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Key Points:

- An intense bottom boundary current originating from the Iceland-Faroe Ridge and the Faroe Bank Channel flows along the Icelandic Shelf
- The rough topography and the intensity of the current lead to bottom mixing and sustain a large bottom mixed layer
- Submesoscale structures generated locally participate in the spreading of deep water masses in the Iceland Basin

Supporting Information:

Supporting Information may be found in the online version of this article.

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Structure of the Bottom Boundary Current South of Iceland and Spreading of Deep Waters by Submesoscale Processes

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Abstract The northeastern part of the North Atlantic subpolar gyre is a key passage for the Atlantic Meridional Overturning Circulation upper cell. To this day, the precise pathway and intensity of bottom currents in this area is not clear. In this study, we make use of regional high resolution numerical modeling to suggest that the main bottom current flowing south of Iceland originates from both the Faroe-Banks Channel and the Iceland-Faroe Ridge and then flows along the topographic slope. When flowing over the rough topography, this bottom current generates a 200 m large bottom mixed layer. We further demonstrate that many submesoscale structures are generated at the southernmost tip of the Icelandic shelf, which subsequently spread water masses in the Iceland Basin. These findings have major implication for the understanding of the water masses transport in the North Atlantic, and also for the distribution of benthic species along the Icelandic shelf.

Plain Language Summary Water masses formed in the Arctic Ocean overflow into the North Atlantic at the bottom of the ocean, forming the so-called upper cell of the Atlantic Meridional Overturning Circulation (AMOC). The pathway of the currents carrying these water masses is still under debate due to a lack of observations. In this study, we discuss in details the pathway of these bottom currents in the specific area south of Iceland. We show that a steady current flows along the Icelandic continental shelf, and then divide in smaller structures when reaching the southernmost tip of Iceland. We also show that on its way, the current mixes the bottom layer of the ocean. These findings have major implication in the understanding of heat and carbon transport at depth in this area, which constitute an important response of the climate to anthropogenic forcing.

1. Introduction

The northeastern part of the North Atlantic subpolar gyre is a key part of the Atlantic Meridional Overturning Circulation (AMOC, Buckley & Marshall, 2016). Its so-called “upper cell” ventilates the upper 2 km of the Atlantic Ocean, and it transports heat and carbon from the surface to depths (Kostov et al., 2014; Marshall et al., 2014). It therefore plays a determinant role in the response of the climate to anthropogenic forcing (Drijfhout et al., 2012; Meehl et al., 2014; Winton et al., 2013). The main sources of dense water into the upper cell are overflows from the Nordic Seas (Chafik & Rossby, 2019; Lozier et al., 2019; Tsubouchi et al., 2021). There, intense heat loss in winter transforms the water into colder and denser water masses that subsequently flow southward through gaps in topography (Brakstad, Gebbie, et al., 2023).

While it is the crossroad of this global circulation, the region south of Iceland has been poorly studied in details (see Figure 1a for the location of the places mentioned below). At this place, there is no consensus on the shape and intensity of bottom currents. Studies agree for an overall southwestward flow from the Iceland-Faroe Ridge (IFR) and the Faroe-Bank Channel (FBC) regions toward the Iceland Basin, following the Reykjanes Ridge, see for example, Stow and Holbrook (1984); Bianchi and McCave (1999). When looking at it more precisely, opinions diverge a lot, due to the lack of available data in the area. Investigators sometimes only consider the IFR, the FBC, include an overflow over the Western Valley, or assume a pathway across the deep waters of the Iceland Basin, see for example, Bowles and Jahn (1983); Hansen (1985); Perkins et al. (1998); Hansen and Østerhus (2000, 2007); Beaird et al. (2013); Logemann et al. (2013); Guo et al. (2014); Ullgren et al. (2014); Danialt et al. (2016); Zou et al. (2017); Zhao et al. (2018); Hansen et al. (2018); Petit et al. (2019); Chafik and Rossby (2019); Koman et al. (2022); Brakstad, Gebbie, et al. (2023). Understanding the actual properties of local

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geophysical processes at depth is therefore timely. It will allow to better target future in situ observations aiming at quantifying water mass transport and mixing by the bottom currents, and thus better assess deep storage of anthropogenic-induced tracers.

Beyond this slowly-varying and averaged picture, it has been shown in the past years that small-scale processes have an important role in modulating the global ocean properties. This includes submesoscale balanced currents such as Submesoscale Coherent Vortices (SCVs), Intrathermocline Eddies, or fronts (McWilliams, 2019). These structures have been shown to be key for the global heat budget (Su et al., 2018) and the distribution of marine ecosystems (Lévy et al., 2018) via deep-reaching vertical and horizontal transports (Siegelman et al., 2020; Zhong & Bracco, 2013). Small-scale processes also include fine-scale vertical mixing, induced by deep-reaching currents and internal tides flowing over the topography (Gula et al., 2022; Polzin & McDougall, 2022; Vic et al., 2019). These processes are of major importance to regulate the transport of heat and biogeochemical tracers, and they are suggested to be a good candidate for the closing of the oceanic energy budget (de Lavergne et al., 2022; Ferrari & Wunsch, 2009; Jayne, 2009). The contribution of these submesoscale processes in the south Icelandic dynamics has yet not been studied. However, it is likely that they play an important role in the transport of water masses there. Note that the submesoscale is defined here as the scale at which processes happen on horizontal scales smaller than the average deformation radius (here $\mathcal{O}(20 - 30)$ km (LaCasce & Groeskamp, 2020)), and on vertical scales smaller than the bottom mixed layer (here $\mathcal{O}(100)$ m, see Section 3.2).

In the present paper, we discuss in details the bottom circulation south of Iceland using regional high resolution numerical modeling. In particular we discuss the shape and intensity of the bottom boundary current flowing at $\sim 1,000$ m depth along the Icelandic shelf. This current is the connection between Nordic Seas and the north-eastern part of the North Atlantic subpolar gyre. In the following, mention to the “bottom boundary current” refers to this current. We further show that this latter generates numerous submesoscale features on its path and where it overshoots sharp bathymetric structures. These processes are shown to be of importance for the distribution of water masses in the area. In Section 2 we present the methods used to investigate these processes. In Section 3 we present the analysis of the numerical simulations. In Section 4 we discuss and conclude on our results.

2. Methods

2.1. Numerical Simulation of the North Atlantic

We use outputs of a realistic simulation of the North Atlantic Subpolar Gyre, already used and validated in previous studies, for example, Le Corre, Gula, Smilenova, and Houper (2019); Le Corre, Gula, and Treguier (2019); de Marez & Le Corre (n.d.); Smilenova et al. (n.d.); de Marez et al. (2021); Wang et al. (2022). It is performed using the Coastal and Regional Ocean COmmunity model (CROCO, Shchepetkin & McWilliams, 2005, a version of ROMS model). This model solves the hydrostatic primitive equations using the full equation of state for seawater (Shchepetkin & McWilliams, 2011). The horizontal advection terms for tracers and momentum are discretized with third-order upwind advection schemes (UP3), see for example, Klein et al. (2008) for a further description. This parameterization considers implicit dissipation and it damps dispersive errors.

A one-way nesting approach is used. A first simulation of the whole North Atlantic is implemented with a $\Delta x \sim 6$ km horizontal resolution and 50 topography-following levels, such that mesoscale eddies are reasonably well resolved. It is initialized and forced at boundaries with the SODA dataset (Carton & Giese, 2008). At the surface, the forcing is obtained from the daily ERA-INTERIM dataset (Dee et al., 2011). The bathymetry is constructed from the SRTM30 PLUS dataset (Becker et al., 2009). Then, this simulation is used as boundary forcing and initialization for a second —child— simulation in the Subpolar region, with $\Delta x \sim 2$ km horizontal resolution and 80 topography-following levels. This higher resolution resolves small scale bathymetric features. In particular, it allows an accurate description of the FBC and the IFR.

We make use of this high resolution simulation in the present study, for the period 2002–2009 (after a 2-year spin up). Reference to time averaged quantities over this period are denoted $\langle \cdot \rangle_t$. The general features of the simulation throughout the subpolar North Atlantic gyre dynamics, both in the model output and in gridded observations-based products have been validated by Le Corre, Gula, and Treguier (2019). In our domain of interest, a slight average temperature and salinity offset is seen in the whole water column (constant throughout depth). However, it does not affect the average stratification (see Figures 2c and 2d) which is here the main parameter for the study

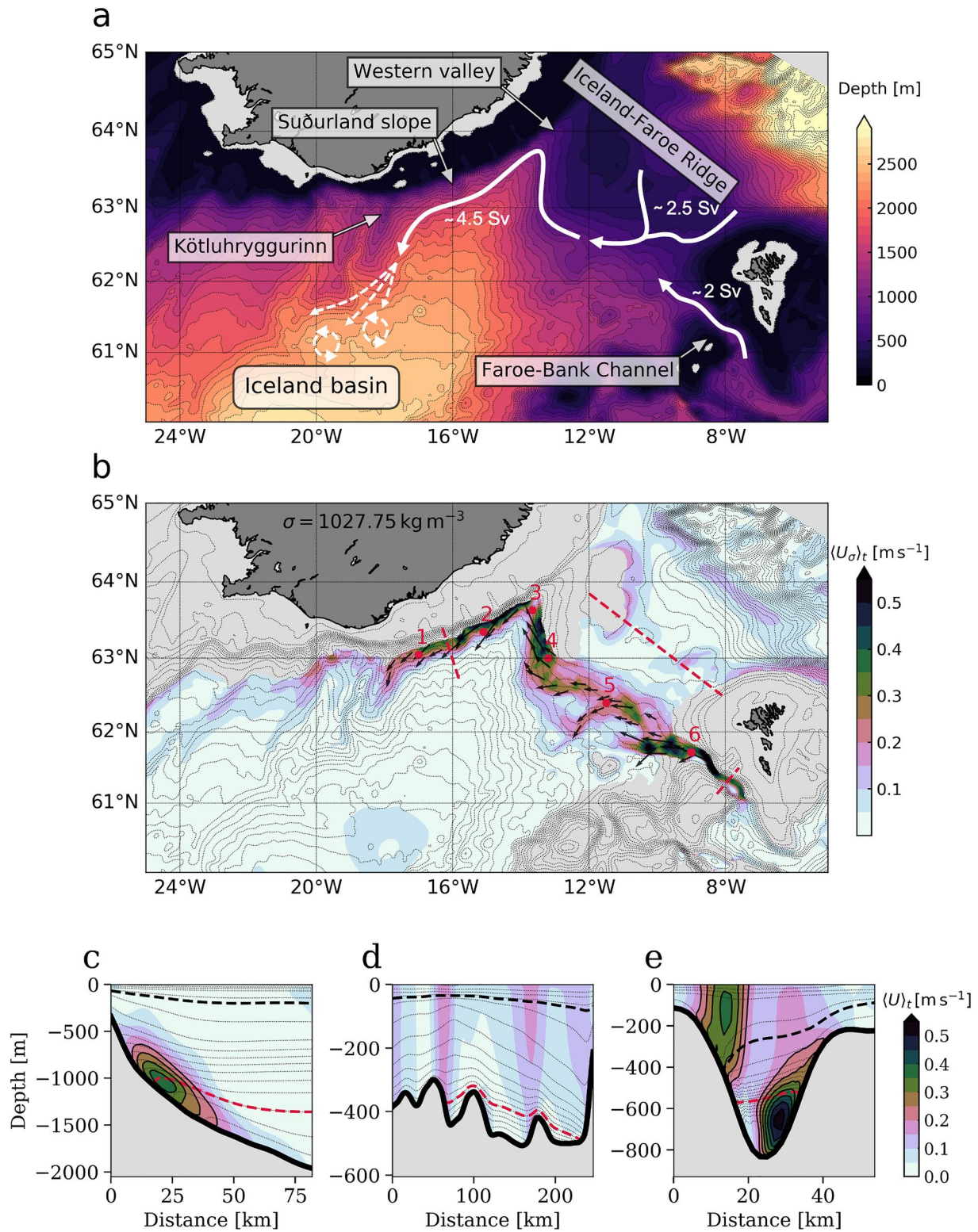


Figure 1. (a) Region of interest, bathymetry, and schematic path of the bottom current; white numbers indicate the transport through the three sections shown in panels (c)–(e). (b) Velocity magnitude on the $1,027.75 \text{ kg m}^{-3}$ isopycnal; position of sections shown in panels (c)–(e), position of profiles shown in Figure 2, and bathymetry (thin black lines). (c)–(e): Vertical sections of the velocity magnitude and isopycnals (thin dashed every 0.05 kg m^{-3} , red dashed $1,027.75 \text{ kg m}^{-3}$, and thick dashed $\sigma_{\text{top}} = 1,027.3 \text{ kg m}^{-3}$).

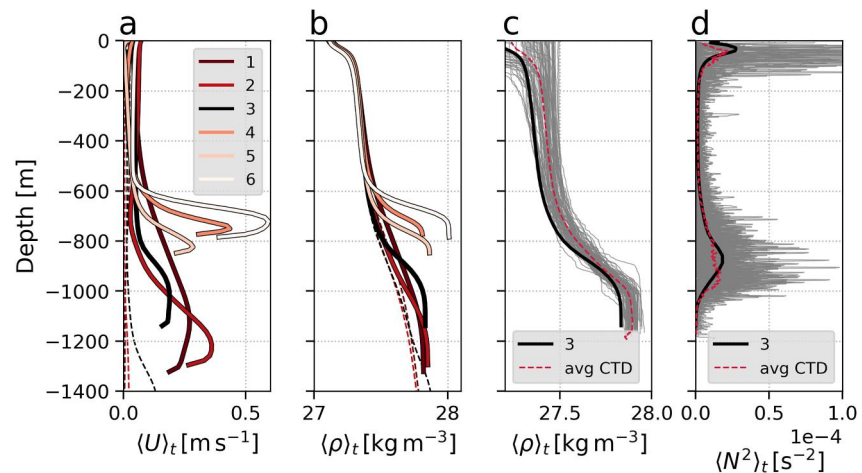


Figure 2. (a) (resp. b): Time-averaged velocity magnitude (resp. potential density) profiles at the locations shown in Figure 1b; thin dashed profiles show profiles ~ 50 km off-shore of the same-color profiles. (c) (resp. d): Comparison of potential density (resp. Brunt–Väisälä frequency) profiles between simulation (thick black) and CTD station (thin gray and thick dashed red) at location 3 (Figure 1).

of the dynamical processes. For further details, we refer the reader to Le Corre, Gula, and Treguier (2019)’s description and validation of the simulation, and their Figure 1 that presents the simulation domain.

2.2. Particule Advection Simulations

We perform three offline particle advection simulations, using the velocity field from the numerical simulation on the $1,027.75 \text{ kg m}^{-3}$ isopycnal, implementing the set of python classes Parcels (Probably A Really Computationally Efficient Lagrangian Simulator). This tool has been widely used in the past few years and it is fully described in Lange & van Sebille (2017), Delandmeter & van Sebille (2019), and in references therein. The three simulations are designed such that they all are one year long. We arbitrarily chose the year 2005 of the CROCO simulation for the currents.

2.3. In Situ Data

The data used for validation and comparison was obtained from SeaDataNet and the Norwegian Marine Data Center (Brakstad, Våge, et al., 2023) for the region southeast of Iceland, corresponding to 80 CTD profiles from 1996 until 2019 covering the 4 seasons. Most of the profiles were uploaded to these open source databases by the Hydrography Observational Program carried out by the Icelandic Marine and Freshwater Research Institute (Ólafsdóttir et al., 2020). The CTD profiles were used to validate the simulation at the virtual location of 13.7°N and 63.6°W (Stokksnes 5), shown in Figure 1b as the point labeled 3.

3. Results

3.1. General Description of the Bottom Current

Time-averaged simulation outputs show that the bottom boundary current originates from two branches at the northeast boundary of the Iceland Basin. A first branch consists of a northwestward flow coming from the FBC. There, an intense current with average maximum velocity of 0.53 m s^{-1} located below 500 m depth flows along the northern slope of the narrow channel, see Figures 1b and 1e. The transport in this channel has been shown in previous studies to be about 2 Sv (Hansen et al., 2016; Hansen & Østerhus, 2007). We determine that this transport is satisfied when integrating the crossing current between the $\sigma_{\text{top}} = 1,027.3 \text{ kg m}^{-3}$ isopycnal and the bottom. A second branch consists of a southwestward flow coming from the IFR. There, two weak currents at $\sim 11^\circ\text{W}$ (on the eastern flank of the westernmost bank) and $\sim 9^\circ\text{W}$ (on the western flank of the easternmost bank) flow over the ridge. The average maximum velocity of 0.19 m s^{-1} at the bottom is seen at the western most location, see Figures 1b and 1d. The crossing overflow transport between σ_{top} and the bottom is about 2.5 Sv, larger than the FBC transport because of the wider section.

When entering the Iceland Basin, the bottom boundary current stabilizes around the $1,027.75 \text{ kg m}^{-3}$ isopycnal, see Figure 1b. It flows northward, constrained along the continental shelf. When reaching the Western Valley, it retroflects following the topography. It then flows southwestward along the continental shelf south of Iceland, namely Suðurland slope, after the name of the Icelandic southern lands. The flow is very well marked along the slope, with average maximum velocity of 0.38 m s^{-1} on the $1,027.75 \text{ kg m}^{-3}$ isopycnal, see Figure 1c. This finding justifies the choice of this particular isopycnal for the further investigation of the current made in this study. The transport induced by the current between σ_{top} and the bottom is about 4.5 Sv , thus satisfying the mass conservation from overflows to the Suðurland slope.

Finally, the current overshoots at a submarine cape located $\sim 18^\circ\text{W}, 62.5^\circ\text{N}$. It is called Kötluhyggurinn, “the Katla ridge”, after the Katla volcano south of Iceland (Shor, 1980). A slight part of the current overflows west over Kötluhyggurinn, creating weak branches of current further west, see Figure 1b. Further examination of the current using particle advection simulations show that these branches have small impact on the water masses trajectory (Section 3.3). Neither seasonal nor inter-annual variability of the bottom boundary current position/intensity/depth are noticed (not shown), thus justifying the use of 7-year overall time averages.

3.2. Vertical Variations and Mixing at the Bottom

Along its path from the FBC to the Suðurland slope, the current velocity magnitude increases with depth to reach a maximum between ~ 0.2 and $\sim 0.6 \text{ m s}^{-1}$, and then decreases toward the bottom floor over an average thickness between $\sim 100 \text{ m}$ and $\sim 500 \text{ m}$, see Figure 2a. It dives from $\sim 700 \text{ m}$ depth at the FBC mouth (profile 6) to $\sim 1,200 \text{ m}$ depth at Kötluhyggurinn (profile 1).

A marked Bottom Mixed Layer (BML) is observed along the current path, see Figures 2b–2d, and is confirmed by 24 years of in situ data. This BML is less than 50 m thick at the FBC mouth. It then becomes thicker along the IFR reaching over 200 m in the Western Valley and along the Suðurland slope. The profile three position coincides with the position of CTD casts performed during a 24 years period in the Western Valley (Stokksness 5, Ólafsdóttir et al., 2020). Average vertical profile of potential density from the simulation matches with in situ observations. The slight offset in density is homogeneous on the vertical and is mainly due to a $\sim 0.5^\circ\text{C}$ temperature offset. Nevertheless, this does not change the dynamics as the stratification (N^2) closely matches thus proving the occurrence of this large BML in the current path, and additionally validating one of the main feature of the simulated current.

The evolution of this BML suggest the combination of frictional and arrested bottom Ekman layers (Brink & Lentz, 2010). The FBC is a narrow-steep-smooth channel which allows the BML to be tightly confined ($\sim 10 \text{ km}$) against the slope; there, the velocity is maximum and the density contrast between the BML and the interior is also the greatest. This bottom boundary current remains confined to the slope throughout the path presented here. First evidence is that this BML is not seen $\sim 50 \text{ km}$ off-shore, outside of the current path, see Figure 2b. Along the path the BML thickness increases coincidentally with the increase in roughness on bottom topography just after the Suðurland slope, which is most likely due to submesoscale viscous processes happening at the bottom, when the current flows over the topography (Polzin et al., 2021).

3.3. The Faith and Spreading of Carried Waters

Two first particle advection simulations confirm that the bottom current originates from both the IFR and the FBC overflows. A total of 6 (resp. 26) particles are released everyday during 300 days along a straight line located in the FBC (resp. on the IFR) on the $1,027.75 \text{ kg m}^{-3}$ isopycnal, see Figure 3a. Remarkably, all particles overflowing in the Iceland Basin eventually get trapped along a very narrow path along the Suðurland slope. We then measure the number of particles from each simulation that cross a section perpendicular to the Suðurland slope, see Figure 3a. Some particles do not reach this region at the end of the simulations (34% and 55%); those particles were either advected too slowly or flowing east of the IFR (see dark blue and dark red trajectories in Figure 3a). Nevertheless, when particles released at both locations get trapped in the bottom current they always travel north toward the Western Valley before retroflecting to the west and crossing the Suðurland slope section. Note that an additional backward advection simulation described in Supplementary Information confirms these findings.

Then, a third simulation is designed in which 15 particles are released everyday during a year along a straight line perpendicular to the Suðurland slope on the $1,027.75 \text{ kg m}^{-3}$ isopycnal, that is, the same section as the one

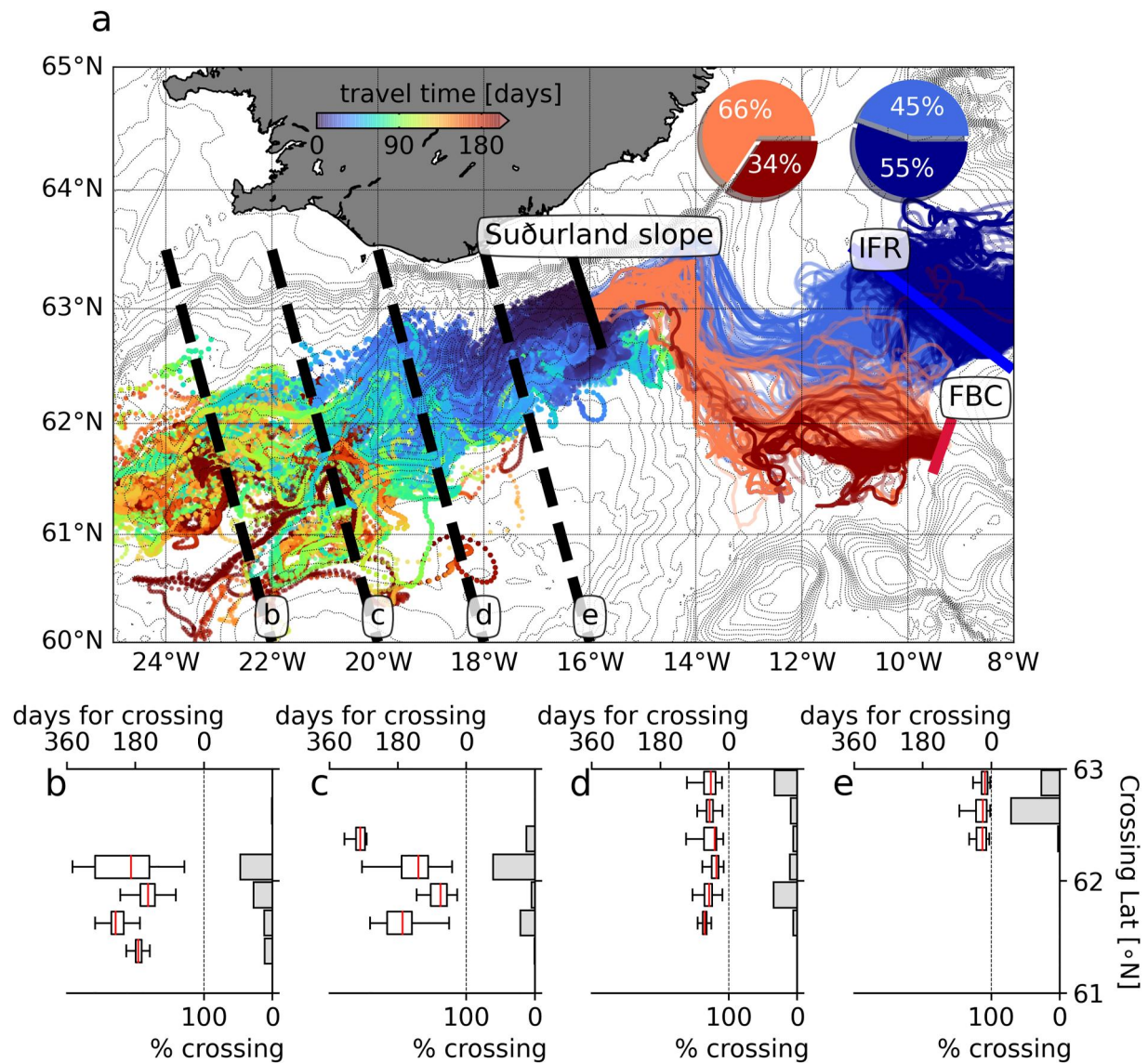


Figure 3. (a) Trajectories of particles (on the $1,027.75 \text{ kg m}^{-3}$ isopycnal) released from the IFR (blue), the FBC (red), and the Suðurland slope (rainbow color that indicates the travel time) sections; darker blue (resp. red) show trajectories of particles released from the IFR (resp. FBC) location that did not cross the Suðurland slope section; for clarity only 1 out of 4 trajectory is shown; pie charts indicate the percentage of trajectories that crossed the Suðurland slope section when released from either the IFR or the FBC locations; black dashed lines indicate the sections used to compute the histograms shown in bottom panels. (b)–(e): Percentage of particles crossing the sections shown in panel (a), and time for the crossing, as a function of latitude.

mentioned previously, see Figure 3a. Particle trajectories from this simulation shows that when reaching Kötluhryggurinn, the waters carried by the bottom current spread out in the Iceland Basin. We measure the latitude and the travel time at which particles cross four different sections, parallel to the launching section, each spaced of 2° in the longitudinal direction, see Figure 3. Particles cross the first section (e) in a few weeks and are concentrated north of 62.5°N , see Figure 3e. After passing Kötluhryggurinn, and as they travel southwestward, they detach from the continental slope, and they cross sections with a large spreading, see Figures 3b–3d. Particles crossing section b are all located south of 62.25°N , and some particles even crossed the 60th parallel North. The spreading is due to turbulent processes, with short time scales, as revealed by the large standard deviations of crossing times. This is also highlighted by the fact that particles are advected by a flow with high values of relative vorticity. In particular, most of the particles have a cyclonic vorticity reaching $\zeta/f > 0.5$ due to the generation of submesoscale structures at Kötluhryggurinn (see Figure S2 in Supporting Information S1). These processes are described in the following section.

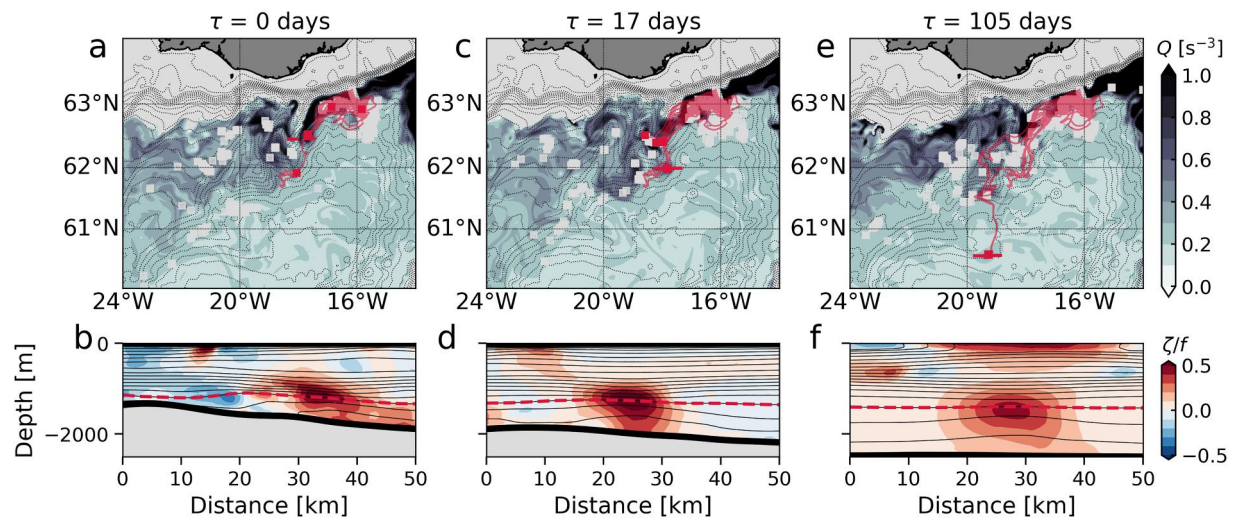


Figure 4. (a), (c) and (e): Snapshots of Ertel Potential Vorticity (divided by 10^9) on the $1,027.75 \text{ kg m}^{-3}$ isopycnal; position of particles trapped (resp. don't trapped) by the SCV at $\tau = 105$ days is shown by the red (resp. white) dots. (b), (d), and (f): vertical section of normalized relative vorticity at the position shown by the red dashed lines in top panels; isopycnals are shown in thin black lines; the $1,027.75 \text{ kg m}^{-3}$ isopycnal is shown by the thick dashed red line.

3.4. Submesoscale Generation at Kötluhryggurinn

Water masses are spread out in the Iceland Basin by submesoscale structures propagating from Kötluhryggurinn. The mechanism is as follows. The bottom current flows along the Suðurland slope, concentrated around the $1,027.75 \text{ kg m}^{-3}$ isopycnal. Viscous interactions (parameterized in the model, see Le Corre, Gula, and Treguier (2019)) with the topography leads to a frictional injection of Potential Vorticity (PV) on this isopycnal. This in turn generates a change of sign of the cross-current PV gradient both horizontally and vertically (see Figure S3 in Supporting Information S1). These are the necessary conditions for Barotropic and Baroclinic instabilities to occur. This results in a highly turbulent flow along the Suðurland slope, as reflected by the high values of Eddy Kinetic Energy (EKE) and Eddy Available Potential Energy (EAPE) on this isopycnal (see Figures S3 and S4 in Supporting Information S1). The flow overshooting at Kötluhryggurinn thus does not follow the slope but meanders south in the Iceland Basin. Water masses are stirred and spread out offshore by intense fronts and rapidly varying flows with —mainly cyclonic— values of vorticity reaching $\zeta/f > 0.5$ (see Figure S2 in Supporting Information S1). Occasionally, the tongue of potential vorticity roll up in the lee of the cape, generating cyclonic SCVs on the $1,027.75 \text{ kg m}^{-3}$ isopycnal. This mechanism has already been intensively studied, and identified as an efficient mechanism to form deep submesoscale vortices in Gradient Wind Balance in many places of the ocean (see *e.g.*, D'Asaro, 1988; Dong et al., 2007; Vic et al., 2015; Gula et al., 2015; Morvan et al., 2019; Srinivasan et al., 2019; McWilliams, 2019; Gula et al., 2019; de Marez et al., 2020). We refer the reader to these studies for further details.

The cyclonic SCVs generated at Kötluhryggurinn enhance the spreading of water masses. A particular event of SCV generation is shown in Figure 4. This structure was generated following the mechanism discussed in the previous paragraph. It then traveled south, hundreds of kilometers, carrying water masses offshore. At $\tau = 105$ days (Figures 4e and 4f), 175 particles (out of 5,464 released in total during the simulation) are trapped in its core and travel southward. This represents more than 3% of the total amount of particles present along the Suðurland slope during a year, that have been spread out by this single event. Counting the number of such events is arduous because most of the time generated SCVs merge between each other making the tracking of single structures hazardous. Nevertheless, we report 15–20 events in the year 2005 of the simulation, suggesting that such spreading by submesoscale events is common in this area.

4. Discussion

In this study, we investigated the bottom boundary current flowing in the north of the Iceland Basin. We showed that it originates from both the Faroe-Bank Channel and the Iceland-Faroe Ridge. It then follows the topography on the $1,027.75 \text{ kg m}^{-3}$ isopycnal where it induces bottom mixing creating a large Bottom Mixed Layer. It finally

overshoots at Kötluhryggurinn, where submesoscale structures are generated and spread water masses in the open Iceland Basin.

In the past decades, circulation in the northern Iceland Basin has been investigated due to its role in the Atlantic Meridional Overturning Circulation, and numerous schematized views of the bottom circulation have emerged. The present paper aims at suggesting that the bottom circulation of the northern Iceland Basin is as schematized as in Figure 1a, with a current coming from both the Faroe-Bank Channel and the Iceland-Faroe Ridge (with about equal contributions) and flowing along the topographic slope. The exact transport induced by each branch of the current is yet to be determined, and the present study does not allow to provide precise estimates. Future deployment of in situ measurements will go in the way of such transport estimates.

Our study also put forth the fact that when overshooting at Kötluhryggurinn, the bottom current somehow disappears and let place to a submesoscale processes-driven spreading of the water masses in the Iceland Basin, thus unvalidating the view of a current steadily flowing along the Reykjanes Ridge. In particular, a significant amount of water is spread out by locally generated cyclonic SCVs. Even if only a few in situ experiments succeeded in measuring SCVs with a sufficient horizontal resolution (see *e.g.*, L'Hégaret et al., 2016; Meunier et al., 2018; Gula et al., 2019, and references therein), only a few observations of *cyclonic* SCVs were reported (*e.g.*, Bosse et al., 2016; de Marez et al., 2020), suggesting that anticyclonic SCVs are predominant in the deep ocean. Our findings thus further suggest that Kötluhryggurinn is an efficient generation spot for deep intense *cyclonic* SCVs. This result is to be confirmed by in situ measurements in the area to allow further analysis of these submesoscale structures.

The region described in this manuscript is of great importance for the future of the AMOC. Indeed, the dense water carried by the bottom current has enormous importance as it significantly contributes to the lower limb of the AMOC. Moreover, the winter convection there can create surface mixed layer depths over 700 m (Brakstad, Gebbie, et al., 2023), which in some regions allows the exchange of surface waters with dense bottom waters. The upper ocean in this region is warming up and IPCC projections suggest this will continue at even higher rates than other basins (Shu et al., 2022). South of Iceland, the combination of deep mixed layers with warmer surface waters, and thick bottom boundary currents with cold-dense waters may exchange this excess of heat resulting in changes of these dense waters in a warming climate.

The bottom boundary current described in this study also appears to be a key phenomenon to sustain biological activity in the area. Indeed, the distribution of several Cold Water Coral species, in particular *Lophelia pertusa*, strongly coincides with the position of the bottom current we described (see Figure 4 of Buhl-Mortensen et al., 2015). It has been shown in the past that the presence of benthic species, such as Cold Water Coral, is strongly correlated to the physical and chemical properties of seawater. In particular, they rely on a renew of suspended food sources and oxygenated waters, that is, feeding currents (Mienis et al., 2019). The bottom current described here has the potential to act as an enhancement-nutrient-supply current. Its strong intensity efficiently renews the bottom water. The interaction of the current with the topography south of Iceland induces strong vertical gradients, locally enhancing vertical mixing of cold nutrient-rich bottom water to the upper layers. The bottom mixing induced by the current also enhances this water flushing, and contributes in increasing the bottom temperature, necessary condition for this species to survive. This current may have implication to a broader spectrum of benthic species, but more investigation in this direction, and a better sampling of physical-biology-related quantities at the bottom is needed to pursue this question.

Finally, even if it is mainly speculations, it is interesting to draw the question of Kötluhryggurinn formation. Studies have discussed the fact that “The Katla Ridges are smooth features with accumulation of sediment beneath the crests in excess of 1.5 km. Their mode of formation is inferred to result from the rapid denudation of Iceland during the Neogene, sediment transport to the base of the slope by turbidity currents and subsequent entrainment and transport southwestward by the flow of Iceland-Scotland Overflow Water.” (Shor, 1980). Even if some other exchanges from the shelf into the canyons may contribute to the sediments, several sources (see *e.g.*, Bowles & Jahn, 1983), suggest that the bottom current has lead to the formation of this bathymetric feature. Taking a step back, this suggests that the bottom current formed Kötluhryggurinn topographic anomaly, which in turn contributed to the generation of submesoscale at this particular place. This could be the signature of geological-timescale forced submesoscale process.

Data Availability Statement

CTD data were provided through SeaDataNet Pan-European infrastructure for ocean and marine data management (<https://www.seadatanet.org>), and can be downloaded as part of the SDC_ARC_DATA_TS_V2 dataset. Due to the large size of simulation outputs, they are available upon request. A script to reproduce particle advection simulations can be obtained online (de Marez, 2023).

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