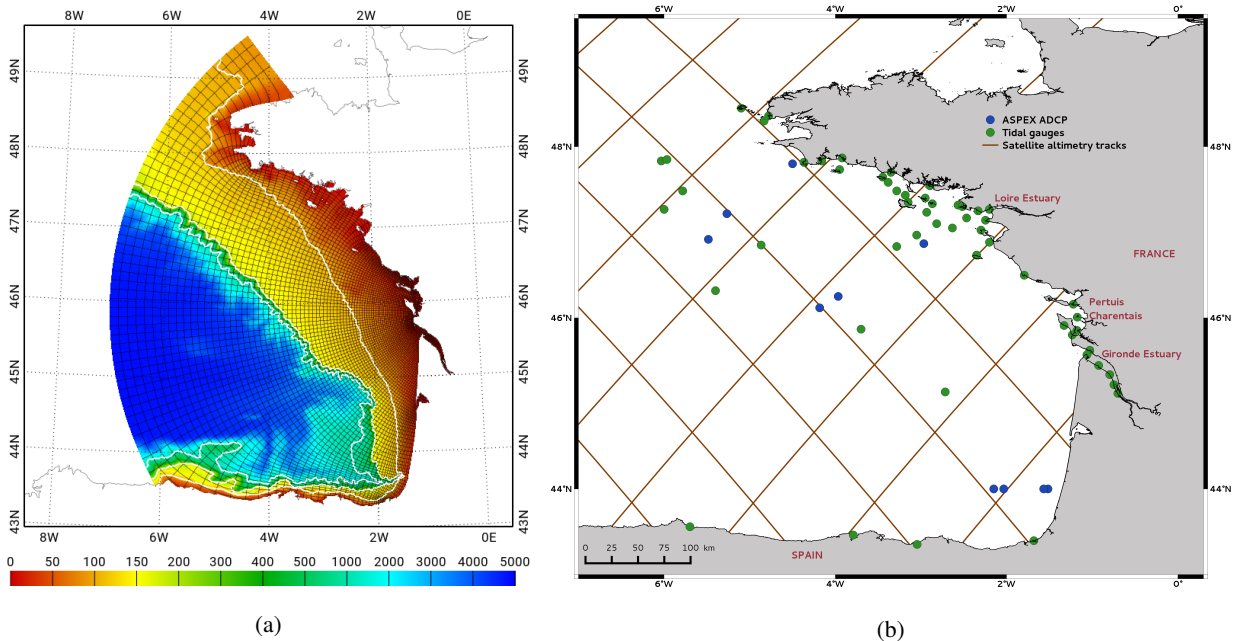


Figure 2: (a): Bathymetry and grid used for the BOBSHELF configuration. In white : 100, 200 and 1000m isobaths. The displayed grid corresponds to the resolution divided by 5. (b): Location of the main in situ data of satellite altimetry tracks used for the model assessment



143 For the 3D configuration used in SYMPHONIE, generalized sigma coordinates are used on the vertical, with  
144 55 levels. The atmospheric forcing variables are provided by the ECMWF operational analyses (6-h fields, 10 km  
145 resolution). The daily river discharges for 6 rivers are obtained from hydrological stations and retrieved through the  
146 French national service 'Banque Hydro' (<http://www.hydro.eaufrance.fr/>).

147 For the tidal circulation, tidal elevations and tidal horizontal currents averaged over the water column are pre-  
148 scribed at the open boundaries, for 9 constituents: M2, S2, N2, K2, K1, O1, P1, Q1 and M4.

149 For the 3D non-tidal circulation, the operational product of MERCATOR-Ocean (the 'IBI' product, with a reso-  
150 lution of  $1/36^\circ$ , see for instance Maraldi et al. (2013)) is prescribed at the open boundaries. The resolution near open  
151 boundaries is about 3km, which is matching the IBI resolution. The IBI forcing consists of daily fields of temperature,  
152 salinity, horizontal velocities and SSH. These fields are averaged over 25h hours to remove the M2 signal. Conse-  
153 quently, most of the tidal signal is filtered out, but we expect some to remain; this question is addressed in section  
154 5.

155 The BOBSHELF configuration used for this study was conceived to study fine scale processes occurring in the Bay  
156 of Biscay, especially the interactions between the Gironde estuary and circulation on the shelf. Going forward, this  
157 modelling setup will be used in particular to assess the observability of such processes by the future satellite altimetry  
158 mission SWOT (Fu et al., 2012; Rodriguez et al., 2017) for oceanography and continental hydrology. In this context,  
159 the question of tides is central, for two reasons. First, tides play a key role in estuarine dynamics by influencing  
160 mixing, resulting in a stronger or weaker stratification at the outlet, and determining the characteristics of the water  
161 masses that can interact with the shelf circulation. It is then essential to have the best possible representation of tides  
162 to study the estuary to ocean continuum. Secondly, as the SWOT mission is being prepared, increasing efforts are  
163 made to provide the best corrections for the different sources of error, including tides. Since one of the objectives of  
164 SWOT is to observe small scale processes close to the coast, i.e. where the tides influence on the sea surface height is  
165 increased, the quality of satellite altimetry data is dependent on the quality of tidal corrections. A good representation  
166 of tidal downscaling up to very coastal areas is then essential to the success of SWOT.

167 Although the prime focus of this study is tides, we also discuss the performance of the simulations regarding  
168 salinity and temperature. Comparisons with buoy SST and SSS are included in Table 2, to show the ability of the  
169 model to correctly reproduce salinity and temperature. The mean error compared to satellite SST (L3S product, Orain  
170 (2016)) is also estimated at 0.53C for the years 2011 and 2012. For this calculation, only the dates when the data  
171 coverage was higher than 50% are considered. Finally, the averaged error on SSS when compared to the gridded  
172 monthly SSS product in the Atlantic Ocean (Reverdin et al., 2007; Alory et al., 2015) is estimated at 0.19 psu, also  
173 for 2011 and 2012.

## 174 2.2. *T-UGOm hydrodynamic model*

175 T-UGOm is a 2D/3D unstructured grid model developed at LEGOS. It can accommodate a variety of numerical  
176 discretization (continuous and dis-continuous finite element, finite volumes) on triangle or quadrangle elements. It

Table 2: Model - data comparison with SST and SSS buoys

	Oléron	Yeu	Houat
RMSE SST (°C)	0.41	0.56	0.41
RMSE SSS (psu)	0.38	0.37	0.61

177 can be used in time-stepping (TS) or frequency-domain (FD) mode. Both TS and FD are routinely run (Pairaud et al.,  
 178 2008; Stammer et al., 2014) for operational ocean high frequency signal (i.e. tides and storm surges) corrections  
 179 in satellite altimetry and gravimetric observations. Initially, the frequency-domain mode has been implemented in  
 180 the original time-stepping T-UGOm code to dynamically downscale tidal boundary conditions for the time-stepping  
 181 simulations. For instance, a Flather open boundary conditions setting needs both tidal elevation and currents to be  
 182 known at the open limits. Where direct interpolation from a global atlas for elevation will be quite suitable, tidal  
 183 currents (when available) are much less likely to be consistent with the nested grid and bathymetry. The frequency-  
 184 domain solver, which is based on a wave equation where tidal currents are not necessarily prescribed along open  
 185 boundaries, allows for reconstructing at a very limited numerical cost a consistent tidal currents field on the nested  
 186 configuration, open boundaries included.

187 The T-UGOm FD 2D solver is originally inspired from the CEFMO frequency-domain tidal model that was earlier  
 188 used for the FES atlases (such as FES2004). FD solver is run for each tidal component separately, it basically assem-  
 189 bles a frequency-domain wave equation and the solution is obtained by a simple inversion of the system. Naturally,  
 190 FD solver is based upon linearized equations, and subsequently non-linear processes require an iterative approach to  
 191 converge toward the fully non-linear solutions. The number of iterations is rather limited for the major astronomical  
 192 tidal components; it tends to increase when addressing compound and non-linear tides. In any case, the numerical  
 193 cost of the FD solver is extremely small compared to the TS solver cost (more than 1000 times smaller). In terms of  
 194 solution accuracy, the FD and TS solvers are quite equivalent, with of course a limited advantage to the TS solver in  
 195 non-linear tides cases. Another major advantage of the FD solver reduced numerical cost is the possibility to conduct  
 196 a wide range of experiments in order to (globally or regionally) calibrate the model parameters such as bottom fric-  
 197 tion and internal tide drag coefficients, verify bathymetry improvements, or test numerical developments. It must be  
 198 noticed that the optimal parameters setting for the FD mode will also meet TS mode requirements.

199 The most commonly used elements in T-UGOm are triangles elements, as they offer the most flexible way to  
 200 discretize the modelling domain with locally adapted resolution. The obvious purpose of implementing quadrangle  
 201 elements is to be able to run T-UGOm FD solver on structured grids, enabling T-UGOm tidal solver to be run on  
 202 most of present structured model configurations. In addition, and in the objective to perform the most consistent tidal  
 203 downscaling, the elevation and current discretization must fit as close as feasible the usual C-grid discretization.

### 204 2.3. *In situ and satellite data for model assessment*

205 Several datasets are used to evaluate the performance of the different simulations (Figure 2b). Along-track tidal  
206 harmonics obtained from a 21-year long time series of satellite altimetry data from TOPEX/Poseidon, Jason-1 and  
207 Jason-2 missions are provided by the CTOH-LEGOS (Birol et al., 2016). The data coverage, extending over the  
208 whole domain, from the deep plain to the shelf, makes it a very valuable dataset of 'spatial tidal gauges' for model  
209 assessment. We also use tidal constituents computed from shorter time-series stemming from the T/P and Jason-1  
210 tandem missions (10/2002-10/2005 and 02/2009-03/2012) on the 'interleaved' tracks; the latter are located midway  
211 along the original tracks thus improving the spatial resolution temporarily. Because of the shorter time-series, the  
212 accuracy on those tidal constituents is degraded but still convenient for the purpose of our comparisons. In the Bay of  
213 Biscay, the harmonic analysis provides M2 tidal elevations with uncertainties of 0.26cm and 0.39cm for the nominal  
214 and tandem mission respectively; such values are very low compared to the M2 elevation that ranges between 1.20  
215 m and 2.10 m in our domain (with a mean value of 1.36 m). These estimations represent the error on the harmonic  
216 analysis in itself. The error is linked to the method that is used to extract the aliased frequencies as precisely as  
217 possible from the ocean background signal. Altimetric data close to the coast undergo a strong loss of accuracy for  
218 several reasons due to instrumental errors and inaccuracies on geophysical corrections; therefore no data is available  
219 at a distance of roughly 50km from the coast.

220 Tidal gauges from the REFMAR, SPC Gironde and Puertos del Estado networks provide both tidal harmonics  
221 and, for a few of them, SSH time series. Other tidal gauges, previously used in Pairaud et al. (2008), also provide tidal  
222 harmonics. These data were obtained through the French Navy and the OHI, but they do not include any indication of  
223 error on the tidal analysis.

224 Harmonic analyses are performed on ADCP current data provided by IFREMER and obtained during the ASPEX  
225 campaign (Le Boyer et al., 2013; Kersalé et al., 2016), for comparison with tidal current harmonics calculated at the  
226 same positions in the model. 2D mean daily currents, projected on the along-shore and cross-shore axes are also  
227 calculated and compared.

228 Because tides impact the regional hydrology through mixing and current rectification, we also compare model  
229 outputs with temperature and salinity data. We use the CORA-IBI database (temperature and salinity profiles) by  
230 IFREMER (Szekely et al., 2017), the Islands network (IFREMER), Météo-France buoys, and the Puertos del Estado  
231 network (sea surface temperature and salinity). A gridded ( $1^\circ \times 1^\circ$ ) database giving monthly estimates of the SSS (sea  
232 surface salinity) in the Atlantic Ocean (Reverdin et al., 2007; Alory et al., 2015) is also used.

## 233 **3. Tidal open-boundary conditions**

### 234 *3.1. The different strategies for tidal OBC tested in this paper*

235 The OBC in SYMPHONIE are based on Flather and radiation conditions whose implementation is described in  
236 the Appendix. External information is needed to specify incoming information. For the barotropic tides, the external

237 information consists in the elevations and horizontal currents averaged over the water column for the nine main tidal  
 238 constituents in the Bay of Biscay: M2, S2, N2, K2, O1, P1, K1, Q1, M4. For the non-tidal circulation, it consists  
 239 in sea surface elevation, 3D temperature and salinity fields, 3D horizontal currents fields for the residual circulation,  
 240 provided by MERCATOR-Ocean (as described in 2.1.2).

241 In this paper, we compare three 7-month simulations of SYMPHONIE 3D, each one forced by a different tidal  
 242 solution (elevation and horizontal depth-averaged currents): 1/ FES2012 atlas, 2/ solution from a SYMPHONIE 2D  
 243 simulation, 3/ solution from a T-UGOm spectral simulation. The performed simulations are summarized in Table 3.  
 244 We emphasize the fact that the only differences between the three runs (S3D\_FES, S3D\_Tugo and S3D\_S2D) are on  
 245 the source of the tidal elevations and currents prescribed at the open boundaries. The OBC numerical scheme is the  
 246 same for the three 3D SYMPHONIE runs.

Table 3: Summary of the performed simulations

<b>2D - forcing solutions</b>			
Tidal forcing (elevations)	Model used	Simulation name	Characteristics
FES2012	SYMPHONIE	S2D	2D clamped No OGCM or atmospheric forcing
FES2012	T-UGOm	Tugo	2D Spectral
<b>3D - circulation solutions</b>			
Tidal forcing (elevations and currents)	Model used	Simulation name	Characteristics
FES2012	SYMPHONIE	S3D_FES	3D OGCM and atmospheric forcing
S2D	SYMPHONIE	S3D_S2D	
Tugo	SYMPHONIE	S3D_Tugo	

### 247 3.2. FES2012 atlas

248 FES2012 is a recent version of the FES (Finite Element Solution) global tidal model (Carrère et al., 2012), follow-  
 249 ing the FES2004 version (Lyard et al., 2006). This model is based on the T-UGOm model (frequency-domain solver  
 250 for the astronomical tides, and time-stepping solver for the non-linear tides) and assimilates tide gauges and satellite  
 251 altimetry derived harmonic constants. Errors both in prior and assimilated solutions have been significantly reduced  
 252 compared to FES2004, especially on the coastal and shelf areas, thanks to a longer time series of altimetric data, a  
 253 more precise bathymetry, and the use of improved data assimilation schemes. The latest FES atlas (i.e. FES2014),

254 despite of superior accuracy, was not used to keep consistent with some already existing simulations forced with  
255 FES2012.

### 256 3.3. 2D simulations

257 As an alternative to the FES2012 tidal atlas, two other tidal forcings are generated to force the 3D circulation  
258 model. The main advantage of these forcings is that they are generated on the same grid (BOBSHELF) and with the  
259 same bathymetry as the ones used in the 3D simulations.

260 To generate tidal boundary conditions, both for tidal elevations and currents, it is possible to run simplified (no  
261 atmospheric or OGCM forcing) 2D simulations with SYMPHONIE. These simulations are performed with clamped  
262 (or Dirichlet) conditions, meaning that only tidal elevations are used at the boundaries (in our case, FES2012 tidal  
263 elevations). Tidal currents are not considered. The model is run for 7 months, from October 2010 to April 2011, with  
264 9 tidal constituents (M2, S2, N2, K2, K1, O1, P1, Q1 and M4). The harmonic analysis is run online. 7 months is the  
265 necessary period to be able to separate the different waves.

266 For the 2D T-UGOm spectral simulations (called Tugo hereafter), the clamped conditions and FES2012 tidal  
267 elevations are also used at the open boundaries. For consistency, the bottom friction is set at the same value in  
268 T-UGOm as in SYMPHONIE.

269 In the next section, we compare and evaluate the three tidal boundary conditions (FES2012, S2D and Tugo) by  
270 comparing tidal elevations to the available observations. In section 4, the results of the 3D simulations are detailed,  
271 first focusing on the tidal solution (elevation and currents), then on a wider range of parameters, to assess the influence  
272 of the tidal forcing on the global circulation.

### 273 3.4. Assessment of the tidal forcing solutions

274 Comparisons with tidal harmonics from satellite altimetry data and tidal gauges are presented in Figure 3. These  
275 bar plots represent the mean complex error, which accounts for errors both in amplitude and in phase. The complex  
276 error  $H_s$  is calculated as follows:

$$277 H_s = \sqrt{h_1^2 + h_2^2} \quad (1)$$

$$278 h_1 = H_m \cos(G_m) - H_o \cos(G_o) \quad (2)$$

$$279 h_2 = H_m \sin(G_m) - H_o \sin(G_o) \quad (3)$$

277 With  $H_m$  and  $G_m$  the amplitude and phase of the modelled constituent, and  $H_o$  and  $G_o$  the amplitude and phase of  
278 the observed constituent.

279 Satellite altimetry results (Figure 3b) show a better agreement with FES2012 for M2 and M4. This is expected  
280 because of the assimilation of these data in FES2012. The S2 tide is better represented by S2D and Tugo. This is  
281 because the S2 signal captured by the altimetry can be divided into two parts: an astronomic one, and an atmospheric

282 one, which is partly non-stationary. As a consequence, the signal assimilated in FES2012 corresponds to a sort of  
 283 S2 'average' that does not take into account the seasonal variability of this constituent. This can explain why the  
 284 performance of the S2 assimilation is relatively limited compared to that of other constituents. More details can be  
 285 found in Lyard et al. (2006).

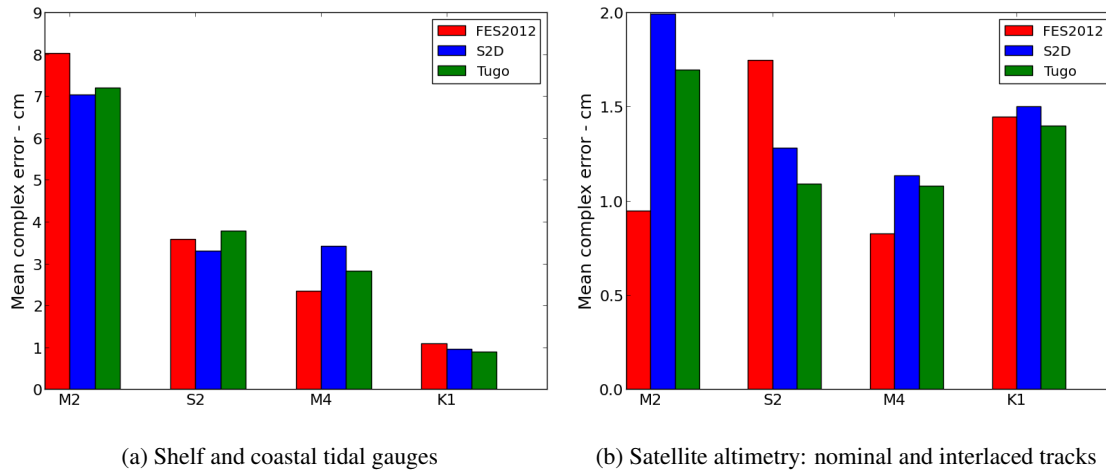


Figure 3: Mean complex errors (in cm) over the BOBSHELF domain for the M2, S2, M4 and K1 harmonics, between the forcing solutions FES2012, S2D and Tugo, and available observations

286 Tidal gauges results (Figure 3a) suggest that the regional models are more accurate near the coast for the M2  
 287 tide. However, the M4 tide is better represented by FES2012. M4 is generated by the interaction of M2 with itself.  
 288 Therefore, a part of the errors on M4 is a direct consequence of the errors on M2. In FES2012, the assimilation is  
 289 performed independently for each constituent, meaning that the M2 solution is not directly impacting the M4 solution,  
 290 reducing the level of errors on this constituent. In S2D and Tugo, there is no assimilation, which means that the M4  
 291 errors are partly inherited from M2, and squared.

292 For K1, the level of error is similar for the three solutions and for the two datasets (tidal gauges and satellite  
 293 altimetry). In the Bay of Biscay, the signal to noise ratio in satellite altimetry is smaller for K1 than for M2, because the  
 294 K1 tide has a much weaker amplitude and a much lower aliased frequency (180 days against 62). As a consequence,  
 295 assimilating the K1 altimetric signal has almost no impact in this region.

296 As shown in Figure 3, the level of error between the two data-model comparisons (tidal gauges and satellite  
 297 altimetry) is significantly different. For example, the M2 complex error is multiplied by more than 4 (for Tugo). Tidal  
 298 gauges are mostly located on the coast, and a few are on the shelf. In these areas, the tidal signal is also amplified when  
 299 compared to the open ocean, where most of the satellite data are obtained from. For instance, the M2 amplitude is  
 300 doubled between the open ocean and certain coastal areas. In addition, near the coast, the evolution of the tidal signal  
 301 is very sensitive to the geometry of the area (coastline, bathymetry). This means that the signal can be significantly  
 302 different between two close locations, whereas in the open ocean, tides are relatively homogeneous in space. This

303 strong spatial variability is difficult to reproduce in a model where the mesh size is larger than the length scale of  
 304 bathymetric features. One model cell is then susceptible to discretize an area where the tidal signal would present  
 305 gradients at a smaller scale.

#### 306 4. Assessment of the 3D SYMPHONIE simulations

307 In this section, we evaluate the impact of using different tidal boundary forcing on the 3D SYMPHONIE model re-  
 308 sults by comparing the simulations to observations from different data sets. The simulation closest to the observations  
 309 (within the data uncertainties range) is obviously identified as the most realistic one. These comparisons also stand  
 310 for a more general assessment of the BOBSHELF configuration: we show indeed that the model-data misfits for the  
 311 different variables are low over the period of study which makes us confident in the ability of the BOBSHELF con-  
 312 figuration to simulate the main processes of the Bay of Biscay circulation with the needed accuracy for our purposes  
 313 (section 2.1.2).

##### 314 4.1. Tidal elevations

315 Complex errors between modelled and observed tidal amplitudes and phases are given in Table 4. Misfits from M2  
 316 altimetry are reduced by more than 75% between S3D\_FES and S3D\_Tugo (70% for S3D\_S2D). For the tidal gauges,  
 317 the error is reduced by 20% (13 % for S3D\_S2D) . Altimetry errors are also lower for S2 with the Tugo and S2D  
 318 forcings. For M4 and K1, the errors on altimetry are slightly higher for S3D\_S2D and S3D\_Tugo than for S3D\_FES  
 319 (between 8 and 15%). However, the difference is more significant for the tidal gauges comparison on M4 (more than  
 320 +38%). For K1, S3D\_S2D and S3D\_Tugo simulations perform a little bit better than S3D\_FES (between -11% and  
 321 -17%).

Table 4: Mean complex errors (cm) between the 3D circulation solutions and the available data for M2, S2, M4 and K1. SA = Satellite altimetry ;  
 TG = Tidal gauges.

3D simulations	M2		S2		M4		K1	
	SA	TG	SA	TG	SA	TG	SA	TG
S3D_FES	7.49	9.30	1.89	3.42	0.993	2.94	1.29	1.28
S3D_S2D	2.10	8.03	1.56	3.80	1.15	4.40	1.49	1.06
S3D_Tugo	1.84	7.39	1.55	3.73	1.11	4.07	1.42	1.13

322 A more detailed view of the complex errors on M2 is available in Figure 4: all the tidal gauges (shown in Figure  
 323 2b) used to compute the mean error are represented, from the northern to the southwestern limit of the domain. All  
 324 tidal gauges between Ouessant and Gijon are either directly on, or very close to the coast (depth < 50m). The T1 to  
 325 'COURIR5 3' gauges are located on the shelf (data from Le Cann (1990) and SHOM dataset). 'MGFCOR large' is  
 326 the only gauge in the deep plain (depth > 4000m).

327 This figure is consistent with the results presented in Table 4: M2 errors are globally much lower in S3D\_Tugo  
328 and S3D\_S2D than in S3D\_FES. Out of 52 tidal gauges, the M2 complex error is the highest for S3D\_FES in 33 of  
329 them. These errors also show a strong spatial variability. Tidal gauges within a few kilometers of each other can  
330 display a very different trend. For example, the 'Birvideaux' and 'Vilaine P1' tidal gauges are distant of less than  
331 10km (see Figure 5 for locations). The first one shows that S3D\_FES has the highest error, whereas the second one  
332 suggests the opposite. This underlines the difficulty to compare single-point data like tidal gauges records to numerical  
333 simulations, as already discussed in the previous section.

334 In this sense, satellite altimetry data appear to be more useful. First, it provides a spatially homogeneous dataset  
335 that covers a larger area, allowing us to check if the data is consistent within a few kilometers. In terms of data  
336 quality, it is also important to notice that tidal gauges data are obtained from different providers, with differences in  
337 instrumentation and data processing. It is of course necessary to use tidal gauges to evaluate regional models close  
338 to the coast. In this modelling configuration in particular, the resolution is increased near the coast. Tidal gauges  
339 comparisons must be made, because satellite altimetry cannot yet provide reliable data in these areas. However, it is  
340 important to be aware of the challenges associated with these data, that can seem easier to use than satellite altimetry  
341 at first.

342 The Royan and Port-Bloc tidal gauges are located at the mouth of the Gironde estuary, each on one side of the  
343 river. The Richard, Lamena, Trompeloup and Fort-Médoc gauges are distributed within the estuary, Fort-Médoc being  
344 the most upstream point. At the Royan, Port-Bloc and Lamena stations, S3D\_FES seems to provide the best solution,  
345 but this tendency is reversed when moving upstream, from Lamena to Fort-Médoc. The Gironde estuary is a particular  
346 environment, with the presence of a turbidity maximum and fluid mud (Sottolichio and Castaing, 1999). The latter  
347 can induce large variations in bed roughness, which influences tidal propagation and distortion within the estuary: in  
348 the presence of fluid mud, the bed roughness is very low, inducing an increase of velocities, and a reduced damping  
349 of the M2 tide. In order to correctly represent tides within the estuary, bottom friction tuning will be necessary.  
350 In our configuration where no specific tuning is done (that would be out of the scope of this paper) and in which  
351 sedimentary processes are not taken into account, we cannot expect to represent accurately tides in the estuary. The  
352 fact that S3D\_FES is closer to the observations in the lower estuary than S3D\_Tugo and that the opposite is found in  
353 the upper estuary seems paradoxical. However, given the limitations of the configuration there, we suggest that error  
354 compensating effects are responsible for the good match of S3D\_FES with observations at the estuary mouth.



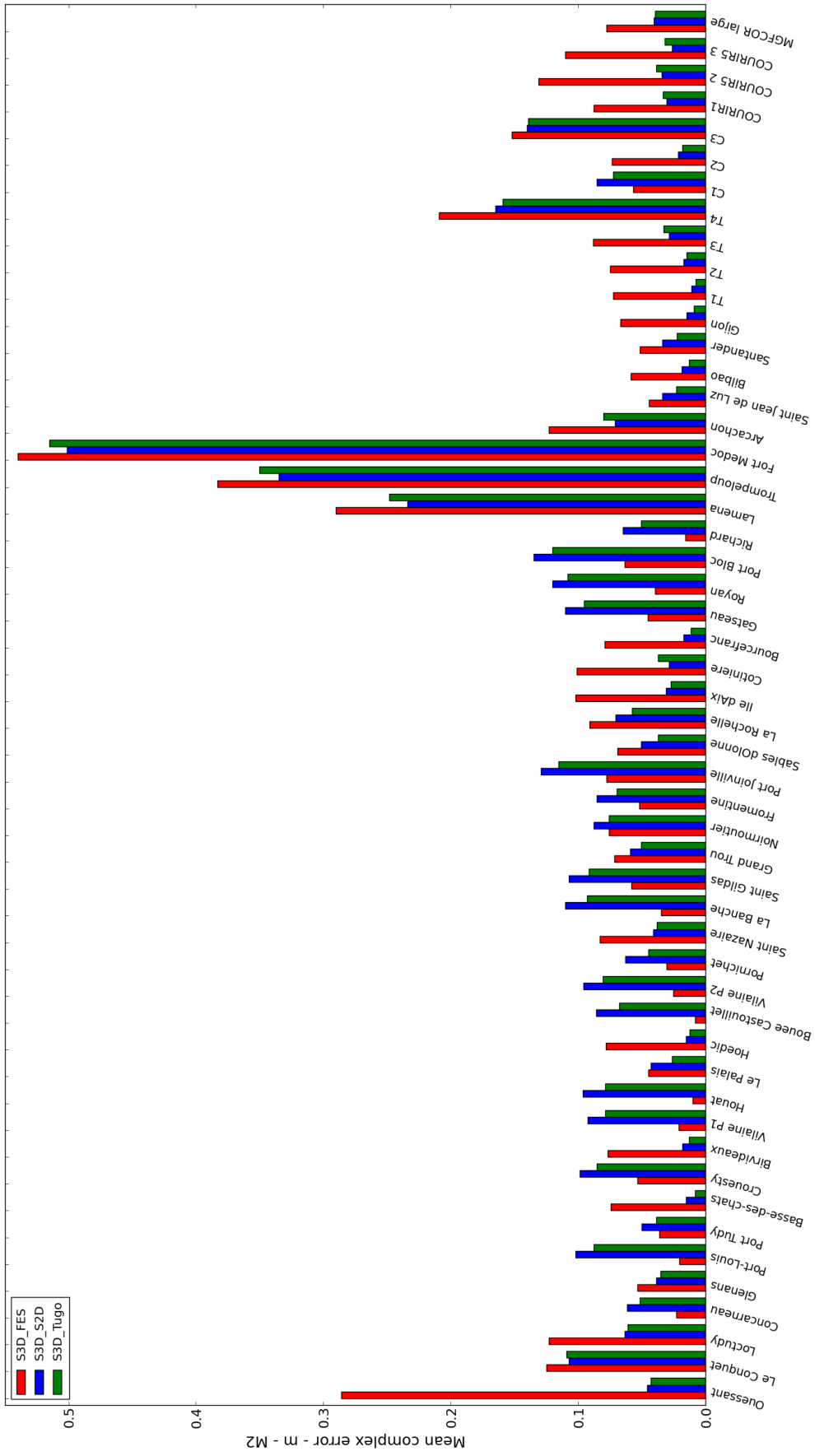


Figure 4: Complex errors for the M2 tide, between the 3D simulations and all the available tidal gauges

355 Figure 5 represents the differences in amplitude and in phase, for M2, between S3D\_Tugo and S3D\_FES. In  
 356 amplitude, the highest differences (more than 7 cm) are obtained on the shelf, in the northern part of the domain,  
 357 where the tides are highly energetic. Overall, the difference is close to 5 cm in amplitude. This value is decreasing  
 358 within the Gironde estuary, suggesting again that the dynamics inside the estuary are more constrained by local effects  
 359 (bathymetry, friction), than by the remote forcing. The same behavior is obtained for the M2 phase difference, which  
 360 is the lowest inside the estuary.

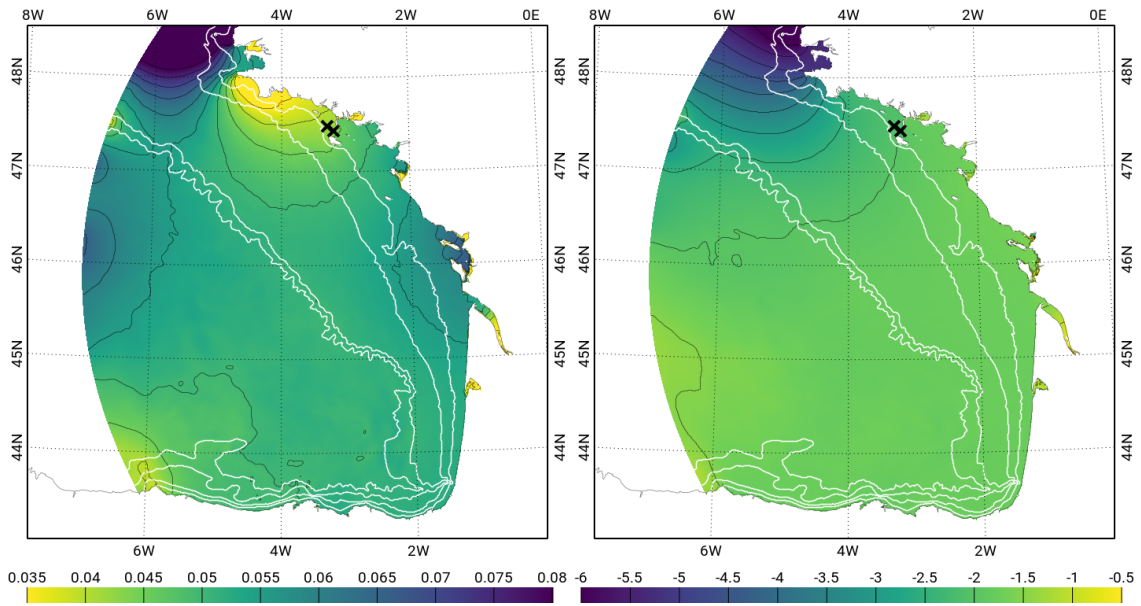


Figure 5: Difference between S3D\_Tugo and S3D\_FES for the M2 amplitude (left, m) and phase (right, °). In white: 50, 100, 200 and 1000m isobaths. Black crosses: 'Vilaine P1' and 'Birvideaux' tidal gauges.

#### 361 4.2. Tidal currents

362 Tidal analyses are performed on ADCP data obtained during the ASPEX campaign (Le Boyer et al., 2013; Kersalé  
 363 et al., 2016), and tidal ellipses parameters are then calculated. 3D velocity currents are averaged over depth to perform  
 364 a 2D analysis. No evidence of a significant vertical structure of the tidal currents was found. For consistency, only the  
 365 ASPEX ADCP data covering the 7 months simulation period are used for comparison.

366 Figure 6 shows a global good agreement between the three S3D simulations and ASPEX data, both for M2 and  
 367 M4, making it difficult to draw conclusions on the best simulation. S2 ellipses (not shown) exhibit results close to  
 368 the M2 ellipses. Significant direction differences between model and data for ASPEX1 and ASPEX3 are found for  
 369 M2 (Figure 6a). The M4 tidal ellipses seem a little bit more different between the three solutions (Figure 6b). For  
 370 ASPEX1, the S3D\_S2D and S3D\_Tugo ellipses are closer to data than S3D\_FES. For ASPEX3, the direction of the  
 371 ellipse is better reproduced in S3D\_Tugo. For ASPEX5 and ASPEX9, the semi-major axis is also closer to data in  
 372 S3D\_Tugo than in the other two simulations.

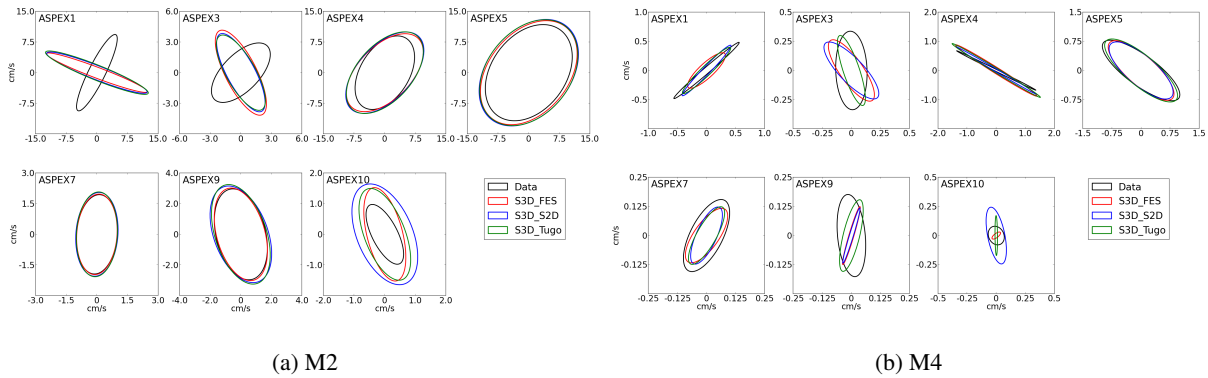


Figure 6: Tidal current ellipses for the M2 and M4 tide: comparison between data and S3D simulations at each ASPEX mooring.

### 373 4.3. Sea surface height

374 The SSH time series of 11 tidal gauges are compared to each 3D simulation. Not all the gauges used for the tidal  
 375 harmonics comparisons provide SSH time series, explaining why the number of data points used here is reduced.  
 376 Standard deviations (STD) and root mean square errors (RMSE) are normalized (divided by the data standard devia-  
 377 tions), to represent all the tidal gauges used on the same Taylor diagram (Figure 7). The three simulations give good  
 378 results, with correlation coefficients greater than 0.99 and normalized RMSE between 5 and 15. Normalized STD  
 379 show a greater agreement between data STD and S3D\_S2D or S3D\_Tugo than S3D\_FES.

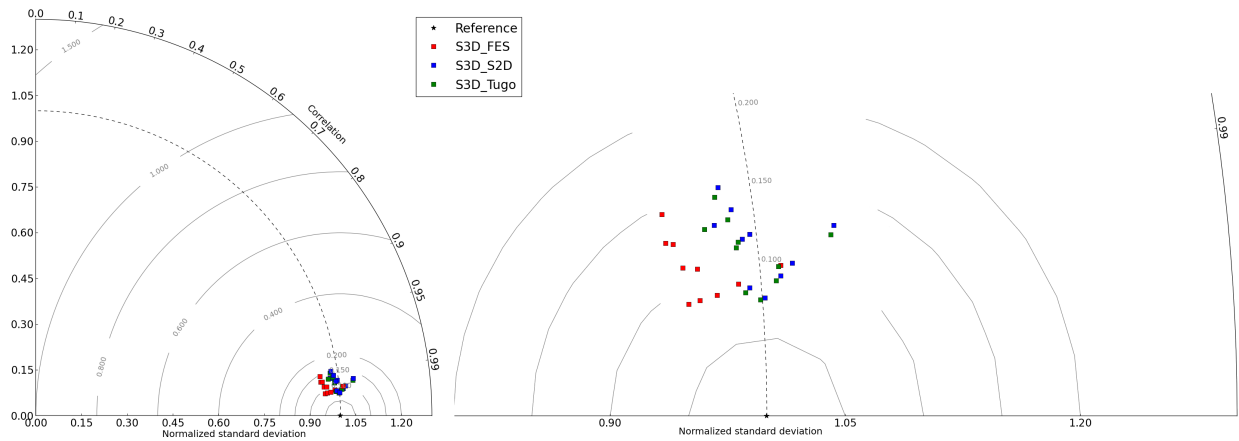


Figure 7: Taylor diagram obtained from the comparison of 11 tidal gauges time series with model outputs. On the right: zoom on the lower right corner.

### 380 4.4. Current velocities

381 The ADCP data used previously to compare tidal ellipses are now processed to compare total currents. 2D daily  
 382 means are calculated in the along-shore and cross-shore directions and compared. Mean RMS errors and correlations  
 383 are calculated over 9 ADCPs (Table 5).

Table 5: Model - data comparison between 2D daily means of cross-shore and along-shore current velocities

	S3D_FES		S3D_S2D		S3D_Tugo	
	<i>RMSE</i>	<i>cōrr</i>	<i>RMSE</i>	<i>cōrr</i>	<i>RMSE</i>	<i>cōrr</i>
Along-shore ( $cm.s^{-1}$ )	4.96	0.448	4.77	0.510	4.55	0.585
Cross-shore ( $cm.s^{-1}$ )	2.94	0.125	2.89	0.188	2.95	0.202

384 For the along-shore currents, the mean RMSE is lower for S3D.Tugo and S3D\_S2D than for the S3D\_FES simu-  
385 lation. The difference is rather small (less than 10 %). However, the mean correlation is increased by more than 30  
386 % between S3D\_FES and S3D\_Tugo (14 % for S3D\_S2D). The correlations are significant (95 % confidence level) in  
387 8 out of 9 moorings for S3D\_Tugo, against 7 for S3D\_S2D and 6 for S3D\_FES. Besides, the global numbers of table  
388 5 do not represent local behaviors. The impact of the tidal forcing is indeed different on the different moorings and  
389 changing in time. This is illustrated on Figure 8 for the ASPEX9 mooring which is located over the slope at the 44°N  
390 section: the variability of the signal is better reproduced by the S3D\_Tugo simulation, particularly at the beginning  
391 and the end of January 2011, and at the beginning of March 2011. To make sure that these differences are not due to  
392 very local features, the results are plotted representing the median value of 9 grid points (data co-located grid point  
393 and 8 surrounding grid points). The minimum and maximum values from these 9 points are also represented by the  
394 shaded area. This type of data-model comparison can be called 'fuzzy verification', and comes from the fact that high  
395 resolution models often score poorly with single-point comparisons, while their performance is often very satisfying.  
396 Thus, fuzzy verification allows the model to be slightly displaced (here in space), and still be valuable (Ebert, 2008).

397 Cross-shore currents are very weak (0-1 cm/s on average). Correlation values are not significant and cannot be  
398 compared between the different simulations.

#### 399 4.5. Temperature and salinity

400 Modelled temperature and salinity are evaluated with respect to several databases of surface measurements (Is-  
401 lands network, Puertos del Estado network, Météo-France buoy) and in situ profiles from the CORA-IBI database  
402 (Szekely et al., 2017). Time series of model-data comparisons and statistics (RMSE, STD, correlation, not shown) are  
403 inconclusive, because the differences between the S3D simulations are within the in-situ measurements uncertainties.

404 Comparisons with the gridded monthly sea surface salinity product in the Atlantic Ocean (Reverdin et al., 2007;  
405 Alory et al., 2015) show some differences between the 3D simulations. For the 7 months runs, considering a 3 months  
406 spin-up for the 3D circulation, comparisons are made from January to April. Monthly means of the model SSS are  
407 computed over the same grid as the observations. Noticeable differences between the simulations are found for two  
408 points only, located on the shelf, and under the influence of the Gironde estuary (Figure 9a, hereafter 2W46N) and  
409 Loire estuary (Figure 9b, hereafter 3W47N) plumes.

410 For the 2W46N location, in January and February, the S3D\_FES simulation seems to be in better agreement with

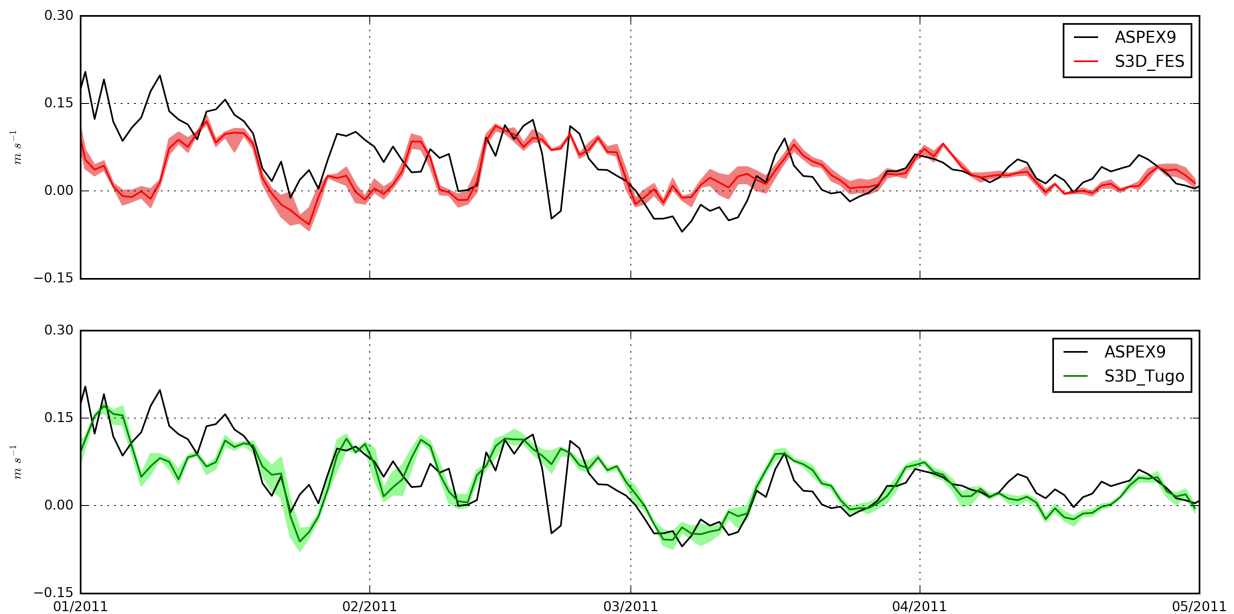


Figure 8: 2D mean daily current velocities for the S3D.FES (top) and S3D.Tugo simulations (bottom), compared to ADCP data from Le Boyer et al. (2013). Solid line: median of 9 grid points. Shaded area: minimum and maximum values from 9 grid points

411 the observations. On the other hand, in March, the S3D.S2D and S3D.Tugo results are very close to observations. In  
 412 April, all three simulations show a similar level of performance, although the S3D.S2D and S3D.Tugo simulations  
 413 are slightly better. For the 3W47N point, S3D.Tugo has the overall lowest error budget, while there is a systematic  
 414 underestimation of the sea surface salinity for the 3 simulations.

415 In summary, comparisons of data-model misfits between the 3D simulations do not show strong differences.  
 416 However, it does not mean that there are indeed no differences in temperature and salinity between these simulations.  
 417 A more extensive dataset would be necessary to reach a more straightforward conclusion.

## 418 5. Discussion

### 419 5.1. A downscaling challenge: an accurate representation at all scales

420 Based only on the results obtained on the tidal forcings (FES2012, S2D and Tugo), the FES2012 forcing seemed  
 421 to be the best option, because of a clear lower error budget for M2 when compared to satellite altimetry data. How-  
 422 ever, the results obtained with the 3D simulations show that the Tugo forcing provides an overall better solution, both  
 423 for tidal elevations and currents. The difference of performance between the forcing FES2012 and the 3D solution  
 424 forced by FES2012 (S3D.FES) is particularly striking with respect to satellite altimetry: the complex error is almost  
 425 multiplied by 8 (compare Figure 3b and Table 4). This gap in performance between FES2012 and S3D.FES can be  
 426 imputed to the differences in resolution and bathymetry at the open boundaries. In other words, there is an inconsis-  
 427 tency between the FES2012 currents at the open boundary conditions and the tidal dynamics inside the domain that

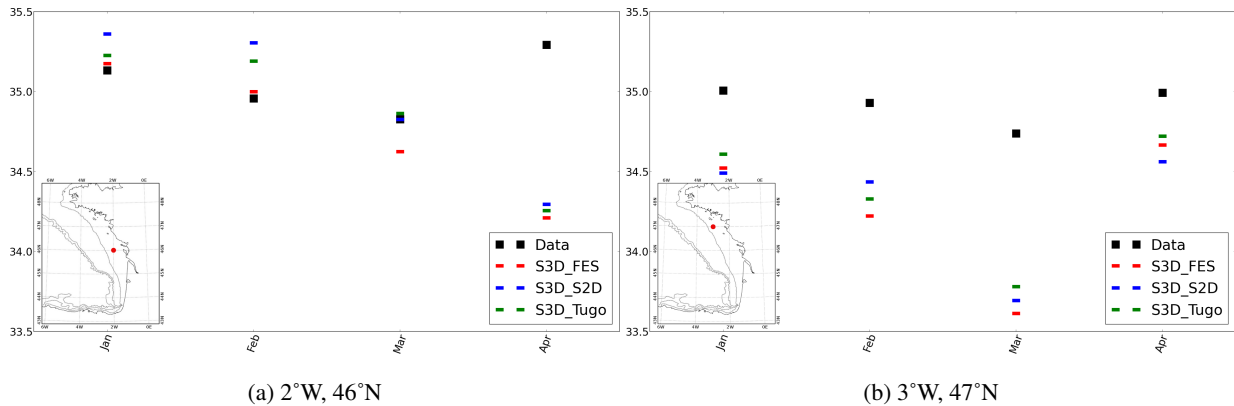


Figure 9: Monthly sea surface salinity (psu) comparisons between the Atlantic Ocean database and the S3D simulations, at the two points shown on the maps

428 are constrained by the resolution and bathymetry of the BOBSHELF configuration. These over-specification errors  
 429 are mainly impacting the M2 tides because the M2 currents are largely dominant.

430 Over-specification is a recurrent pitfall with OBC. It occurs here because of the high sensitivity of tidal currents to  
 431 bathymetry and to the detailed representation of the coastline inside the domain. To our knowledge, there is no OBC  
 432 scheme that prevails to systematically reduce the over-specification errors. Continuous efforts are developed in the  
 433 community as configurations evolve with higher and higher resolution representing more and more complex processes.  
 434 The recent study of Herzfeld and Andrewartha (2012) for instance proposes a method based on the Dirichlet conditions  
 435 in conjunction with a local flux adjustment for volume conservation that has been successfully tested in different  
 436 configurations and that requires little tuning. Robustness and simplicity are of course highly attractive qualities for a  
 437 scheme to be implemented in complex systems. Over-specification may be enhanced in the presence of tides together  
 438 with a low-frequency circulation (such as in our case) because active boundary conditions for tides can be reflective  
 439 for the low-frequency circulation as noted by Herzfeld and Gillibrand (2015). These authors investigate a method  
 440 based on dual-relaxation time-scales to solve possible conflicts on the passive/active nature of open boundaries for  
 441 the different components of the flow. Such new research offers interesting perspectives of improvement for the OBC  
 442 schemes that we may want to explore in future studies. In the present one, we have chosen to adjust the external  
 443 forcing field in order to reduce the inconsistencies with the interior solution rather than modifying the OBC scheme.

444 The modelled M2 field for S3D\_FES and S3D\_Tugo is shown in Figure 10: the complex error between these simu-  
 445 lation results and the tidal harmonics extracted from satellite altimetry data is represented by the circles superimposed  
 446 on the maps. Figure 10 shows that the errors are well distributed over the whole domain, and not only at the open  
 447 boundaries. In other words, the change of tidal open-boundary forcing has a significant impact on SSH all over the  
 448 basin. More precisely, the M2 amplitude is globally underestimated in S3D\_FES, and the M2 phase is overestimated.

449 The better performance with S3D\_Tugo results from the use of forcing tidal currents that are consistent with the  
 450 interior resolution and bathymetry, since the Tugo model has been run on the BOBSHELF grid. This is illustrated

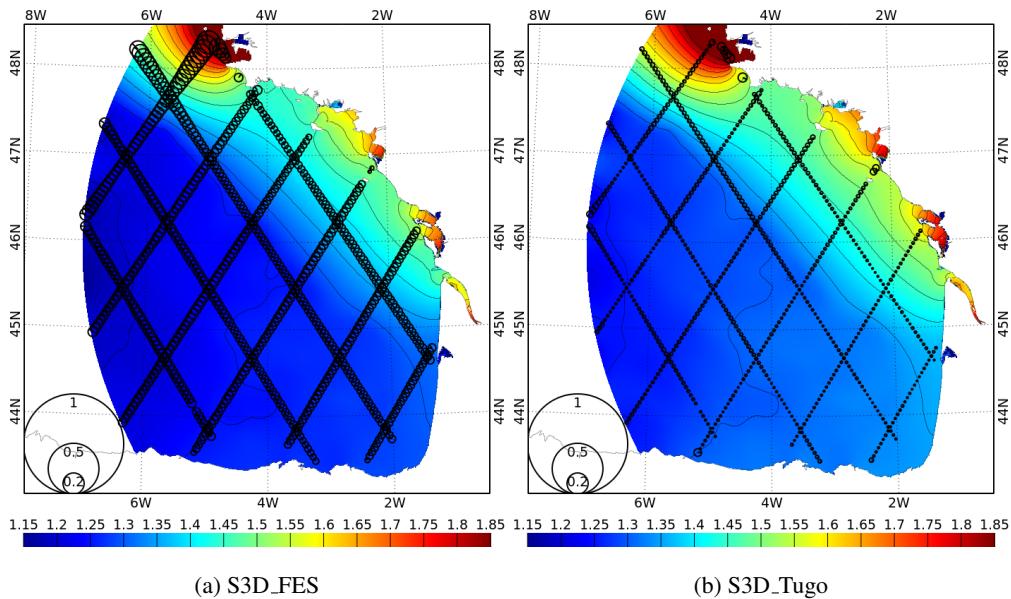


Figure 10: Comparison between the M2 tide obtained from satellite altimetry and from S3D. In the background: M2 amplitude (m). The circle size is proportional to the complex error (m).

451 on Figure 11a, showing the difference between the M2 tidal current amplitude in FES2012 and Tugo. Differences  
 452 larger than 10 cm/s are observed over the northern slope and shelf (north of 45°N) where the tidal currents are strong  
 453 and the resolution increased. In particular, the difference between the 2D M2 tidal current in FES2012 and Tugo can  
 454 reach more than 20 cm/s close to the Armorican slope, at the north-west open boundary of our domain (Figure 11a).  
 455 Between S3D\_FES and S3D\_Tugo (Figure 11b), these differences are reduced, but they can still reach more than 10  
 456 cm/s. In the deep plain, the differences are much smaller in the 3D simulations. S3D\_S2D also gives better results  
 457 than the S3D\_FES simulation, but with a level of error that is slightly above the S3D\_Tugo simulation. As a reminder,  
 458 T-UGOm was specifically designed to model and study tides, unlike SYMPHONIE, which is dedicated to circulation  
 459 simulations. In this sense, it is logical that the performance of the Tugo forcing would be higher, although the results  
 460 obtained with S2D are already an improvement from using the FES2012 atlas.

461 As an attempt to evaluate the impact of the tidal forcing on the hydrology and residual dynamics inside the  
 462 domain, we have compared the simulations to the observations at our disposal. We found a significant impact on the  
 463 representation of SSS at large scale ( $1^\circ \times 1^\circ$ ) and for monthly averages over the shelf. Possible mechanisms include the  
 464 advection by surface current and the vertical mixing which is strongly influenced by tides over the shelf. On the other  
 465 hand, comparisons to T, S profiles from single point measurements were inconclusive; this is not surprising as model-  
 466 data misfits at single points comprehend many possible sources of error, such as co-localization, that dominates the  
 467 influence of the tidal solutions on stratification. Similarly, SSH time series comparisons showed very little differences  
 468 in terms of RMSE or correlation, but the STD of the observed signal was better reproduced by S3D\_Tugo ; this  
 469 suggests that the SSH variability at the coast is better represented in this simulation. Finally, the daily 2D means of

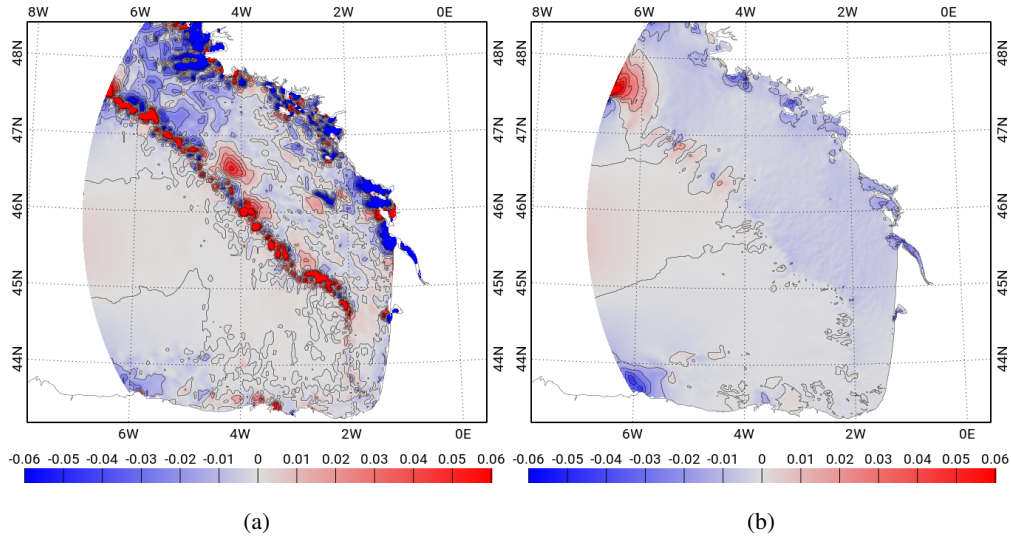


Figure 11: Difference between the amplitude of the 2D M2 current in: (a) FES2012 and Tugo, (b) S3D\_FES and S3D\_Tugo (m/s)

470 the along-shore current velocities from the ASPEX campaign time series are better correlated to the S3D\_Tugo field  
 471 than to the S3D\_FES field; we evidenced episodes of a strong improvement of the model slope current in S3D\_Tugo  
 472 with respect to the data. An attempt is made to provide an interpretation to this improvement in the next section.

### 473 5.2. Tidal forcing influence on the circulation

474 In section 4.4, 2D daily currents obtained from ADCP data are compared to modelled currents. These comparisons  
 475 clearly show a better agreement with the S3D\_Tugo simulation, even though the tidal ellipses calculated at the same  
 476 location (section 4.2) are very close between the three simulations. This suggests that, at these locations (far from  
 477 the open boundaries), the tidal forcing seems to have a greater influence on the global circulation than on the tidal  
 478 currents themselves. The difference between the S3D\_FES and S3D\_Tugo M2 tidal current (Figure 11b) confirms this  
 479 observation: on the shelf and slope in the southeastern part of the Bay, the differences are very small compared to the  
 480 large differences observed at the open boundaries.

481 To explore further the impact on the mean circulation, we compare the mean currents over ten days, from daily  
 482 detided fields. Figure 12 shows the results obtained at a 50m depth, in the southern part of the domain, for the  
 483 S3D\_FES and S3D\_Tugo simulations, over the first 10 days of January. This period has been chosen because it  
 484 corresponds to the one where the differences at the ASPEX 9 mooring are the largest (Figure 8). It also coincides with  
 485 the occurrence of an eastward along slope current with an amplitude locally larger than 25 cm/s. The associated SST  
 486 field displays a warm water tongue along the Spanish coast that extends northward along the slope up to 45° N off the  
 487 French coast. This SST patterns is consistent with the classical view in the literature of a poleward slope current in  
 488 early winter advecting warm water masses (see for instance Pingree and Le Cann (1992)). On Figure 12, differences  
 489 between the simulations, at small scales in the mesoscale field, for both current and SST, can be interpreted as a



490 'stochastic' response of the turbulent flow to the small perturbations of the open-boundary conditions. On the other  
 491 hand, the mean current over the Spanish slope seems to be affected by the change in tidal boundary conditions. In the  
 492 S3D.Tugo simulation, the mean current is more constrained to the upper part of the slope, especially between 5°W  
 493 and 4°W, and 3°W and 2°W.

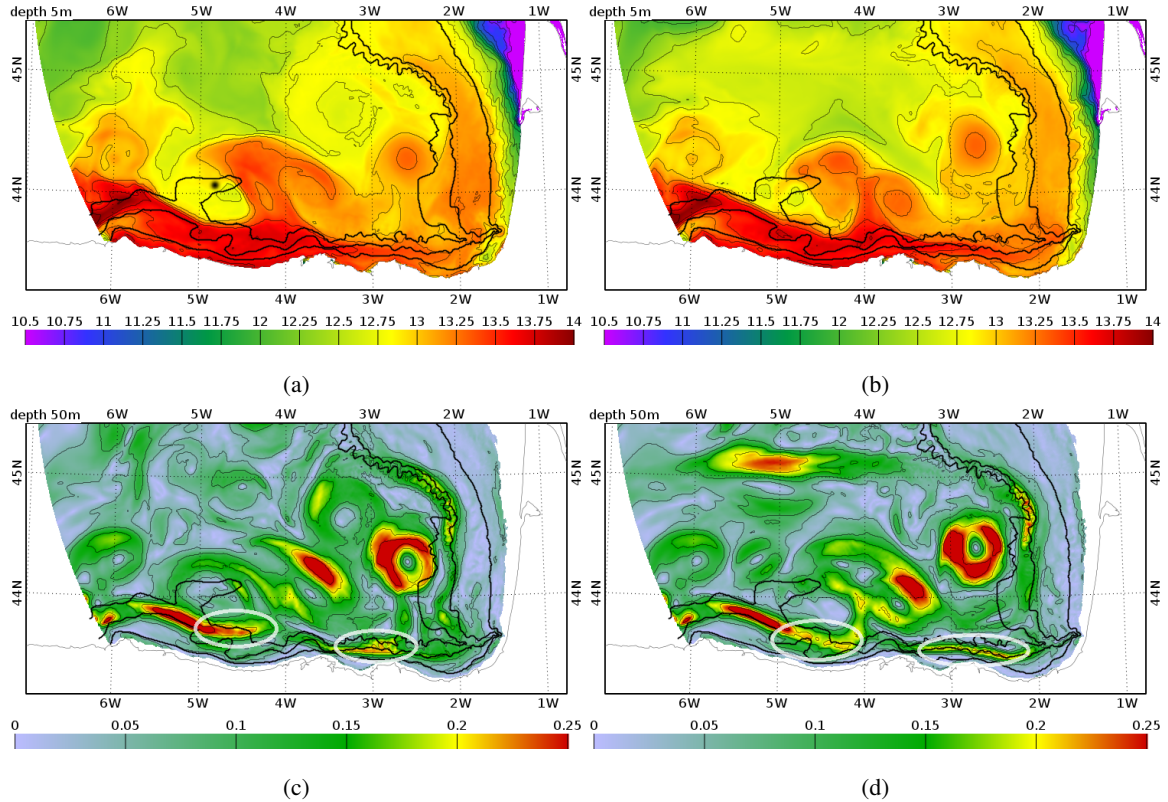


Figure 12: Comparison between 10-day means (January 1 to 10) of sea surface temperature (°C) and current velocities (m/s) at 50m from S3D.FES (a,c) and S3D.Tugo (b,d). The Le Danois Bank is marked with a black dot in (a). In black: 100, 200 and 1000m isobaths.

494 These different features are observed just east from the Le Danois Bank (approximately at 44°05'N, 4°50'W),  
 495 which is a seamount-like topographic feature at about 60 km of the northern Spanish coast (González-Pola et al.,  
 496 2012), quite close to the open boundaries of our configuration. From in-situ measurements, González-Pola et al.  
 497 (2012) show that diurnal tides (K1 and O1) are strongly amplified, both on the northern and southern sides of this  
 498 seamount. The authors interpret the amplification of the diurnal tides as resulting from the generation of resonant  
 499 seamount trapped waves. They also argue that the amplification of K1, O1, and, to a lesser extent, of M2, could be an  
 500 indicator of topographic tidal rectification, that can generate a mean residual current.

501 Tidal rectification occurs when the nonlinear terms in the momentum equation become of the order or greater than  
 502 other forcing terms; the processes involved are a combination of continuity and Coriolis effects and bottom friction  
 503 (Loder, 1980). Shelf breaks or seamounts in macrotidal environments, with a strong cross-isobath tidal flow, are

504 favourable sites for the occurrence of topographic tidal rectification. For instance, Garreau and Maze (1992) derive  
505 analytical solutions for eulerian currents generated by the rectification of the M2 tides over a slope: they find that the  
506 solution is consistent with observed residual current at the top of the shelf break in the northern Bay of Biscay.

507 The Le Danois Bank western and northern slopes are a priori favorably oriented so that M2 tides can indeed  
508 rectify. Based on rough scaling arguments, González-Pola et al. (2012) show that M2 rectified flow could reach there  
509 a few centimeters per second.

510 In our simulation, because of the multiple forcing terms that are resolved by our model, it is difficult to isolate tidal  
511 rectification processes (a thorough analysis of tidal rectification in our simulations is beyond the scope of this study).  
512 But in view of the studies in the literature and of the general topographic characteristics, we find very likely that topo-  
513 graphic tidal rectification indeed occurs in our runs in the Le Danois area. In our simulations, the amplification of K1  
514 and O1 does occur, with a stronger intensity in S3D\_Tugo. Because of the proximity to the open boundary, the tidal  
515 current shows significant differences there between the S3D\_FES et S3D\_Tugo runs (Figure 11). We therefore expect  
516 the rectified flow to be different as well. The early January period is characterized by a mean along-slope current, at  
517 least partially originated outside of our domain. We suggest that tidal-mean current interactions and topographic rec-  
518 tification mechanisms impact the along-slope mean flow; the observed difference further downstream, at the location  
519 of the ASPEX 9 mooring (Figure 8) would then result from this impact. We also performed a simulation without tides  
520 (not shown), and found that the 10-day mean currents and SST were significantly modified in this part of the Bay, thus  
521 confirming the strong impact of tides on circulation.

522 Tidal rectification appears as a likely propagator within the domain of the differences on tidal currents observed at  
523 the open boundaries (Figure 11).

### 524 5.3. *Combining tidal forcing and OGCM forcing: potential impact of imperfectly detided fields?*

525 In this section, we open the discussion to another issue, which is not directly related to the two previous ones but  
526 that remains central in the general problem of downscaling tides.

527 In an attempt to downscale in a regional (child) model both the tidal and non-tidal (i.e. general circulation) dy-  
528 namics, one faces two possible strategies: 1/ chose a parent model that simulates both tides and the general circulation  
529 and use the total parents fields (tidal plus non-tidal), 2/ use separate boundary conditions for the tidal and non-tidal dy-  
530 namics. The first strategy requires that the tidal signal prescribed at the open boundaries and the one generated inside  
531 the domain are in phase which is very likely not to happen. Reasons for this situation not to happen are, as illustrated  
532 in this study and commented in section 5.2, the usual inconsistency between the bathymetry and the difference of  
533 resolution between the parent and child models. Other possible causes include the difference of parameterizations  
534 (for instance for the bottom friction) and other forcings. This is the reason why the coastal modelling communities  
535 (MacCready et al., 2009; Dong et al., 2011; Katavouta and Thompson, 2016) usually consider separately the open-  
536 boundary forcing fields for tides and for the general circulation. This raises however another issue that we address  
537 below.

538 Indeed, in case the second strategy is chosen, the forcing fields for the general circulation (from the parent model)  
 539 must be completely detided. If they are imperfectly detided, the residual tidal signal may be aliased and impact the  
 540 interior solution. In our case, the parent fields from the IBI/NEMO model are provided every 24h as averages over  
 541 25 hours. Averaging over 25 hours is a very efficient way to remove the M2 signal (period of 12.4 hr) which is by  
 542 far the largest constituent in the Bay of Biscay. It also partly removes the other semi-diurnal and diurnal components.  
 543 A thorough look at the effect on S2 tides however shows a significant residual at the MSf frequency (i.e. at the  
 544 fortnightly period of the spring-neap cycle). In the IBI solutions, we find a residual signal of up to 3 cm in the  
 545 northern part of the domain, and approximately 1 cm on average in the whole domain (Figure 13a). For comparison,  
 546 the same tidal analysis has been performed on other daily fields from the MERCATOR-Ocean operational system  
 547 with the NEMO model running without tides: the so-called PSY2V4R4 product with a 1/12° horizontal resolution and  
 548 data assimilation. The analysis results in a much weaker signal at the MSf frequency, of less than 1 cm in the whole  
 549 domain (and likely due to the assimilation of a residual MSf signal in the altimetric data) (Figure 13b).

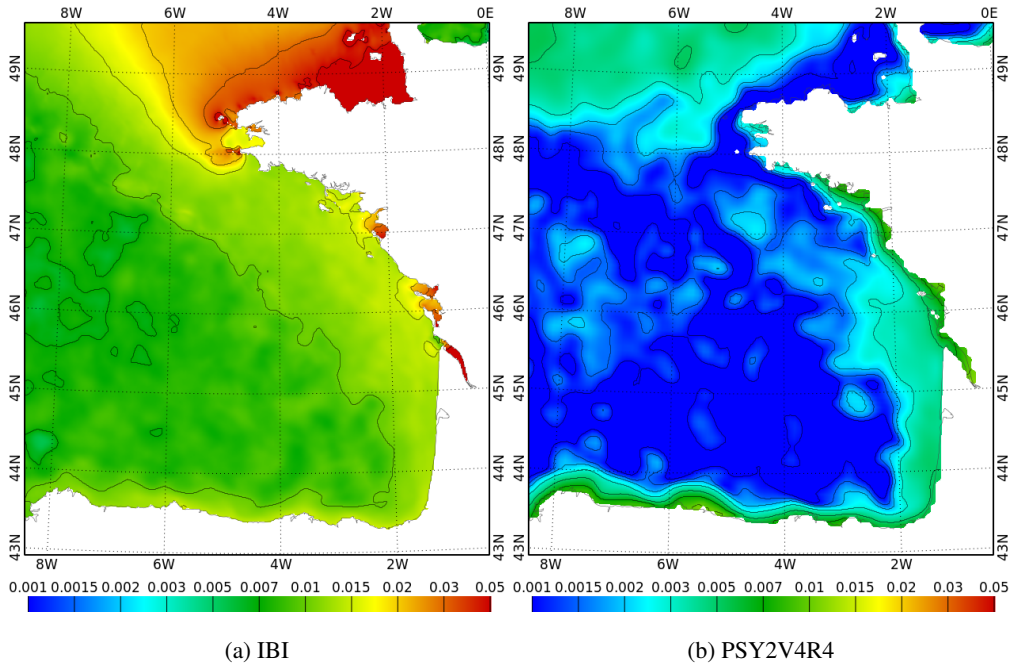


Figure 13: MSf residual (m) after a tidal analysis of daily fields of SSH

550 A 3D simulation forced by PSYV4R4 was performed to compare the results, in terms of tidal elevations, to the  
 551 simulation forced by the IBI operational product. When compared to the tidal gauges (same dataset as the one used  
 552 in section 3.4), the differences in terms of model-data misfits between the two runs are small on average for M2, S2  
 553 and M4 (0.1 to 0.2 cm in complex error). However, differences can be quite high locally, especially in, or close to  
 554 the Gironde estuary for the M4 constituent. At the Royan and Port-Bloc and Richard tidal gauges, the complex errors  
 555 are reduced by 2.2, 1.9 and 4.8 cm respectively. On the other hand, the errors are increased by 6.3, 6.1 and 6.1 cm

556 respectively for the three following upstream stations (Lamena, Trompeloup and Fort-Médoc). Thus, although the  
557 overall error budget remains almost unchanged, the influence of the global forcing can have a strong influence locally.

558 In conclusion, the use of imperfectly detided 3D fields from the parent model leads here to a residual signal in  
559 SSH at the neap-spring frequency all over the domain (1.1 cm on average when forced by IBI, 0.42 cm when forced  
560 by PSY2V4R4 (not shown)), and enhanced over the shelf. The impact on the tidal signal itself is weak on average  
561 except locally. Our comparison to tidal gauges shows that, in spite of this residual signal, there is no evidence of any  
562 degradation on the tides representation in the child model.

## 563 **6. Conclusion**

564 In regional and coastal modelling, a common way to handle open-boundary conditions consists in using external  
565 forcing fields from a model at basin scale in one-way nesting approach. Differences between parent and child models  
566 bathymetry and resolution mainly lead to inconsistencies between the parent forcing and the child dynamics (i.e.  
567 over-specification error as defined for instance by Herzfeld and Gillibrand (2015)). In this study, we propose a robust  
568 and simple approach to improve the downscaling of barotropic tides in a 3D circulation model. We have chosen to  
569 develop an approach that reduces the inconsistencies between the external fields and the interior solution. We do not  
570 try to improve the OBC scheme itself because we aim at proposing a generic approach that can work in many different  
571 configurations with different OBC equations and numerical implementations. By pre-processing the FES2012 tidal  
572 atlas with a 2D simulation (S2D or Tugo), we produce tidal fields generated with the same bathymetry and on the  
573 same grid as the 3D model.

574 Another solution could be to add a few cells in the 2D pre-processing configurations, in order to prevent the  
575 child model (3D) from inheriting any possible errors from the 2D model due to the open-boundary scheme (e.g. rim  
576 currents). This is not the option chosen here as we aimed to avoid defining and handling an extra configuration (the  
577 coastal domain of interest with extra cells). We have checked in the 2D solutions that there was indeed no evidence  
578 of reflection for M2 currents (which are the dominant tidal currents).

579 Results show the potential benefit of using a tailored tidal forcing to force a 3D circulation model. The generation  
580 of a 2D tidal solution on the same grid and bathymetry as the ones used in the 3D model reduces the errors due to  
581 interpolation and bathymetry inconsistencies at the open boundaries, especially on the tidal currents. Compared to  
582 the classical approach, that consists of using a tidal atlas, the use of the tidal model T-UGOm brings about a clear  
583 improvement in the tidal solution of the 3D simulation. This approach and method can be transposed to other 3D  
584 circulation models, particularly in areas where the tides play a key role on coastal dynamics.

585 Thanks to several datasets, the influence of the tidal forcing on the tidal solution and the circulation in a 3D cir-  
586 culation model was studied. The results show the benefit of considering both single-point data (tidal harmonics from  
587 tidal gauges, SSH time series, current velocities and tidal currents from ADCP data) and wide spatial coverage data  
588 (satellite altimetry tidal harmonics) to determine the best forcing. These findings underline the importance of using

589 several datasets and diagnostics to validate a numerical model. Even though this configuration was conceived to study  
590 fine scale processes, the validation of the large scale circulation is crucial. This study particularly emphasizes the ben-  
591 efit of satellite altimetry, which provides regular time series homogeneous in space. Tidal gauges, although valuable,  
592 can exhibit significant differences in tidal elevations within a few kilometers, due to differences in instrumentation or  
593 the time series lengths for example.

594 Moreover, regional circulation model are often designed to reproduce and study small scale dynamics that occur  
595 very close to the coast. In shallow waters, the propagation and the distortion of the tide are strongly influenced  
596 by the topography and by bottom friction. For example, nonlinear interactions occurring between the tide and the  
597 topography result in the generation and/or amplification of overtides such as M4. Bottom friction strongly impacts  
598 the propagation of M2. To reproduce this behavior in numerical models, several tests are often required to tune the  
599 bottom friction. With 3D circulation models, this calibration can take a lot of time, both in terms of CPU and running  
600 time. With T-UGOm, a large number of tests can be performed in a day, because the running time is of the order of  
601 minutes, against hours or days for the 3D simulations. Even if we use the S2D approach presented in this study, the  
602 running time is still significantly higher than with T-UGOm in frequency-mode. Although the friction formulations  
603 are obviously different between models, the bathymetry and grid being exactly the same, a first approximation can  
604 easily be obtained, before adapting this tuning to the circulation model.

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616 tidal elevations have been made using POCViP, developed at LEGOS by D. Allain and F. Lyard.

## 617 **Appendix: Open-boundary conditions (OBC) in SYMPHONIE**

618 The numerical schemes used at the OB in Symphonie are described in Marsaleix et al. (2006) and are summarized  
619 in this appendix.

620 The state variables of Symphonie are: the horizontal ( $u, v$ ) and vertical ( $w$ ) currents, temperature and salinity ( $T, S$ )  
621 and the free surface elevation anomaly ( $\eta$ ) with respect to the state at rest. Other model variables are those used in  
622 the turbulence closure scheme and are not considered here. Barotropic ( $\bar{u}, \bar{v}$ ) and baroclinic ( $u', v'$ ) components of the  
623 current ( $u = \bar{u} + u', v = \bar{v} + v'$ ) are computed separately using the time-splitting technique described by Blumberg  
624 and Mellor (1987). The barotropic velocities are computed as the depth-averaged velocities. The OBC are based on  
625 distinct formulations for the different variables and on the use of external forcing fields along the open boundaries.

### 626 *Barotropic variables*

627 For the barotropic variables, the OBC consist in:

- 628 • C1- a Flather condition applied to the free surface elevation anomaly ( $\eta$ ) and taking into account the external  
629 forcing
- 630 • C2- a radiative condition for the tangential component of the transport
- 631 • C3 - the transport component normal to the boundary is deduced from the continuity equation, from  $\eta$  (C1) and  
632 from the tangential transport (C2)

633 Let us take the example of a domain with a western open boundary. Figure 14 shows the discretization and the  
634 location of the open boundary. Conditions C1-C3 write as follows at time  $t$ :

$$635 \eta_{1,j} = \eta_{1,j}^f + \frac{1}{\sqrt{gH}}(U_{2,j} - U_{2,j}^f) \quad (4)$$

$$636 \left[ \frac{\partial V}{\partial x} \right]_{2,j} = \left[ \frac{\partial V^f}{\partial x} \right]_{2,j} \Rightarrow V_{1,j} - V_{1,j}^f = V_{2,j} - V_{2,j}^f \quad (5)$$

$$637 \frac{\partial \eta}{\partial t} + \frac{\partial U}{\partial x} + \frac{\partial V}{\partial y} = 0 \Rightarrow U_{1,j} = U_{2,j} - \frac{dx}{dy}[V_{1,j+1} - V_{1,j}] + dx \frac{\partial \eta_{1,j}}{\partial t} \quad (6)$$

638 where ( $\eta^f, U^f, V^f$ ) are the external forcing fields, ( $U = (H + \eta)\bar{u}, V = (H + \eta)\bar{v}$ ) are the barotropic transports,  $H$   
639 the bathymetry and  $g$  gravity.

640 Besides, the barotropic velocity is relaxed towards the external forcing within a sponge-layer in order to reduce  
641 possible reflection of the outgoing flow. There is no additional constraint on the global mass conservation. Indeed,  
642 Marsaleix et al. (2006) show that the mean sea surface elevation over the domain tends toward the mean sea surface  
643 elevation of the forcing field over a time scale of about  $S/Lc$ , where  $S$  is the area of the domain,  $L$  the length of  
644 the open boundaries and  $c$  the mean barotropic phase speed. Taking rough values for our domain ( $L = 500km, H =$   
 $5000m, S = LxL, c = \sqrt{gH}$ ) we obtain a time scale of 0.6 hours.

645 The reason for the effectiveness of the Flather condition is discussed in Blayo and Debreu (2005). Besides, we note  
646 that with such a choice of conditions and of implementation on the Arakawa C grid (Figure 14), for any point inside  
647 the domain, the first order (linear) terms in the equation of motions depend on the elevation and tangential velocity  
648 at the boundary only. The normal velocity at the boundary enters the second-order advection and diffusion terms. As

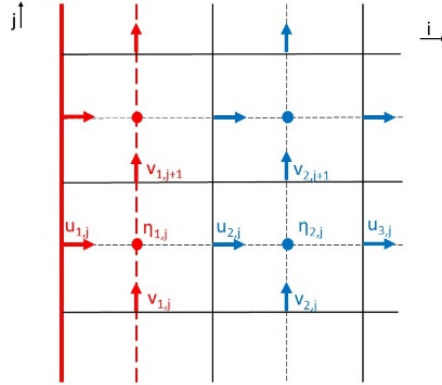


Figure 14: Implementation of the OBC for a western boundary on the SYMPHONIE grid. The boundary points are shown with red symbols, while interior points are shown in blue.

649 noted by Herzfeld and Andrewartha (2012), this allows to minimize the impact of uncertainties in the normal velocity,  
 650 which is critical since the normal velocity is responsible for import/export of mass and energy through the boundary.

#### 651 *Baroclinic variables*

652 For the baroclinic variables, a radiation condition is applied to the perturbations of  $(u', v')$  from the external forcing  
 653 as explained by Marsaleix et al. (2006) (see their equation 28).  $(u', v')$  are also relaxed toward the external fields in a  
 654 sponge-layer. The overall condition writes:

655 C4-  $Bu' = Bu^f$  and  $Bv' = Bv^f$  where the operator  $B$  is given by  $B = \frac{\partial}{\partial t} + c \frac{\partial}{\partial n} - \frac{\div}{\tau}$

656 Condition C4 ensures that the child solution and the external forcing give the same response to the boundary  
 657 operator  $B$ , which is a way of imposing consistency, as recommended by Blayo and Debreu (2005). The speed  $c$  is  
 658 constant and of an order of magnitude comparable to that of the phase speed of the internal waves (here  $c = 1m/s$ ).

659 In the BOBSHELF configuration, the sponge-layer is 30 points wide (about 60km considering the horizontal  
 660 resolution near the open boundaries of the domain) and  $\tau$  decreases from 1 day at the closest point to the open-  
 661 boundary to one hundredth of its value at the 30th point inside the domain.

#### 662 *Temperature and salinity*

663 The temperature and salinity conservation equations are used to compute the boundary conditions in  $T, S$ , as  
 664 justified by Marsaleix et al. (2006). These involve the velocities at the OB as computed from conditions C1-C4. The  
 665 upwind advection scheme adjusts its calculation according to the sign of the current component that is normal to  
 666 the open boundary. In incoming conditions, temperature and salinity open boundary fluxes are calculated using the  
 667 external  $T, S$  fields and, in the opposite case, considering the  $T, S$  fields of the interior solution.

#### 668 *External forcing*

669 The external forcing comes from a tidal model (or tidal atlas) and from a general circulation model (parent model)  
670 for the residual circulation, as explained in section 3.1. To be more precise:

- 671 •  $\eta^f$  is the sum of the elevations for the nine tidal constituents provided by the tidal atlas and of the non-tidal sea  
672 level elevation provided by the parent circulation model ;
- 673 •  $(U^f, V^f)$  are the sum of the horizontal transport for the nine tidal constituents and of the horizontal non-tidal  
674 transport from the parent model ;
- 675 • the external forcing fields for  $(u', v', T, S)$  are given by the parent model.

676 *Dirichlet (clamped) condition used for the 2D SYMPHONIE simulation*

677 In the 2D SYMPHONIE run (S2D) performed to generate the tidal forcing (see section 3.3), the Flather condition  
678 C1 is replaced by a Dirichlet condition that writes:

$$679 \quad \eta = \eta^f$$

680 In this case,  $\eta^f$  is the tidal elevation given by FES2012, since for the 2D run we exclude any other forcing than  
681 tides (no atmospheric forcing field, no river runoff, no residual circulation at the open boundaries). Conditions C2-C4  
682 are unchanged and the barotropic velocity is relaxed towards the external current (here from FES2012) in the sponge  
683 layer as described above; the relaxation is expected to reduce possible reflection in case of outgoing conditions, while  
684 the absence of any other forcing than tides limits the generation of additional waves propagating towards the open  
685 boundary. We have checked that the main tidal current (M2 constituent) does not indeed show any spurious patterns  
686 close to the open boundary.

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