

Improving the Sensitivity of the Plasmon-Based Sensor by Asymmetric Nanoarray

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Abstract

In this work, we investigate the effect of the symmetry in 2D arrays of gold nanoparticle on the sensitivity to the refractive index change. We demonstrate a generalized result that an asymmetric periodic arrangement of metallic nanoparticle leads to a higher sensitivity than a regular square of nanoparticle. Further decreasing the symmetry of the system by using asymmetric nanoparticle (nanorods, triangle) rather than symmetric nanoparticle (nanocylinder) will further improve this sensitivity. Finally, we suggest that such asymmetric nanostructure could operate as a SERS and LSPR plasmon-based sensor by changing the polarization of the incident light.

Keywords Biosensors · Nanoparticles · Optical sensors · Plasmons · Resonance

Introduction

Metallic nanoparticles exhibit very interesting properties in the near-field and the far-field region that are used for sensing applications. When excited by light at proper frequency, a collective oscillation of the electron in the conduction band at a finite interface metal/dielectric arises leading to an electromagnetic mode (solution of Maxwell equations) called localized surface plasmon (LSP). LSP resonance frequency, i.e., frequency at which maximum of electron are driven by the electric field, is extremely sensitive to the nanostructure size, shape and type, and vicinity [1]. Because of this sensitivity, they are of great interest for the development of highly sensitive biosensors for detection of ultrasmall quantity of molecules. Two sensing approaches have emerged in the last two decades: SERS and LSP sensing. Surface-enhanced Raman spectroscopy (SERS) was discovered in the 1970 [2]; it relies on the enhancement of the Raman signal of an active molecule in close vicinity of the nanoparticle by the LSP

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near field. The LSP sensing method employs the wavelength shift of the plasmon resonance to detect change in refractive index due to the deposition of the targeted molecule. Several groups have reported measurements of the sensitivity for isolated nanoparticle with different shapes and sizes [3-9]. In [10], it was demonstrated that for isolated nanoparticle, the sensitivity depends only on the wavelength of the plasmon resonance and the dielectric properties of the metal. Gold nanoparticle arrays are another class of plasmon-based sensor which has attracted the greatest interest in sensor applications. In fact, they are easily fabricated by a variety of techniques enabling mass production. Moreover, they are tunable, versatile, and simple. On a glass substrate, they can be easily functionalized enabling more efficient adsorption of the studied molecules. Gold nanograting is also interesting due to their biocompatibility and wide availability of linked molecules. They are not affected by oxidation; rather, they are stable and can handle high temperatures. They could present a large surface area available for binding and detection. Finally, gold nanograting can sustain localized surface plasmon mode, as well as photonic mode or hybrid mode, leading to an improvement of the quality factor of the resonance peak [11, 12]. Sensitivity improvement of subwavelength nanograting-based SPR sensor has then been studied in several articles [13–23]. It was demonstrated in [23] how the design of an ordered array of gold nanoparticle leads to a tunable plasmon-based sensor with resonance frequency that has any desired value within the visible and near-infrared

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spectrum. Also, in [24], it was demonstrated that coupling the localized surface plasmon resonance to the diffracted wave leads to an improvement of the sensitivity. It was even stated that the quality of the sensor is governed by the frequency difference between the surface lattice resonance and the Rayleigh anomaly of the array. In [25], it was suggested that arrays with unit elements formed by two touching particles display a much higher sensitivity than single-particle arrays. Finally, in [26], it was demonstrated that narrow transmission resonances at near-infrared wavelength arise by breaking the symmetry of coupled split-ring resonators.

None of references cited above investigate the effect of symmetries broken on the sensitivity. In this letter, we comprehensively investigate the effect of the two-grating constant on the sensitivity to the refractive index change in an array of gold nanoparticle. We will demonstrate that by using gold asymmetric array, we could enhance the sensitivity of the plasmon-based sensor, as well as the detected signal strength. We further demonstrate that by reducing the symmetry of the nanoparticle in the array, we can further increase the sensitivity. The polarization of the incident light is an additional influent parameter in the asymmetric array that could tune the plasmon resonance wavelength enabling the proposed sensor to be used as SERS or LSP sensor simultaneously.

Description of the Structure and Numerical Method

We use the FDTD (finite-difference time-domain) technique to solve Maxwell's curl equations in time domain [27]. A home-made code was developed where the computational volume is truncated in the x and y directions to simulate only a unit cell of the grating with the periodic boundaries' conditions. In the z direction, absorbing boundary conditions are implemented to simulate the extension of the lattice to infinity using the CPML (convolutional perfectly matched layer) technique. The CPML technique as compared to the UPML technique (uniaxial perfectly matched layer) has many advantages as it avoids specular reflection error from the front PML interface at low frequency (late-time). Moreover, it allows for more efficient absorption at high angle of incidence and for evanescent waves [28]. This is particularly necessary for our simulations, as the grating periods are large leading to high angle of incidence at the CPML border.

The permittivity of gold is described using the Drude Lorentz model and implemented in the FDTD algorithm using the ADE (auxiliary differential equation) technique [29]. Figure 1 shows the LSP-based sensor under interest consisting of a subwavelength grating structure sitting on a glass substrate. We define the distance p_x to be the edge-to-edge distance between neighboring nanocylinders

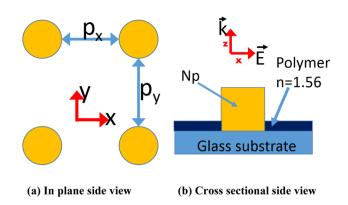


Fig. 1 Nanoparticles array for sensor applications. **a** In plane side view. **b** Cross-sectional side view

along the x direction, and p_y to be similarly defined for nanoparticles along the y direction. The incident light is at normal incidence and is polarized parallel to the x axis. To calculate the shift of the plasmon resonance, and thus to evaluate the sensing ability of the system, we assume that a thin dielectric film (12 nm) with a refractive index n = 1.56 (corresponding to an average value for modeling a layer of biological molecules) is deposited on the glass substrate.

Sensitivity of the Asymmetric Array as Compared to the Square Array

Qualitative Analysis

In the following, influences of the variation of the grating constant along the x and the y directions on the sensitivity will be presented qualitatively. For this goal, we calculate the plasmon resonance shift $\Delta \lambda$ for a nanocylinder with a diameter d = 150 nm, and a height h = 60 nm as well as for a nanocylinder with diameter d = 100 nm and the same height. In Fig. 2, we represent in color the shift $\Delta \lambda$ for different p_x and p_y . It shows that the variation of the distance p_{y} has greater impact on the sensitivity (S is proportional to $\Delta\lambda$) as compared to p_r . The highest sensitivity is observed for an asymmetric array with approximately $(p_y = 2 p_x)$. We explain our result qualitatively by the 2D periodicity of the nanoparticle array. In fact, the geometrical arrangement may reduce for certain distance $p_{\rm v}$ the radiative losses as the energy scattered could be captured into plasmon leading to a higher sensitivity to the refractive index change [30]. In the following, we will focus on the effect of the distance p_{y} on the sensitivity S as well as the effect of reducing the symmetry in each cell in the array.

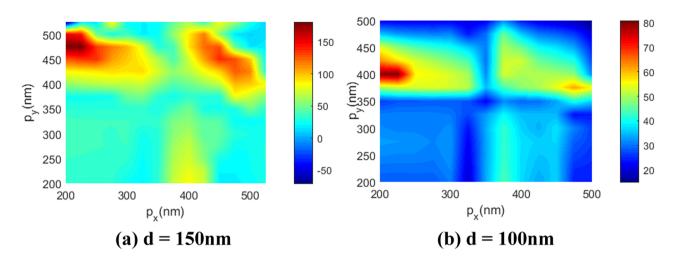


Fig. 2 Effect of the periodicity along the y direction on the sensitivity: The highest shift is observed for the asymmetric array with $p_y > p_x$. **a** d = 150 nm. **b** d = 100 nm

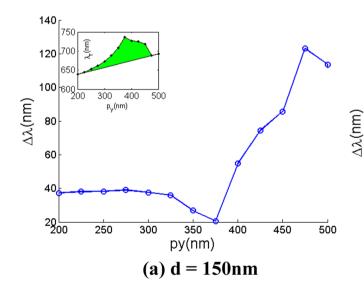
Quantitative Analysis

In this section, we investigate the effect of the variation of p_y (perpendicular to the incident light polarization) on the sensitivity for an array of nanocylinders with diameter d=150 nm.

Results are plotted in Fig. 3a, where we observe a nonmonotonic variation with a little variation of the sensitivity for $p_y < 350$ nm, then a strong increase for $p_y \ge 375$ nm. We attribute the enhancement of S to constructive interference between nanoparticles, especially the diffracted wave that ensure the coupling between proximal nanoparticle in the far-field region. Note that *S* increases especially when the interparticle spacing multiplied roughly by the local refractive index (effective index of substrate, NP, and polymer) approaches the wavelength of the isolated nanoparticle plasmon resonance (we reach threefold enhancement).

Similar results are observed for smaller nanoparticle (Fig. 3b).

Our numerical results show that the variation of the grating constant in the direction perpendicular to the polarization of the incident wave has significant influences on sensitivity. Unlike the isolated particle [10], the increase of the sensitivity here is not correlated to the red shift of the



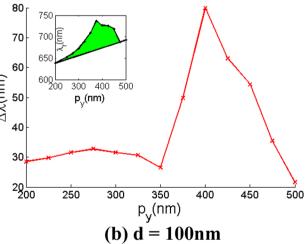


Fig. 3 Effect of the periodicity along the *y* direction on the sensitivity: computations are made for a fixed $p_x = 200$ nm. A significant increase of the sensitivity is observed with an optimal value of $p_y = 200$ nm.

475 nm. The inset, red shift of the plasmon resonance wavelength while increasing p_v . **a** d = 150 nm. **b** d = 100 nm

plasmon resonance wavelength (inset of Fig. 3). The asymmetric array with $p_y > 2$ p_x will exhibit the highest plasmon resonance shift. Our calculations are made for $p_y > 200$ nm; below 200 nm, the far-field coupling is very weak as well as the near –field except for very short distance less than 10 nm where we expect that the near-field becomes dominant; however, such array will be more challenging for fabrication and is out of our interest.

Another interesting aspect of the asymmetric array is that the density of the metal in the nanoarray is greater than the symmetric array leading to a stronger signal detected in the sensor.

In the following, we further reduce the symmetry of the system, by investigating asymmetric nanoparticle in the asymmetric array. For this goal, we investigate the plasmon resonance shift for three types of nanorods as well as for a nanotriangle. The major axis (parallel to the x direction) for each nanorod is fixed to 150 nm. The aspect ratio will be varied from 1 to 3. In order to make a comparison between those nanorods, we normalize the shift $\Delta \lambda$ with the surface area of each nanorod. Results are depicted in Fig. 4a where we observe that the sensitivity could be enhanced by 3 times if we increase the aspect ratio of the nanorod. In Fig. 4b, we further confirm our result by comparing the sensitivity of an array of nanotriangle and an array of nanocylinder. We choose the dimension of the nanotriangle such us it has approximately the same surface area of the nanocylinder, i.e., same contact surface with the target molecule. We observe that with a nanotriangle, we could reach the highest plasmon resonance shift. We attribute the increase of sensitivity to the larger confinement of the electron cloud due to the geometry of the nanoparticle.

In conclusion, we suggest that adding to the asymmetry of the array, the asymmetry of each nanoparticle improves significantly the sensitivity of the nanosystem. The decrease of the symmetry in the nanograting could lead to higher near-field confinement in each nanoparticle as well as efficient far-field coupling between neighboring nanoparticles. The combination of these two factors leads then to the increase of the sensitivity (up to 6 times when comparing to a symmetric array).

Use of the Developed Sensor as SERS and LSPR Sensor

In this section, we investigate the near-field properties and SERS enhancement when increasing the period along the y direction. For this goal, we calculate the average enhancement factor over a small surface area (12 nm) surrounding a nanoparticle with d=150 nm (see the inset of Fig. 5). The excitation wavelength of the laser is assumed to be 785 nm. The SERS EM enhancement factor **G** resulting from electromagnetic mechanism of SERS is defined usually as:

$$G = \frac{\left|E(\omega).E(\omega + \delta\omega)\right|^2}{\left|E_0(\omega).E_0(\omega + \delta\omega)\right|^2}$$

where ω is the incident laser frequency, and $\delta \omega$ is the Stokesshifted frequency; *E* is the total field at the molecule location, and *E*₀ is the incident field. We demonstrate in Fig. 5 that the SERS enhancement could be increased by modifying the period in one direction perpendicular to the polarization of the incident light. Our results seem to be coherent

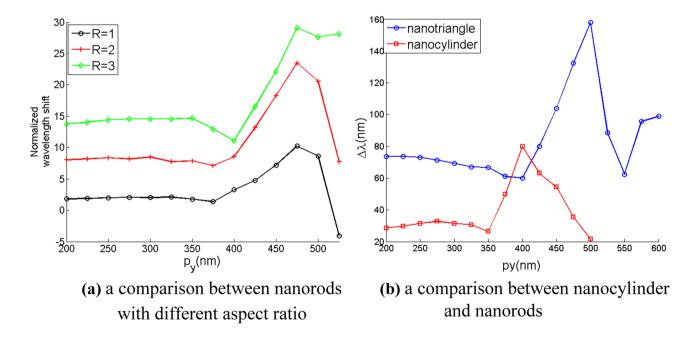


Fig. 4 Calculation of SERS enhancement for different polarizations. **a** A comparison between nanorods. **b** A comparison between nanocylinder with different aspect ratio and nanorods

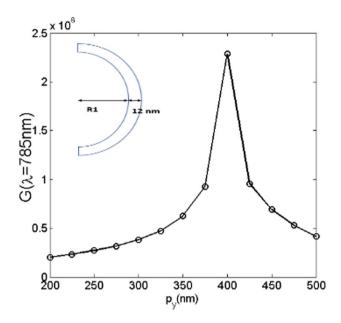
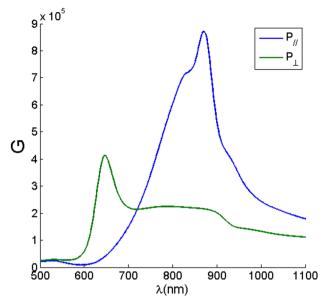


Fig.5 Effect of the symmetry along the y direction on the SERS enhancement. The inset, region surrounding the NP where the enhancement was calculated

with the one found in [31] where the authors demonstrate that higher enhancement factor was obtained for a rectangular array rather than square array of gold slit nanostructure.

One of the characteristics of an asymmetric array is that they are sensitive to the polarization of the incident light. Indeed, the polarization could be an additional degree of freedom that can change the properties and thus the application of the asymmetric nanostructure. To study this effect, we calculate for a grating with dimensions $p_x = 200$ nm and $p_{\rm v} = 500$ nm the enhancement factor when the array is illuminated with a light polarized along the x axis (denoted p//) and when the incident light polarization is along the y axis denoted $p\perp$. Those array dimensions were shown previously to be very suited for LSP sensing (Fig. 2a). If we consider now that the excitation wavelength of the laser used for the SERS sensing is $\lambda_{\text{Laser}} = 633$ nm, we demonstrate in Fig. 6 that the SERS enhancement factor for an excitation wavelength $\lambda_{\text{Laser}} = 633$ nm could be enhanced by 4 orders of magnitude by simply rotating the sample, i.e., by changing the polarization of the incident light leading then to a significant enhancement factor required for SERS sensing. Note that the switching in between polarization orientations results in switching between LSP and SLR mode. Moreover, the resonance wavelength found in the G calculation is red shifted compared to the resonance observed in the far-field. This could be explained among others by the near-field nature of the EM mode present at nanoscale that is greatly affected by the presence of the evanescent waves. Another explanation may result in the fact that we made of the computation of G a spatial average in a very short distance close to the nanoparticle. We expect



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Fig. 6 Effect of the periodicity polarization on the SERS enhancement

that resonance in the near-field could be further shifted if we consider a larger region around the nanoparticle.

Consequently, an array with dimensions $p_x = 200$ nm and $p_y = 500$ nm is very suitable for LSPR sensing at nearinfrared wavelength as well as for SERS applications in the visible range. More generally, we believe that similar nanostructure based on asymmetric array and polarization change could be designed and optimized to operate as SERS and LSPR sensing at different wavelength range. This is of great technological interest as such device could be easily fabricated and optimized.

Conclusion

We have demonstrated in this work that reducing the symmetry of an array of nanoparticle by variation of the period in the direction perpendicular to the incident light polarization or by reducing the symmetry of the nanoparticle leads to a sensitivity of the resonance peak to the change in refractive index that is significantly higher than can be produced by a perfectly symmetric array. Moreover, such asymmetric nanostructure will be sensitive to the polarization of the incident light enabling them to be used as LSPR and SERS sensors. Our result has a general and predictive character that could help to develop simpler and cheaper surface plasmon-based sensor for biodetection and environmental applications.

Author Contribution Aymen Bouali (email: aymen.bouali@ipest.ucarthage.tn Tel: + 21,652,376,349): writing—original draft preparation, reviewing and editing, investigation, methodology, software. Montacer Dridi: data curation, writing—visualization. Florent Colas: conceptualization, methodology, software. Chantal Compere: supervision, validation.

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Data Availability All data are in the manuscript.

Code Availability Yes.

Declarations

Ethics Approval Yes.

Consent to Participate Yes.

Consent for Publication Yes.

Competing Interests The authors declare no competing interests.

References

- 1. Maier SA (2007) Plasmonic Fundamentals and Applications. Springer
- Jeanmaire DL, Van Duyne RP (1977) Surface Raman spectro electrochemistry: Part i. heterocyclic, aromatic, and aliphatic amines adsorbed on the anodized silver electrode. Electroanal Chem 84:1–20
- McFarland AD, Van Duyne RP (2003) Single nanoparticles as real-time optical sensors with zeptomole sensitivity. Nano Lett 3:1057–1062
- Haes AJ, Van Duyne RP (2002) A nanoscale optical biosensor: sensitivity and selectivity of an approach based on the localized surface plasmon resonance spectroscopy of triangular silver nanoparticles. J Am Chem Soc 124:10596–10604
- Haes AJ, Zou S, Schatz GC, Van Duyne RP (2004) Nanoscale optical biosensor: short range distance dependence of the localized surface plasmon resonance of noble metal nanoparticles. J Phys Chem B 108(22):6961–6968
- Duval Malinsky M, Kelly KL, Schatz GC, Van Duyne RP (2001) Nanosphere lithography: effect of substrate on the localized surface plasmon resonance spectrum of silver nanoparticles. J Phys Chem B 105(12):2343–2350
- Raschke G, Brogl S, Susha AS, Rogach AL, Klar TA, Feldmann J (2004) Gold nanoshells improve single nanoparticle molecular sensors. Nano Lett 4:1853–1857
- 8. Jensen TR, Lazarides AA, Kelly KL, Schatz C (2000) Optical properties of metal nanoparticles and nanoparticle aggregates important in biosensors. J Mol Struc 529:52959–52963
- 9. Schatz GC, Yang W, Van Duyne RP (1995) Discrete dipole approximation for calculating extinction and ra- man intensities for small particles with arbitrary shapes. J Chem Phys 869–875
- Miller MM, Lazarides AA (2005) Sensitivity of metal nanoparticle surface plasmon resonance to the di- electric environment. J Phys Chem B 109(46):21556–21565 (PMID: 16853799)
- Auguié B, Barnes WL (2008) Collective resonances in gold nanoparticle arrays. Phys Rev Lett 101:143902

- Chu Y, Schonbrun E, Yang T, Crozier KB (2008) Experimental observation of narrow surface plasmon resonances in gold nanoparticle arrays. Appl Phys Lett 93:181108
- Gwinner MC, Koroknay E, Fu L, Patoka P, Kandulski W, Giersig M, Giessen H (2009) Periodic large-area metallic split-ring resonator metamaterial fabrication based on shadow nanosphere lithography. Small 5(3):400–406
- Tóth E, Sipos Á, Fekete OA, Csete M (2021) Spectral engineering via complex patterns of circular nano-object miniarrays: II. Concave patterns tunable by integrated lithography realized by circularly polarized light. Plasmonics 16:599–617
- Bouali A, Haxha S, AbdelMalek F, Dridi M, Bouchriha H (2014) Tuning of plasmonic nanoparticle and surface enhanced wavelength shifting of a nanosystem sensing using 3-d-fdtd method. IEEE J Quantum Electron 50(8):651–657
- 16. Jensen TR, Duval ML, Kelly KL, Lazarides AA, Schatz GC, Van Duyne RP (1999) Nanosphere lithography: effect of the external dielectric medium on the surface plasmon resonance spectrum of a periodic array of silver nanoparticles. J Phys Chem B 103(45):9846–9853
- Sonnefraud Y, Verellen N, Sobhani H, Vandenbosch GAE, Moshchalkov VV, Van Dorpe P, Nordlander P, Maier SA (2010) Experimental realization of subradiant, super radiant, and fano resonances in ring/disk plasmonic nanocavities. ACS Nano 4(3):1664–1670 (PMID: 20155967)
- Hao F, Sonnefraud Y, Van Dorpe P, Maier SA, Halas NJ, Nordlander P (2008) Symmetry breaking in plasmonic nanocavities: subradiant LSPR sensing and a tunable fano resonance. Nano Lett 8(11):3983– 3988 (PMID: 18831572)
- Henzie J, Lee MH, Odom TW (2007) Multiscale patterning of plasmonic metamaterials. 2:549–554
- Hicks EM, Zhang X, Zou S, Lyandres O, Spears KG, Schatz GC, Van Duyne RP (2005) Plasmonic properties of film over nanowell surfaces fabricated by nanosphere lithography. J Phys Chem B 109(47):22351–22358 (PMID: 16853911)
- Lee S-W, LeeK-S, Ahn J, Lee J-J, Kim M-G, Shin Y-B (2011) Highly sensitive biosensing using arrays of plasmonic au nanodisks realized by nanoimprint lithography. ACS Nano 5(2):897– 904 (PMID: 21222487)
- 22. Charles DE, Aherne D, Gara M, Ledwith DM, Gun'ko YK, Kelly JM, Blau WJ, Brennan-Fournet ME (2010) Versatile solution phase triangular silver nanoplates for highly sensitive plasmon resonance sensing. ACS Nano 4(1):55–64 (PMID: 20030362)
- 23. Lévi G, Krenn JR, Salerno M, Schider G, Lamprecht B, Leitner A, Félidj N, Aubard J, Aussenegg FR (2002) Controlling the optical response of regular arrays of gold particles for surface-enhanced raman scattering. Phys Rev B 15:075419
- Rodriguez SRK, Zhang Y, Crego-Calama M, Brongersma SH, Offermans P, Schaafsma MC, Gomez Rivas J (2011) Universal scaling of the figure of merit of plasmonic sensors. ACS Nano 5:5151–5157
- 25. Quidant R, Enoch S, Badenes G (2004) Optical sensing based on plasmon coupling in nanoparticle arrays. Opt Express 12:4604
- Aydin K, Pryce IM, Atwater HA (2010) Symmetry breaking and strong coupling in planar optical metamaterials. Opt Express 18:13407–13417
- 27. Taflove A, Hagness SC (2000) Computational electrodynamics the finite-difference time-domain method
- Gvozdic BD, Djurdjevic DZ (2017) Performance advantages of CPML over UPML absorbing boundary conditions in fdtd algorithm. J of Elec Eng 68:47–53
- Vial A, Grimault AS, Macias D, Barchiesi D, De La Chapelle ML (2005) Improved analytical fit of gold dispersion: application to the modeling of extinction spectra with a finite- difference timedomain method. Phys Rev B 71:085416

- Li M, Cushing SK, Wu N (2015) Plasmon- enhanced optical sensors: a review. Analyst 140
- Ausman LK, Li S, Schatz GC (2012) Structural effects in the electromagnetic enhancement mechanism of surface-enhanced raman scattering: Dipole reradiation and rectangular symmetry effects for nanoparticle arrays. J Phys Chem C 116(33):17318–17327

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