

1 **Diagnosis of Ocean Near-Surface Horizontal**
2 **Momentum Balance from pre-SWOT altimetric data,**
3 **drifter trajectories, and wind reanalysis**

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8 **Key Points:**

- 9 • Horizontal ocean surface momentum reconstruction from colocated drifters and
10 along-track altimetric data gives closure of up to 80%.
11 • Ageostrophic dynamics account for about one third of the global balanced momen-
12 tum variance.
13 • Errors preventing closure are mainly resolution mismatch errors, followed by colo-
14 cation and instrumental errors.

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15 Abstract

16 Along-track and gridded altimetric observations of sea level are colocated and com-
 17 bined with data from drifter observations and wind reanalysis to reconstruct global in-
 18 stantaneous near-surface horizontal momentum balance. This reconstruction includes
 19 not only geostrophic terms, but also Lagrangian accelerative terms and turbulent stress
 20 terms. The methodology developed quantifies the degree of closure, distinguishes bal-
 21 anced signals from errors, and estimates dynamical compensation between pairs of terms.
 22 Overall, the residual variance of the momentum balance is about 20% of the sum of in-
 23 dividual terms variance. We carry out a detailed exploration of the misclosure, which
 24 is dominated by unbalanced signals in drifter observations (resolution mismatch accounts
 25 for 41% of the total error), followed by instrumental and spatial colocation errors. Al-
 26 though geostrophy is the leading order equilibrium, ageostrophic contributions associ-
 27 ated with non-linear balanced motions, internal tides and near-inertial waves account
 28 for one third of the global balanced signal variance. Momentum balance reconstructions
 29 and the methodology developed here for that purpose hold promise for validating SWOT
 30 sea level observations, for quantifying our ability to estimate the ocean circulation from
 31 these observations, and for improving our understanding of ocean near-surface dynam-
 32 ics.

33 Plain Language Summary

34 Estimates of the ocean surface circulation are routinely gleaned from sea level mea-
 35 surements, based on the assumption of a balance between the Coriolis force felt by any
 36 moving object (including water parcels) traveling at the surface of Earth and the force
 37 induced by spatial variations in sea level. Due to its unprecedented resolution, sea level
 38 observations from the new satellite-based Surface Water Ocean Topography (SWOT) sur-
 39 vey led by NASA and CNES will contain substantial signatures from high-frequency or
 40 shorter-scale motions, where the equilibrium between these two forces is disrupted. These
 41 disruptions hinder our ability to make accurate predictions of ocean surface currents from
 42 sea level measurements. Other sources of observation are required to reconstruct the ocean
 43 surface dynamics more completely. In this paper, we supplement pre-SWOT altimetric
 44 data with surface drifter trajectory data and a wind reanalysis product to diagnose the
 45 ocean surface dynamical equilibrium beyond geostrophy. Our approach indicates that
 46 one third of the global dynamical balance is due to forces that are ignored in the two-
 47 force geostrophic balance. We quantify different sources of error that prevent perfect dy-
 48 namical closure. Our approach will help take full advantage of SWOT sea level measure-
 49 ments.

50 1 Introduction

51 Near-surface ocean variability plays a pivotal role in regulating air-sea interactions
 52 and ocean circulation, which in turn redistribute heat and all ocean-suspended materi-
 53 als. A comprehensive study of this variability is therefore essential for understanding and
 54 forecasting the evolution of the ocean on climatic scales (Ferrari, 2011; Cronin et al., 2019;
 55 Elipot & Wenegrat, 2021). The physical and geochemical environment (e.g. nutrient avail-
 56 ability) is subject to significant control by near-surface dynamics, which in turn exert
 57 a strong influence on the development of marine life, and ultimately affect human ac-
 58 tivities (Taylor & Ferrari, 2011; Lévy et al., 2018). More thorough knowledge and more
 59 precise modeling of surface ocean dynamics can facilitate more accurate operational es-
 60 timations of the ocean surface circulation with numerous potential applications includ-
 61 ing rescue strategies, oil spill containment, forecasting of plastic drift, and enhanced en-
 62 vironmental management (Röhrs et al., 2023).

63 The dynamics of the near-surface ocean is complex, resulting from a combination
 64 of a large range of interacting processes characterized by different temporal and spatial
 65 scales. These processes include mesoscale eddies with $O(100\text{ km})$ horizontal scales, sub-
 66 mesoscale motions with $O(10\text{ km})$ horizontal scales, internal tides and Lagrangian waves
 67 ($100\text{ km}-1\text{ km}$), and quasi three-dimensional turbulence scales ($1-10\text{ m}$). At small Rossby
 68 numbers i.e. large spatial scales and slow motions, near-surface dynamics reduce to the
 69 so-called geostrophic balance between the pressure gradient force and the Coriolis force.
 70 Under this regime, sea level observations provide access to ocean surface currents. How-
 71 ever, for a more accurate reconstruction, the vertical redistribution of momentum induced
 72 by surface winds can be taken into account using Ekman models, for instance (Ekman
 73 and Vagn Walfrid (1905)). In addition, high-frequency and shorter-scale wave-like mo-
 74 tions, such as internal tides and wind-driven near-inertial waves, do not obey geostrophic
 75 balance and require the integration of accelerative and/or advective effects (Yu et al.,
 76 2021).

77 In recent decades, advancements in observational capabilities have enhanced our
 78 understanding of the various processes that contribute to near-surface ocean dynamics.
 79 Up to the 1980s, efforts to reconstruct near-surface ocean dynamics relied on local in situ
 80 measurements (e.g. moorings or oceanographic cruises). These efforts yielded results that
 81 were in good agreement with a Lagrangian-geostrophic-Ekman balance (Davis et al., 1981;
 82 Johnson & Luther, 1994). The addition of satellite-based Global Positioning System (GPS)
 83 for tracking drifting devices in the 1970s and the advent of altimetric satellite missions
 84 in the late 1980s paved the way for investigations of near-surface dynamical balances at
 85 the global scale. Later, in the 2000s, reconstructions of Ekman-geostrophic momentum
 86 balance from these improved observations of sea level, gravity, drifter velocities, and sur-
 87 face winds allowed to assess these observations and to estimate the first global mean dy-
 88 namic topography (MDT), defined as the contribution to spatial sea level fluctuations
 89 that is in balance with the mean surface circulation (Niiler et al., 2003; Rio & Hernan-
 90 dez, 2004).

91 Over the past decade, advancements in altimetric satellite accuracy have enabled
 92 the resolution of features with spatial scales of approximately 65 km in along-track ob-
 93 servations (Dufau et al., 2016; Stammer & Cazenave, 2017). Furthermore, the precision
 94 and resolution of drifter tracking have also improved, resulting in the production of an
 95 hourly low-noise global dataset of surface currents (Elipot et al., 2016). These improved
 96 observations have led to refined MDT products (Maximenko et al., 2009; Rio et al., 2011;
 97 Mulet et al., 2021). The recently launched altimetric satellite-based Surface Water and
 98 Ocean Topography (SWOT) survey is now providing observations of sea level with un-
 99 precedented resolution (estimated 2 km) and coverage owing to its wide-swath capabil-
 100 ity.

101 These successive enhancements in observation capabilities can capture motions that
 102 require going beyond geostrophy. Laying the groundwork for SWOT data, this study aims
 103 to diagnose global ocean surface dynamics by reconstructing instantaneous horizontal
 104 momentum conservation (i.e. estimating and combining the different terms), from along-
 105 track altimetry and colocated drifter and wind observations, thereby capturing high-frequency
 106 and submesoscale contributions. The novelty of this work lies in the use of a set of colo-
 107 cated observations rather than geographically gridded products. We compare the use of
 108 along-track and gridded altimetry. We set out to answer the three following questions:

- 109 1. What is the capacity of pre-SWOT observations to close the upper ocean horizon-
 110 tal momentum budget ?
- 111 2. How do all dynamical terms quantitatively and geographically contribute to the
 112 closure of near-surface ocean dynamics?
- 113 3. What are the sources of misclosure in momentum conservation?

114 The article is organized as follows. Section 2 presents the data and statistical methods
 115 employed to reconstruct and diagnose surface momentum balance from observations. Sec-
 116 tion 3.1 presents the global diagnosis of closure and dynamical contributions. Sources
 117 of misclosure are investigated in Section 3.2. Geographical modulations of these results
 118 are the subject of Section 3.3.

119 2 Materials and methods

120 2.1 Materials

121 Surface momentum conservation was reconstructed over the 2010-2020 period based
 122 on a combination of surface drifter trajectories, altimetric observations and surface wind
 123 stress reanalysis. The hourly dataset from the Global Drifter Program (GDP) (Elipot
 124 et al., 2016) provides observations of surface currents and acceleration. GDP drifters are
 125 tracked using two different positioning systems (i.e. GPS and Argos), which differ in their
 126 accuracy. GDP drifter displacements are assumed representative of water motion at 15
 127 m depth when drifters still possess their drogues. However, a drifter always loses its drogue
 128 after some time, then becoming more influenced by wind drift (Poulain et al., 2009). Sen-
 129 sitivity analyses presented in Appendix A motivate the following choices: 1/ only drogued
 130 drifter data are considered; 2/ a 2.5 cpd low-pass filter is applied to drifter data to min-
 131 imize contamination of acceleration estimates by positional noise; 3/ GPS and Argos drifters
 132 are both considered, because the 2.5 cpd low-pass filter effectively mitigates the greater
 133 noise level in the Argos tracking data.

134 Altimetric data comprises along-track L3 data where we consider sea level anom-
 135 alies (SLA), MDT (see Stammer and Cazenave (2017) for proper definitions), and barotropic
 136 and internal tide corrections for Jason-2, Jason-3, Cryosat-2, Sentinel-3A, and Sentinel-
 137 3B satellite tracks. Along-track SLA effectively resolve processes with wavelengths down
 138 to ~ 65 km in mid-latitudes areas, but have limited spatial coverage inherent to along-
 139 track observations (Dufau et al., 2016). We also consider AVISO-gridded SLA which are
 140 provided daily at a $1/25^\circ$ resolution around the globe. The AVISO SLA are estimated
 141 with optimal interpolation using L3 along-track observations from all altimetric missions
 142 available. Its mean effective spatial resolution is lower than that of along-track obser-
 143 vations, resolving scales of ~ 200 km at mid-latitudes (Ballarotta et al., 2019).

144 The ERA* dataset contains hourly and $1/25^\circ$ wind stress data (Trindade et al.,
 145 2020). This dataset is a data assimilation reanalysis product corrected with geolocated
 146 scatterometers, and presumably takes into account processes that were absent or mis-
 147 represented in the original model (e.g. strong current effects, wind effects associated with
 148 mesoscales, coastal effects, and large-scale circulation effects) (Portabella et al., 2021).

149 2.2 Building a colocated dataset

150 This study is based on colocations between different observations (Figure 1). AVISO
 151 and ERA* are gridded products; both allow interpolation and therefore do not consti-
 152 tute constraints. In contrast, along-track altimetry and drifter trajectory observations
 153 are spatially and temporally sparser. Matchup points where both are available were there-
 154 fore identified with prescribed temporal and spatial mismatch tolerances denoted as ΔT
 155 and ΔX , respectively, for what we call drifter-matchup and altimeter-matchup points.

156 Given the hourly resolution of drifter data, a temporal mismatch ΔT of 30 min was
 157 selected. This temporal mismatch is small compared with the characteristic time scales
 158 of the geophysical signals of interest, namely inertial or tidal periods. Therefore, this mis-
 159 match does not yield substantial temporal colocation errors (Figure S1 in Supporting
 160 Information). A spatial mismatch tolerance of 25 km was chosen. It represents a trade-
 161 off between the number of available colocations and thus the resulting statistical relia-

162 bility and accumulation of colocation errors with spatial mismatch (see Section 3.2.1).
 163 With this spatial mismatch of 25 km, the statistical errors on the momentum residual
 164 mean square were less than 50% of its value in 86% of the ocean 5°-geographical bins
 165 (Figure S2 in Supporting Information).

166 The total number of colocations for these given tolerances is about 239,000.

167 2.3 Reconstructing horizontal along-track momentum conservation

168 We then used this colocation dataset to reconstruct the horizontal surface momen-
 169 tum conservation equation in the altimeter along-track direction (x -axis):

$$\begin{array}{c}
 \text{Lagrangian and Coriolis} \\
 \text{accelerations} \\
 \underbrace{d_t u \quad -fv}_{\text{GDP drifters}} + \underbrace{g \partial_x \eta}_{\substack{\text{Pressure} \\ \text{gradient term}}} - \underbrace{\frac{1}{\rho_0} \partial_z \tau_x}_{\substack{\text{Wind term} \\ \text{Vertical stress divergence}}} = \underbrace{\epsilon}_{\text{Residual}}, \quad (1) \\
 \underbrace{\hspace{10em}}_{\substack{\text{Along-track} \\ \text{+ AVISO altimetry}}} \quad \underbrace{\hspace{10em}}_{\text{ERA}^*}
 \end{array}$$

170 where d_t is the material time derivative, f the Coriolis frequency, u and v the along-track
 171 and cross-track velocities, respectively, g gravity, η introduced as the Dynamic Sea Level
 172 (DSL), ρ_0 the seawater density (considered constant), and τ_x the along-track turbulent
 173 stress. The residual ϵ is composed of several different possible errors, such as missing physics
 174 or estimation process errors, which will be examined in Section 3.2.

175 The surface momentum conservation (1) includes terms involved in the geostrophic
 176 balance, namely Coriolis acceleration and the pressure gradient term, but also two ageostrophic
 177 terms, namely Lagrangian acceleration and vertical turbulent stress divergence. Note that
 178 Eulerian advective terms are accounted for in the Lagrangian acceleration. The quan-
 179 tification of the contributions of these two ageostrophic terms is one of our main focus
 180 here.

181 For each colocation, all terms on the left-hand side of (1) are estimated from ob-
 182 servations as follows:

- 183 1. GDP-filtered velocities are rotated in the along-track/cross-track directions and
 184 provide estimates of the Coriolis acceleration ($-fv$) and Lagrangian acceleration
 185 ($d_t u$) via centered time differentiation.
- 186 2. The pressure gradient term can be estimated in three ways, giving three differ-
 187 ent reconstructions that we compare:
 - 188 (a) along-track reconstruction: the pressure gradient term is estimated from along-
 189 track altimetry at the altimeter-matchup point;
 - 190 (b) altimeter-matchup AVISO reconstruction: the pressure gradient term is esti-
 191 mated from AVISO altimetry interpolated at the altimeter-matchup point, i.e.
 192 with the same spatial mismatch as along-track altimetry;
 - 193 (c) drifter-matchup AVISO reconstruction: the pressure gradient term is estimated
 194 from AVISO altimetry, but interpolated at the drifter-matchup point, thus with
 195 no spatial mismatch.

196 In all cases, DSL η was estimated from the sum of SLA and MDT. Along-track
 197 reconstruction DSL values account for ocean tides and internal tide signals, be-
 198 cause we added the related corrections back in. We were not able to do so for AVISO
 199 reconstructions, because these corrections are not available in the product. DSL
 200 was then differentiated to provide an estimate of the pressure gradient term ($g \partial_x \eta$).
 201 The AVISO DSL spatial gradient in the along-track direction was linearly inter-
 202 polated at the drifter-matchup or altimeter-matchup position and time.

- 203 3. ERA* wind stress was spatially and temporally linearly interpolated at the drifter-
 204 matchup point. This surface wind stress was then extrapolated into vertical stress

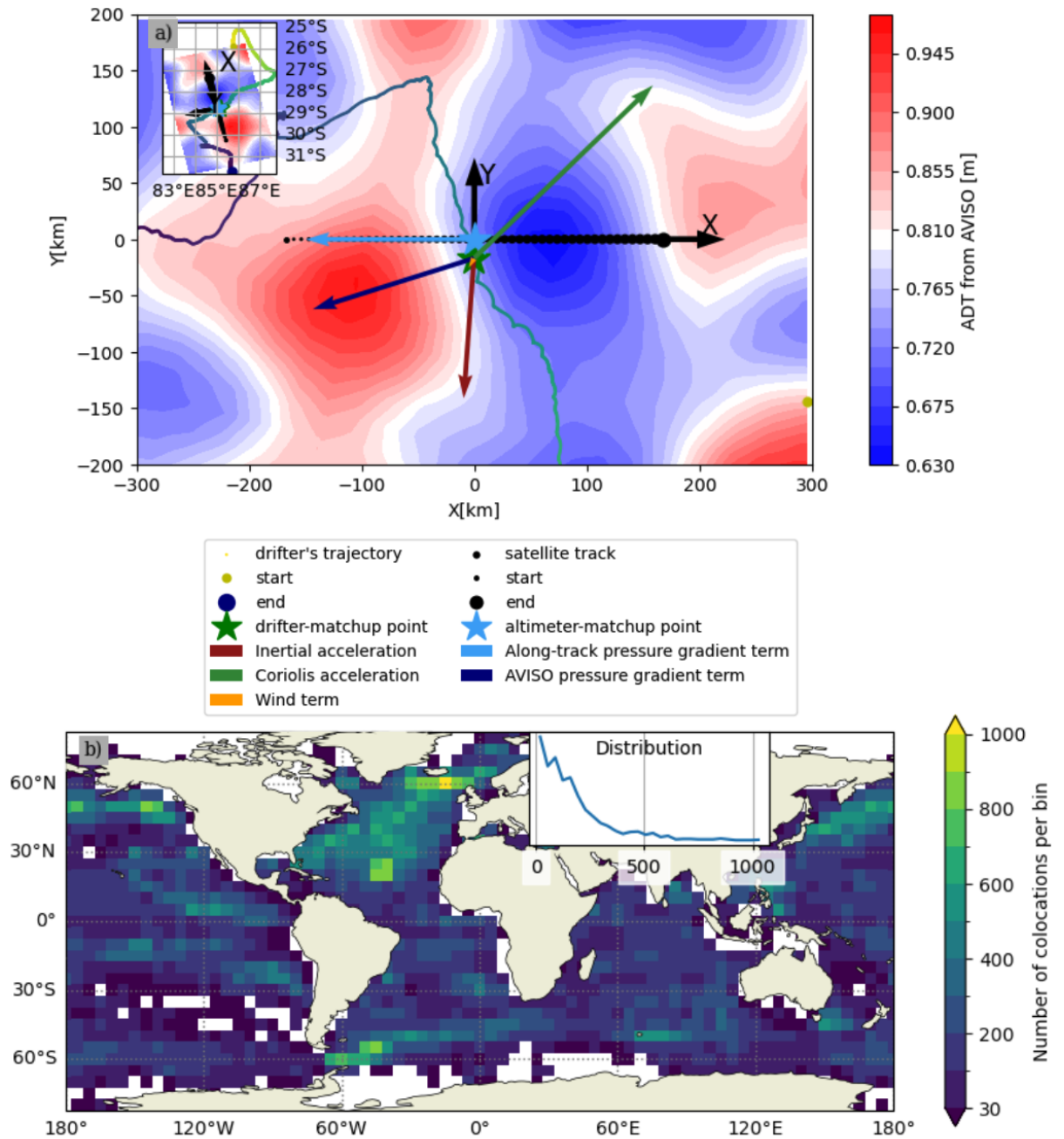


Figure 1. a) Example of a colocation: a drifter trajectory crossing a nadir satellite track in the local coordinates and in the longitude-latitude coordinates (inset). The different terms of the momentum equation (arrows) are estimated either at the drifter-matchup point or at the altimeter-matchup point (stars), which are separated by a ΔX spatial mismatch of 17 km and a ΔT temporal mismatch of 20 min for the represented colocation. The drifter trajectory is represented for 2 months centered around the matchup date. The nadir satellite track is also represented with measurement points on either side of the altimeter-matchup point, measurements being separated by approximately 7 km. The local coordinate system is defined with the along-track x-axis and the cross-track y-axis (black arrows). AVISO absolute dynamic topography (i.e. the sum of MDT and SLA) is interpolated over a 400 km long x 200 km wide x-y oriented box with a spatial resolution of 5 km and used here as a background. Note that, in contrast to the other terms, only the along-track pressure gradient component can be computed and represented here. b) Geographical distribution of colocations in 5°-geographical bins. The inset shows the distribution of the number of colocations per bin.

205 divergence at 15 m using the Rio et al. (2014) empirical Ekman model with their
 206 global 15-m parameters β_{ek} and θ_{ek} ($0.25 \text{ m}^2 \cdot \text{s} \cdot \text{kg}^{-1}$ and $\pm 48.18^\circ$, respectively):

$$\frac{1}{\rho} \partial_z \tau_x = -f \beta_{ek} [\tau_x \sin(\theta_{ek}) + \tau_y \cos(\theta_{ek})] \quad (2)$$

207 **2.4 Statistical methods and diagnostics**

208 All diagnoses in this study will be borne out through the computation of mean square
 209 (MS) values over colocations. The MS of variable x is denoted by its capital letter, e.g.
 210 $\langle x^2 \rangle = X$ (for example $\langle \epsilon^2 \rangle = \mathcal{E}$, $\langle a_i^2 \rangle = A_i$ etc). All terms on the left-hand side of
 211 (1) are denoted as a_i -terms in this section. MSs and variances can be used without distinction
 212 in this analysis, because mean values are smaller than 2% of their standard deviations.
 213

214 Summing the a_i -terms according to equation(1), they will partially balanced each
 215 other out, leaving a residual that contains unbalanced errors. A first way to quantify closure
 216 over the colocation dataset is thus to estimate the MS of this residual :

$$\mathcal{E} = \langle (\sum_i a_i)^2 \rangle. \quad (3)$$

217 While reconstructing the momentum balance, the sum of the MSs of the a_i -terms,
 218 that we denote Σ , is decomposed into two components : a balanced signal component
 219 β , that is closely related to the momentum signal that is effectively canceled in the momentum
 220 conservation reconstruction, and the residual MS \mathcal{E} that contained the errors
 221 preventing a perfect closure. These quantities are related through:

$$\sum_i A_i = \underbrace{\Sigma}_{\text{Total sum of the } a_i\text{-term MSs}} = \underbrace{\beta}_{\text{Balanced component}} + \underbrace{\mathcal{E}}_{\text{Residual MS}}. \quad (4)$$

222 If closure is extremely poor, i.e. if the a_i -terms do not balance out at all, then the
 223 residual MS equals Σ . The sum Σ is thus used to define as second way to quantify closure,
 224 the degree of closure β/Σ , that compares the portion of MS in the balanced signal
 225 component relative to Σ .

226 Then, the balanced signal component can also be decomposed in terms of paired
 227 contributions X_{ij} , that represent the portion of the balanced signal component that is
 228 explained by the equilibrium between two individual terms a_i and a_j :

$$\beta = \sum_{i,j \neq i} X_{i,j} \text{ with } X_{i,j} = -2 \langle a_i a_j \rangle, \quad (5)$$

229 with $\langle a_i a_j \rangle$ the covariance between the a_i -terms a_i , a_j .

230 Finally, the balanced signal component and residual MS are rewritten on a per-term
 231 basis to highlight their origins:

$$\beta = \sum_i \beta_i \quad (6)$$

$$\mathcal{E} = \sum_i \mathcal{E}_i \quad (7)$$

232 with:

$$\begin{cases} \beta_i = \frac{1}{2} [A_i + (\mathcal{E}_{-i} - \mathcal{E})] = - \sum_{j \neq i} \langle a_i a_j \rangle = - \langle a_i \epsilon_{-i} \rangle \\ \mathcal{E}_i = \frac{1}{2} [A_i - (\mathcal{E}_{-i} - \mathcal{E})] = \langle a_i \epsilon \rangle \end{cases} \quad (8)$$

233 where \mathcal{E}_{-i} is the residual MS of a reconstruction without introducing the a_i -term. As
 234 shown by equation(8), the more an a_i -term reduces the residual MS via the difference

235 $\mathcal{E}_{-i} - \mathcal{E}$, the higher its contribution to momentum conservation β_i is and the lower its
 236 residual contribution \mathcal{E}_i is. Alternatively, the balanced signal contribution of a term to
 237 momentum conservation β_i can be understood as the opposite of the sum of its covari-
 238 ances with the other terms and its residual contribution \mathcal{E}_i as its covariance with the resid-
 239 ual.

240 However, the meanings of β_i and \mathcal{E}_i are clear in the particular case in which each
 241 term of the momentum conservation a_i can be decomposed into a balanced physical signal
 242 and an error that is uncorrelated with all other errors or balanced signals, giving

$$\begin{array}{c} a_i \\ \text{Total} \\ \text{signal} \end{array} = \begin{array}{c} b_i \\ \text{Balanced} \\ \text{physical signal} \end{array} + \begin{array}{c} e_i \\ \text{Error} \end{array} \quad (9)$$

243 with

$$\sum_i b_i = 0 \text{ and } \sum_i e_i = \epsilon. \quad (10)$$

244 In this particular case, one can show that contributions β_i and \mathcal{E}_i are equal to the MSs
 245 of the contributions of the balanced physical signals and their errors:

$$\beta_i = B_i = - \sum_{i \neq j} \langle b_i b_j \rangle \text{ and } \mathcal{E}_i = E_i. \quad (11)$$

246 In the more general case, the interpretation of β_i and \mathcal{E}_i as the contributions of balanced
 247 physical signals and their errors is approximate due to residual correlations between er-
 248 ror and physical signals. The accuracy of this terminology for our reconstructions and
 249 the impact of these correlations can be quantified by closely examining the nature of the
 250 different sources of errors that lead to misclosure. We identify these errors as follows:

- 251 1. Resolution mismatch errors: some physical signal in one term, say a_i , remains un-
 252 balanced by other terms because of their lower spatial and/or temporal resolutions.
 253 This unbalanced variance will be contained in the residual contribution \mathcal{E}_i asso-
 254 ciated with a_i , i.e. the most highly resolved term.
- 255 2. Missing physics: some physical signal in one term, say a_i , remains unbalanced be-
 256 cause the term that should provide the balancing signal is simply absent from the
 257 reconstruction (Equation 1). In the present study, such errors can be related to
 258 baroclinic pressure gradients, vertical advection or horizontal dissipation.
- 259 3. Instrumental errors: typically due to noise on drifter position or on altimetric sea
 260 level measurements.
- 261 4. Physical modeling errors: for example resulting from inaccurate representation of
 262 the vertical stress divergence (Ekman dynamics and its parametrization).
- 263 5. Colocation errors: terms in Equation 1 are not estimated at the same position and
 264 time, but with spatial and temporal mismatches of up to 25 km and 30 min, re-
 265 spectively.

266 Resolution mismatch and missing physics errors are not correlated to other terms'
 267 physical signal or error components, because they by definition represent physical sig-
 268 nals that are unbalanced by the other terms. Instrumental errors such as noise on alti-
 269 metric or drifter measurements are reasonably approximated as uncorrelated (Spydell
 270 et al., 2019; Dufau et al., 2016) in which case they are not correlated to both physical
 271 signal or other errors. Lagrangian and Coriolis accelerations come of course from the same
 272 drifter measurements, but also come from orthogonal directions, and related instrumen-
 273 tal errors can thus be considered as uncorrelated as well. Physical modeling errors are
 274 suspected for the wind term. To investigate its importance we quantify the effect of a
 275 scaling error on the wind term by a factor α (see Text 2 of Supporting Information). Scal-
 276 ing factors of 0.5 and 1.5 on the wind term increases the residual MS only by about 1%
 277 and 4% of its value, respectively. This is an indication that modeling errors on the wind

278 term have a limited impact on momentum balance reconstructions. The four previous
 279 error sources are thus expected to produce balanced signal contributions that are of phys-
 280 ical origin. This is not the case for colocation errors (Text 2 of Supporting Information).
 281 These errors indeed affect both the residual and balanced contributions with mirrored
 282 effects : the residual contribution of a term increases as much as its balanced signal con-
 283 tribution decreases compared with the no-colocation-error case (see Supporting Infor-
 284 mation Text 2 for the complete description). These colocation errors can be estimated
 285 through their dependency to the colocation mismatch as done in Section 3.2.1. They rep-
 286 resent a consequent part of the residual contributions and so will be taken into account
 287 in the error budget of Section 3.2.1 along with the other errors. However, they are found
 288 to be small compared with the corresponding balanced signal contributions (represent-
 289 ing at most 7% of the balanced signal contribution for the pressure gradient term). The
 290 balanced signal contributions is thus quasi entirely of physical origin. This finally jus-
 291 tifies the physical meaning we accord to these balanced signal contributions.

292 For clarity, we introduce the unit γ , which is equal to the acceleration related to
 293 a dynamic sea level gradient of 1 mm per km. Thus, γ equals $9.81 \times 10^{-6} \text{ m.s}^{-2}$. For
 294 global results, the 95% confidence intervals on the different metrics are computed assum-
 295 ing Gaussianity for the distribution of the means and using the central limit theorem.
 296 For binned results, a bootstrap method was applied with 9999 resamplings (Efron & Tib-
 297 shirani, 1994).

298 2.5 Global vs. geographical analyses

299 The metrics introduced in Section 2.4 were first the subject of a global computa-
 300 tion as presented in Section 3.1. Then, to highlight regional variability, the same met-
 301 rics were also computed in 5° -geographical bins. Binned estimations are subject to more
 302 substantial relative statistical error for bins with lower numbers of colocations (Figure
 303 S2 in Supporting Information). Geographical bins in which these relative errors are higher
 304 than 50% are not shown. The mean number of colocations per 5° -bin was about 160 (dis-
 305 tribution plotted in the inset of Figure 1).

306 3 Results and discussion

307 3.1 Global average diagnoses

308 The pressure gradient and Coriolis acceleration terms are largest for global aver-
 309 ages, representing together 86% of the sum of MSs (Figure 2). This is expected from the
 310 dominance of sub-inertial motions on ocean surface variability (Yu, Garabato, et al., 2019;
 311 Arbic et al., 2022) and the validity of geostrophy for this class of motions. The Lagrangian
 312 acceleration and wind terms are weaker, representing 12% and 2% of the sum of MSs,
 313 respectively.

314 The closure of momentum balances reconstructed with along-track sea level data
 315 is measured by the global-scale average residual MS, which is $(1.64 \pm 0.03) \gamma^2$. This value
 316 corresponds to a degree of closure of about $(80.5 \pm 0.4)\%$. This high degree of closure
 317 translates into a large majority of individual MSs being balanced (see balanced signal
 318 vs. residual contributions in Figure 2a second bar chart).

319 When considering the drifter-matchup AVISO reconstruction, the global average
 320 residual MS was $(1.27 \pm 0.02) \gamma^2$, with an associated closure degree of $(81.9 \pm 0.4)\%$ (Fig-
 321 ure 2b). The residual MS is thus reduced by about 20% with AVISO data. Using AVISO
 322 data rather than along-track altimetry has two counter balancing effects on the resid-
 323 ual MS. On one hand, the smoothing of altimetric information in AVISO data prevents
 324 from explaining small scale variability contained in the terms estimated from drifter data.
 325 Resolution mismatch errors are introduced in other words. This is shown by the smaller

326 pressure gradient balanced signal contribution with AVISO data (2.08 ± 0.04) γ^2 com-
 327 pared with (2.57 ± 0.04) γ^2 with the along-track sea level data (Figure 2 second line).
 328 On the other hand, the residual MS is reduced by the combination of the same smooth-
 329 ing which partially deletes pressure gradient term instrumental errors and the estima-
 330 tion of the pressure gradient at the drifter-matchup which deletes colocation errors. This
 331 is demonstrated by the quasi absence of residual contribution in the pressure gradient
 332 term, unlike that of the along-track sea level data (9%). The improved performance of
 333 drifter-matchup AVISO reconstruction in terms of residual MS suggests that AVISO map-
 334 ping deletes more instrumental errors and colocation errors than it introduces resolution
 335 mismatch errors. Interestingly, the larger amount of balanced information in along-track
 336 sea level data compared to AVISO helps increase the contribution of the balanced com-
 337 ponent in drifter observations: the Lagrangian acceleration balanced signal contribution
 338 is (0.83 ± 0.03) γ^2 compared with (0.62 ± 0.02) γ^2 for AVISO data and the Coriolis
 339 acceleration balanced signal contribution is (3.20 ± 0.05) γ^2 compared with (2.91 ± 0.04) γ^2
 340 using AVISO data (Figure 2a and b second lines). The processes resolved by along-track
 341 and drifter trajectories explaining these differences are likely high-frequency and/or sub-
 342 mesoscale features that have been filtered out by AVISO mapping process. This result
 343 highlights the additional value of along-track altimetry in accounting for fine scales and
 344 providing more complete momentum balance reconstructions.

345 Using along-track reconstruction, the global dominant dynamical pair is, as expected,
 346 the geostrophic pair (contribution of (4.77 ± 0.09) γ^2 so 57% of Σ) (Figure 2a). The sec-
 347 ond and third most important pairs are the Lagrangian and Coriolis accelerations pair
 348 (1.19 ± 0.03) γ^2 or 14% of Σ) and the Lagrangian acceleration and pressure gradient pair
 349 (0.49 ± 0.03) γ^2 or 6% of Σ). Note that three-term balances such as cyclogeostrophy or
 350 those involved with gravity waves likely contribute to the diagnoses of these pairs, but
 351 three-term balances are beyond the scope of the present analysis. The Coriolis accel-
 352 eration and wind term paired contribution, i.e. Ekman-like dynamics, accounts for ($0.43 \pm$
 353 0.01) γ^2 or 5% of Σ . Contributions from the last two dynamical pairs, the Lagrangian
 354 acceleration-wind and pressure gradient-wind term pairs, are of smaller amplitude and
 355 negative. These negative values may indicate that the two terms in the pairs tend to in-
 356 crease or decrease together and thus must balance out together with a third term. Us-
 357 ing the drifter-matchup AVISO reconstruction instead of the along-track reconstruction
 358 reduces the geostrophic contribution (4.2 ± 0.07) γ^2 and the contribution of the Lagrangian
 359 acceleration-pressure gradient term pair (0.07 ± 0.02) γ^2 , highlighting once again that
 360 AVISO cannot render small-scale variability. Regardless of the altimetric data (e.g. along-
 361 track or AVISO), about 28-30% of the total balanced signal variance is explained by ageostrophic
 362 motions, which can therefore not be ignored for accurate reconstruction of surface mo-
 363 mentum balance.

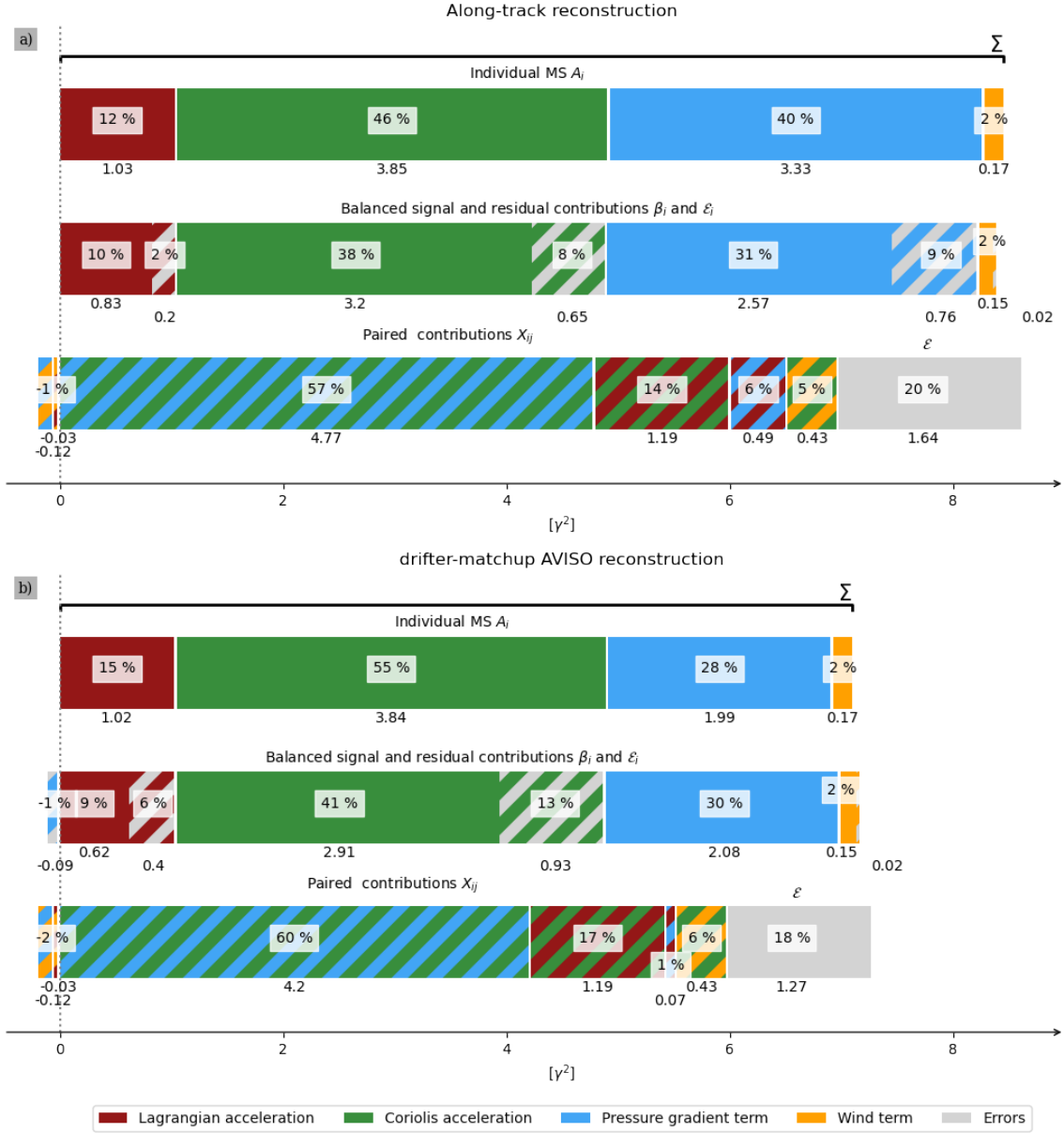


Figure 2. Dashboards of momentum closure decompositions using along-track reconstruction (a) and drifter-matchup AVISO reconstruction (b) (see Section 2.4, Equations 4, 5-7). The top horizontal bar charts show the MSs from each term. Middle bar charts are the decomposition into balanced signal and residual contributions for each term. Bottom bar charts show the decomposition into paired contributions (pairs are indicated by the two-color hatching) and the final residual MS. Each bar is annotated with the corresponding percentage of Σ (total top bar chart) and the corresponding MS expressed in γ^2 . Negative quantities are moved towards negative abscissas.

3.2 Exploring misclosure

We now attempt to characterize and quantify the importance of the different sources of errors enumerated in Section 2.4 that lead to the observed misclosure.

3.2.1 Colocation errors

The temporal mismatch criterion $\Delta T < 30$ min is small compared with characteristic timescales of dominant surface processes (typically ≤ 12 h) and temporal mismatches are thus not expected to lead to substantial temporal colocation errors. This assumption is confirmed by the lack of dependency of momentum closure on temporal mismatches (Figure S1 in Supplementary Information).

We thus surmise that colocation errors are dominated by spatial mismatches. The sensitivity of closure to spatial mismatch is provided by averaged residuals conditioned by spatial mismatch in bins of 2 km width. For this section and this section only, colocations with spatial mismatches up to 200 km were considered. The mean number of colocations per 2 km bin was 19,268, with the distribution being almost uniform.

Along-track and altimeter-matchup AVISO momentum residual MSs steadily increase as a function of spatial mismatch, reaching values 3.6 and 2.7 times their respective minimum values, for mismatches of 100 km (Figure 3a). This increase slows for larger mismatches, indicating that the spatial scales of energetic motions that mainly contribute to momentum balance here are of around 100 km or more (e.g. mesoscales). Ultimately, the pressure gradient is expected to become fully uncorrelated with other terms and the residual MS should converge to $\mathcal{E}_{-g\partial_x\eta} + A_{g\partial_x\eta}$ (max. lines on Figure 3a). This dissociation does not appear to have occurred at scales under 200 km. The increase of residual MS with spatial mismatch translates into a transfer from balanced signal component to residual contribution for all terms except the wind term (Figure 3c and d).

The difference in residual MSs between along-track and altimeter-matchup AVISO increases with spatial mismatch and plateaus at about 100 km (visible on the difference of residual MSs, not shown) to a value that is comparable to the difference in pressure gradient MSs (Figure 3a). The faster increase of the along-track residual reflects the signature of processes that are present in along-track data, but not in AVISO data, presumably due to their finer spatial scales. This pattern is also apparent for the Coriolis-pressure gradient pair contribution from the altimeter-matchup AVISO reconstruction, which becomes comparable to the along-track reconstruction at scales of about 100 km also (Figure 3b).

The drifter-matchup AVISO momentum residual also increases with spatial mismatch, but to much lesser extent (reaching only 1.1 of its value at minimal mismatch at 100 km). The AVISO pressure gradient was interpolated at the exact drifter position and time; therefore, there are no colocation errors, strictly speaking. These errors are rather related to AVISO mapping process, which builds a gridded product from the optimal interpolation of distant along-track altimetric data. These errors are negligible compared to others and thus will not be further discussed.

To estimate the amplitude of colocation errors on the global average residual MSs, the residual MS for spatial mismatch below 4 km (red lines on Figure 3a inset axes) is subtracted from the global [0, 25km] averaged residual MS presented in Section 3.1. The resulting along-track total colocation error is about $0.35 \gamma^2$, i.e. 21% of the corresponding global residual MS.

The sensitivity of the residual MS to spatial mismatch (Figure 3a) can be decomposed as the sum of individual residual contributions sensitivities (Figure 3b). This decomposition is dominated by the Coriolis and pressure gradient residual contributions, whose sensitivities therefore closely resemble residual sensitivities. The Lagrangian term

Table 1. Error budget for along-track reconstruction: MSs, balanced signal and residual contributions and errors for both individual and total values

	Lagrangian acceleration	Coriolis acceleration	Pressure gradient	Wind term	Total
MS	1.03	3.85	3.33	0.17	$\Sigma = 8.38$
Balanced signal contribution $[\gamma^2]$	0.83	3.20	2.57	0.15	6.75
Residual contribution $[\gamma^2]$	0.20	0.65	0.76	0.02	$\mathcal{E} = 1.64$
<hr/>					
Colocation error					
$[\gamma^2]$	0.02	0.14	0.19	negligible	0.35
[% of residual]	1.2%	4.9%	12%	negligible	21%
<hr/>					
Resolution mismatch error					
$[\gamma^2]$	0.18	0.51	-	-	0.69
[% of residual]	11%	31%	-	-	42%
<hr/>					
Instrumental error					
$[\gamma^2]$	negligible	negligible	0.57	-	0.57
[% of residual]	negligible	negligible	36%	-	36%

413 error exhibits some sensitivity for the along-track sea level reconstruction, but no clear
414 sensitivity for the altimeter-matchup AVISO reconstruction. This is consistent with the
415 weak balance between Lagrangian and pressure gradient terms in the AVISO reconstruc-
416 tion (Section 3.1, Figure 2, Figure 3b). The wind term error did not show any clear sensi-
417 tivity to spatial mismatch over the values considered (≤ 200 km). This insensitivity
418 likely stems from the negligible correlation of the wind term with the pressure gradient
419 term (Section 3.1).

420 Applying the same method as for the estimation of total colocation error, individ-
421 ual colocation errors were estimated by subtracting the residual contributions for spa-
422 tial mismatch below 4 km from $[0, 25\text{km}]$ averaged values (Table 1). These colocation
423 errors echo the discussion above on residual contribution sensitivities on spatial mismatch
424 and will be useful for estimating resolution mismatch errors as well as instrumental er-
425 rors.

426 The contributions of the Lagrangian-Coriolis accelerations and Ekman (Coriolis acceleration-
427 wind) pairs to momentum balance closure are insensitive to spatial mismatch, as expected
428 because both terms are estimated at the drifter-matchup point. The geostrophic pair dom-
429 inates by a factor of 4 over the Lagrangian-Coriolis pair at small spatial mismatches for
430 along-track reconstruction, consistently with the individual term decomposition (Fig-
431 ure 3 b). The geostrophic pair decreases with spatial mismatch and contributes less than
432 the Lagrangian-Coriolis accelerations pair for spatial mismatches larger than 150 km. An
433 exponential fit on this decrease leads to decay length scales of about 91 ± 5 km for along-
434 track sea level reconstruction and 119 ± 9 km for altimeter-matchup AVISO reconstruc-
435 tion. This finding emphasizes the dominance of mesoscale motions on momentum bal-
436 ance closure. The difference in decay scale, the larger balanced signal contributions at
437 minimal spatial mismatches with along-track data, and the merging of along-track and
438 AVISO contributions at lags smaller than about 100 km all suggest that finer scale vari-
439 ability is captured with along-track data as argued above. The Lagrangian acceleration-
440 pressure gradient pair contributes at most to about half of the Lagrangian-Coriolis pair
441 for along-track data and smallest spatial mismatch. An exponential fit on the Lagrangian
442 acceleration-pressure gradient pair sensitivity gives a decay length scale of 43 ± 4 km,
443 indicating a reduced spatial scale of the processes contributing to the momentum bal-
444 ance via the Lagrangian acceleration-pressure gradient balance compared with the geostrophic
445 balance. With altimeter-matchup AVISO data, the Lagrangian acceleration-pressure gra-
446 dient pair contributes much more weakly to momentum balance (factor of ~ 6 compared
447 with along-track data), but shows some decay at a comparable length scale. At spatial
448 mismatch larger than about 100 km, along-track and altimeter-matchup AVISO data lead

449 to comparable balanced signal contributions (see also Figure 3d), and AVISO data may
 450 be more useful given its lower noise level.

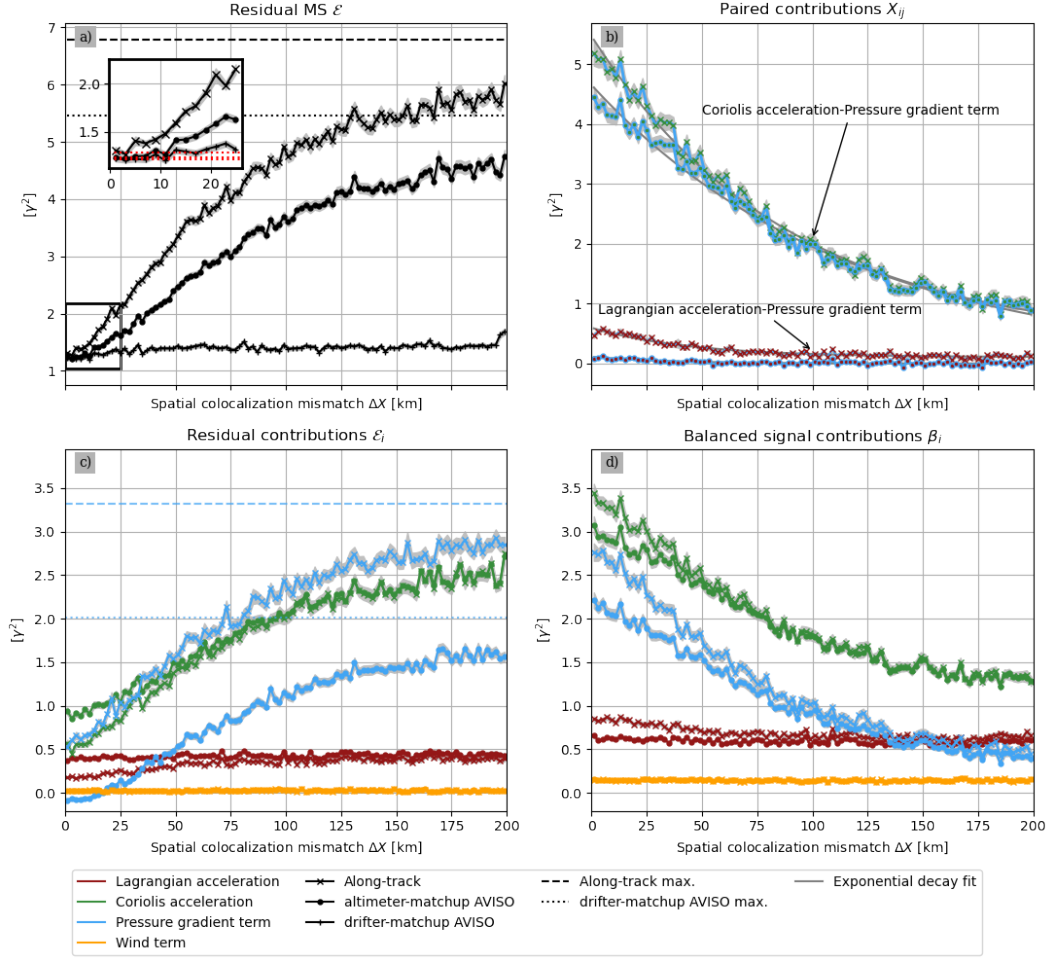


Figure 3. a) Residual dependency on spatial colocalization mismatch for the three different reconstructions (along-track, drifter-matchup AVISO and altimeter-matchup AVISO). The ultimate maximum values $\mathcal{E}_{-g\partial_x\eta} + A_{g\partial_x\eta}$ where the pressure gradient term would be fully uncorrelated are plotted as horizontal lines. Red lines on the inset highlight values for spatial colocalization mismatches of less than 4 km used to estimate global colocation errors. b) Dependency on spatial colocalization mismatch of the contribution of the Coriolis acceleration-pressure gradient term pair and the Lagrangian acceleration-pressure gradient term pair. The exponential decay fit is also plotted. c) and d) are respectively the dependency of the balanced signal and the residual contributions on spatial mismatch. Horizontal lines are the ultimate maximums for the pressure gradient residual contribution for along-track altimetry and altimeter-matchup AVISO reconstructions. Diagnoses were averaged over colocations in 2 km spatial mismatch bins. The 95% confidence intervals are plotted in gray.

451

3.2.2 Resolution mismatch errors

452

Given the high temporal resolution and local footprint of drifter observations compared with sea level and wind sources, we expect Lagrangian and Coriolis accelerations

453

454 to contain large resolution mismatch errors. Assuming both of these terms are devoid
 455 of instrumental noise (Appendix A1), subtracting the previously estimated respective
 456 colocation errors (Section 3.2.1 and Table 1) in the Lagrangian and Coriolis acceleration
 457 errors components leads to estimates of resolution mismatch errors of about $0.18 \gamma^2$ and
 458 $0.51 \gamma^2$, respectively (Table 1). We then explored whether these variances can be explained
 459 by unmeasured pressure gradient or wind terms. Altimetric sea level observations are
 460 smoothed with a 65 km low pass filter that removes all finer scale variability (Quality
 461 Information Document of the along-track product). Assuming the missing pressure gra-
 462 dient variance is the sum of Lagrangian and Coriolis acceleration mismatch errors ($\epsilon_r =$
 463 $0.18 \gamma^2 + 0.51 \gamma^2 = 0.69 \gamma^2$) and that signals missed in altimetric data are characterized
 464 by a spatial scale L of 65 km, a rough estimation of the corresponding sea level standard
 465 deviation is $\sqrt{\epsilon_r} \gamma \times L/g \sim 5$ cm. This estimated standard deviation being a fraction
 466 of the total sea level variability, we conclude that missing information in altimetric sea
 467 level may explain the resolution mismatch errors derived above. Having no knowledge
 468 on the information missing in wind reanalysis or on the quality of the physical model em-
 469 ployed to estimate vertical stress divergence, we cannot provide meaningful estimates
 470 of associated errors at this stage.

471 In principle, the pressure gradient can also contain resolution mismatch error if cor-
 472 related information is missing in the wind term. In the present state, there is little cor-
 473 relation between the pressure gradient and the wind term (Figure 2). In the literature,
 474 Bonjean and Lagerloef (2002)'s diagnostic model shows that equatorial surface dynam-
 475 ics can be simplified to the pressure-gradient-compensating wind term, potentially re-
 476 sulting in resolution mismatch errors in our study. However, this would be relevant only
 477 for a small proportion of the colocations; furthermore, we are not aware of other global
 478 observations or dynamical expectations regarding a systematic correlation between the
 479 pressure gradient and wind terms at the ocean surface. Moreover, the drifter-matchup
 480 AVISO pressure gradient term exhibits negligible errors compared with the balanced sig-
 481 nal component, thereby voiding the possibility of such resolution mismatch error for AVISO
 482 data (Figure 2b). We therefore conclude that resolution mismatch errors between the
 483 pressure gradient and wind terms are unlikely to have a strong signature on the global
 484 residual.

485 **3.2.3 Instrumental errors**

486 Drifter position data errors provided with the GDP dataset have a median of 60 m
 487 and 400 m for GPS and Argos drifters, respectively, and have repercussions on velocity
 488 and acceleration data. As shown in Appendix A1, most of these drifter instrumental er-
 489 rors in Coriolis and Lagrangian accelerations residual contributions are effectively removed
 490 from the studied residual by the LOWESS method (Elipot et al., 2016) and the tempo-
 491 ral 2.5 cpd low-pass filter.

492 For the pressure gradient, the processing that lead to along-track and AVISO sea
 493 level products mitigates instrumental noise via spatial low-pass filtering (65 km cutoff)
 494 and optimal interpolation, respectively (Quality Information Document of the along-track
 495 product). The pressure gradient term estimated from AVISO data is fully balanced out
 496 by other terms (Figure 2). AVISO instrumental noise is therefore necessarily negligible.

497 For the along-track reconstruction, after removing the colocation errors quantified
 498 in Section 3.2.1, about $0.57 \gamma^2$ of the pressure gradient residual contribution are still to
 499 be accounted for. The resolution mismatch error in the pressure gradient term is expected
 500 to be small as argued in Section 3.2.2. By process of elimination, instrumental noise in
 501 the pressure gradient term residual contribution must thus represent up to $0.57 \gamma^2$, which
 502 is also 36% of the residual MS. For the along-track product used, the remaining sea level
 503 error after the 65 km low-pass filtering is reported to lie between 0.85 and 1.1 cm, de-
 504 pending on the altimeter considered (Quality Information Document of the along-track

product). Assuming this error is spectrally white up to wavelengths δ_x of 65 km with a variance ϵ_η of 1 cm, the error on the pressure gradient term is $(g\epsilon_\eta \times 2\pi/\delta_x)^2/3 \sim 0.3 \gamma^2$. This is a loose confirmation that the remaining error on the pressure gradient may be explained by instrumental noise. Going further would require a more detailed analysis of the along-track altimetric noise spectral distribution and processing, its projection on the present diagnostics which fall outside of the scope of the present work.

3.2.4 *Missing physics errors*

Finally, our reconstruction of near-surface momentum balance neglects several physical processes contributions including, for instance, the baroclinic pressure contribution associated with horizontal density gradients $\partial_x \rho$. According to Fox-Kemper et al. (2011), such density horizontal gradient variance $\langle \nabla \rho^2 \rangle$ can reach maximum values of up to $3 \times 10^{-10} \text{ kg}^2 \cdot \text{m}^{-8}$ in simulations of the Southern Ocean. The error induced on momentum at depth scales h of 15 m is $(gh/\rho_0)^2 \times \langle \nabla \rho^2 \rangle = 0.006 \gamma^2$, which is a small fraction (about 0.4%) of the residual. Vertical advection of the horizontal momentum is neglected on the basis of the weakness of vertical velocity near the ocean surface. The estimation of the horizontal turbulence contribution is complicated here by the isolated nature of drifter observations, which do not provide meaningful grounds for scale separation. In conclusion, in this study, we assumed that missing physics errors are small compared with the other error sources.

3.3 Geographical analysis

3.3.1 *Momentum balance closure*

We now turn away from the global-scale average diagnoses to describe how momentum balance varies regionally. In an absolute sense, closure is best in the equatorial zone and in moderately energetic regions such as oceanic gyre centers (Figure 4a). At mid to high latitudes and/or in energetic areas, momentum balances closure is lower, with residual MSs typically exceeding the global average ($1.64 \gamma^2$). A potential explanation involves decreased values of energy-dominant spatial scales at these latitudes as indicated by observed eddy sizes in Chelton et al. (2011), which fall below 75 km at latitudes polewards of 45° . This shortening of the scales of variability can increase the amount of unresolved variability in altimetric data, thereby increasing resolution mismatch error on Coriolis acceleration. This difference in variability scales would also increase collocation errors. Both trends are consistent with the larger values of Coriolis acceleration and pressure gradient residual contributions (Figures 4c and 4d). Furthermore, the larger increase in Coriolis acceleration residual contributions compared with those of the pressure gradient may indicate that resolution mismatch errors weigh more than collocation errors on the residuals in these areas. In the Aghulas, Gulf Stream, and Kuroshio current systems, the residual is driven by pressure gradient term contribution (Figure 4d). As geostrophy is the main contributing equilibrium in these areas, collocation errors in the residual contributions of the pressure gradient term and of the Coriolis acceleration should be of the same order of magnitude. Given the more modest Coriolis residual contribution, the pressure gradient residual contribution observed in these areas is thus unlikely to result from collocation errors (Figures 4c and 4d). Resolution mismatch errors associated with the wind term are also unlikely, because winds are not particularly strong there. These errors thus appear to reflect an increase in altimetric instrumental errors, which may potentially be related to the surface wave field via the heterogeneities induced by its interactions with ocean circulation (Quilfen et al., 2022).

3.3.2 *Dynamical regimes*

We identified characteristic dynamical regimes in selected regions (contours on Figure 4a). These regions were chosen for the contrasting nature of their momentum bal-

554 ance closures, as indicated by dynamical paired contributions and a priori dynamics (e.g.
 555 expected strong geostrophy in the Gulf Stream etc (Yu et al., 2021)). The Gulf Stream
 556 region (GS) is defined by locations between 15°N and 50°N and 85°W and 0°W where
 557 the geostrophic pair contribution is greater than 70%. The Antarctic Circumpolar Cur-
 558 rent region (ACC) is composed of locations southwards of 35°S , where the geostrophic
 559 pair contribution is greater than 40%. The North Pacific region (NP) is the region ex-
 560 tending from 120°E to 150°W and northwards of 40°N where the Lagrangian acceler-
 561 ation - Coriolis acceleration pair contribution exceeds 10%. This definition thus excludes
 562 the Kuroshio Current. Finally, the Equator region (EQ) is defined as the area in between
 563 15°S and 15°N .

564 In GS and ACC, geostrophy dominates the ocean surface dynamics, explaining about
 565 80% of the momentum balance closure, compared with 57% for the global average (Fig-
 566 ure 5a). Remarkably, the same percentage of geostrophic variance in these energetic re-
 567 gions was also estimated through the numerical reconstruction of Yu et al. (2021). The
 568 Lagrangian-Coriolis and Lagrangian-pressure gradient pairs contribute relatively less to
 569 the momentum closure compared with the global averages, emphasizing the relatively
 570 moderate amplitude of cyclogeostrophic corrections, even in these energetic areas. In ACC,
 571 the pressure gradient-wind terms pair is barely higher than the global average in rela-
 572 tive terms (7% vs. 5%), but twice as high in absolute values, which is consistent with
 573 our expectations of strong winds in this region.

574 In NP, the dominant dynamical equilibrium is that associated with the Lagrangian-
 575 Coriolis pair (41%, Figure 5c), reflecting the higher near-inertial variability in this re-
 576 gion as reported in Yu, Ponte, et al. (2019); Liu et al. (2019); Flexas et al. (2019). The
 577 geostrophic pair (30%) is of secondary importance. Internal tides may also contribute
 578 to the balance via the combinations with these pairs.

579 The momentum balance observed in EQ is in stark contrast with those in all other
 580 regions, with the dominance of the Lagrangian-pressure gradient pair (30% of the clo-
 581 sure, Figure 5d). Internal tides are energetic in this area (Buijsman et al., 2017) and may
 582 contribute to momentum balance via the Lagrangian-Coriolis and Lagrangian-pressure
 583 gradient pairs. The modest contribution of the pressure gradient-wind term pair and the
 584 absence of correlations with pressure gradient terms (expected from Bonjean and Lager-
 585 loef (2002)) may most likely be explained by the poor estimation of the wind term. These
 586 poor estimates may arise due to the use of global Ekman parameters and the pronounced
 587 variations of these parameters away from global averages at near equatorial latitudes (Rio
 588 et al., 2014). Given the lower noise level in altimetric data near the equator (Figure 4d),
 589 our inability to close momentum balances there may therefore be driven by this poor es-
 590 timation of the wind term.

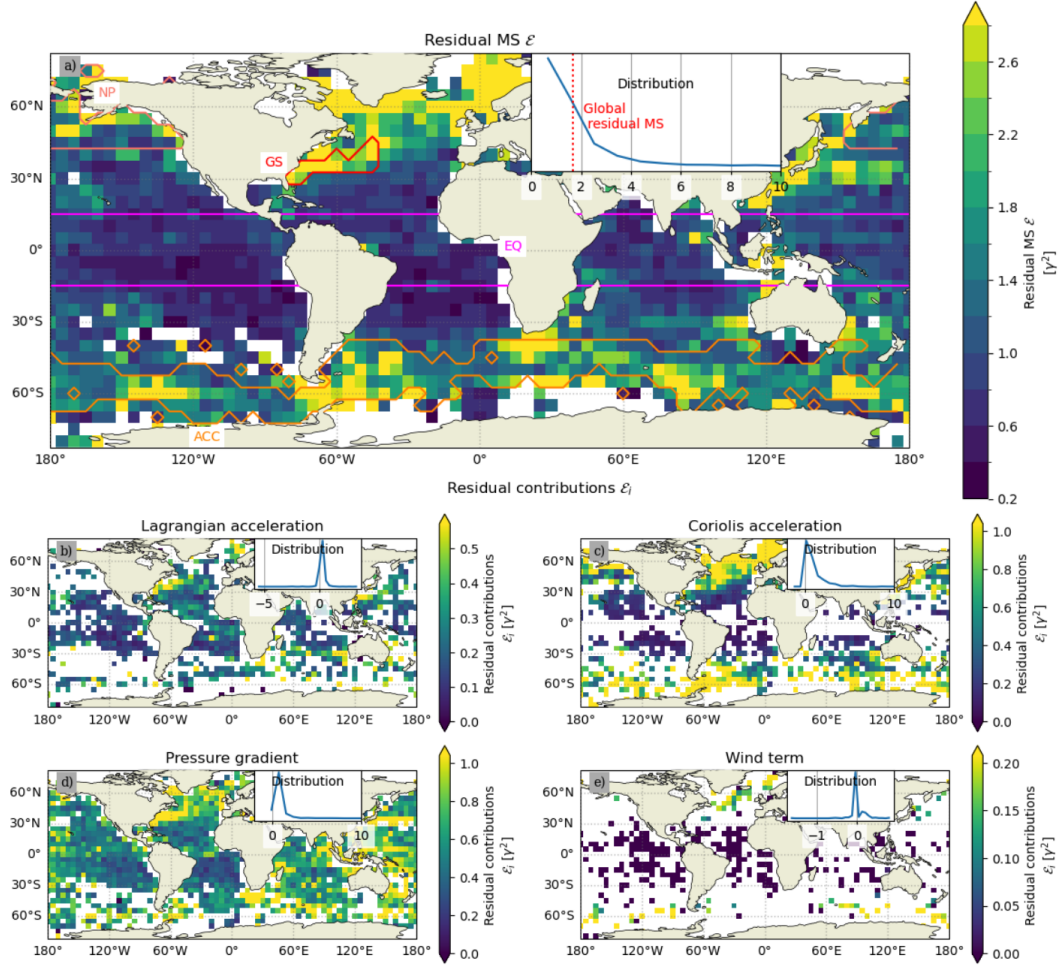


Figure 4. The residual mean square value (a) and residual contribution of Lagrangian acceleration (b), Coriolis acceleration (c), pressure gradient (d) and wind (e) terms mapped in 5° -geographical bins. Only bins in which the relative statistical errors are less than 50% are shown. Contours in a) delimit the different dynamical regimes in Figure 5. Note that the color bar scales differ between graphs.

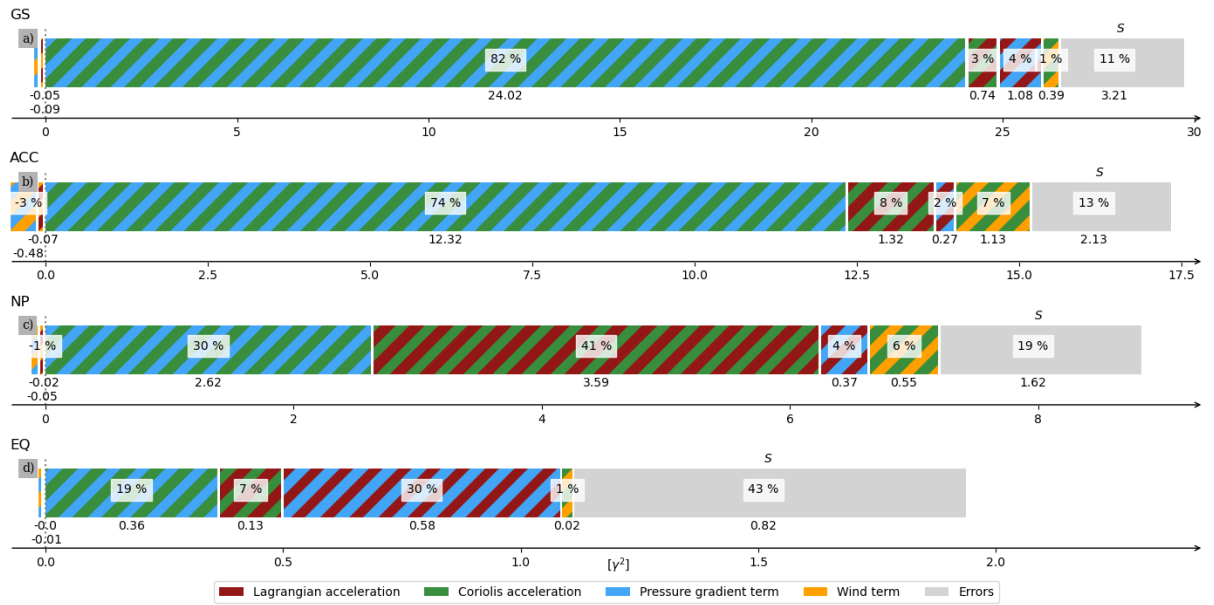


Figure 5. Different decompositions of paired contributions for areas plotted on Figure 4a i.e. the Gulf Stream (GS), the Antarctic Circumpolar Current (ACC), the North Pacific (NP) and the Equatorial Band (EQ). Note the differences between the horizontal axes.

4 Conclusion

By combining current and sea level observations and wind reanalysis, we successfully reconstructed global instantaneous near-surface horizontal momentum balances at a 80% closure degree (β/Σ). This success motivated the investigation of dominant terms and dynamical equilibria, and geographic sensitivities. As expected, geostrophy is the leading order equilibrium, explaining about two thirds of the variance in the balanced signals at a global scale. In western boundary current systems and ACC, geostrophy dominates and explains up to 81% and 91% of the balanced signal variance, respectively. However, our results also highlight that about one third of the ocean surface dynamics resolved here is not geostrophic. As reported by Yu et al. (2021), Lagrangian and Ekman flows thus cannot be neglected everywhere when reconstructing instantaneous near-surface balances. Geostrophy explains only one third of the balanced signal variance in the North Pacific (NP) and the equatorial band (EQ), for instance. We clearly demonstrated this specificity using observationally based quantitative information.

Gridded AVISO sea level reconstructions provide global residual MSs that are about 20% lower than with along-track sea level. This difference suggests that AVISO mapping deletes more instrumental errors than it introduces resolution mismatch errors. In addition, the novel methodology introduced here reveals that along-track altimetry presumably contains the signature of short-scale variability that is not present in AVISO and balances out better with the variability contained in drifter observations. Along-track altimetry effectively enables the compensation of Lagrangian acceleration by the pressure gradient that AVISO cannot.

We detailed the quantitative error budget for global closures using along-track sea level data. Colocation mismatch errors (between drifter and altimetric observations) are mainly spatial, and account for 21% of the global residual MS. Short-scale variability contained in drifter observations, but not in other data sources (altimetry, wind) results in a resolution mismatch errors that we estimate to be about 42%. Instrumental errors on sea level observations likely account for 36% of the residual. The remaining 1% are wind term errors, that are driven either by the quality of the wind reanalysis or by the accuracy of the Ekman model.

One limitation of our analysis stems from the quality of the Ekman model used for the estimation of the wind term. The Ekman model used is indeed implemented with global values of its parameters and therefore neglects the seasonal and regional variability reported in Rio et al. (2014). Departures from global parameter values are especially large at the equator for the amplitude parameter; this can only have an adverse effect on momentum balance closure. These departures may for instance explain the reason why we are unable to observe the equatorial pressure gradient-wind term compensation modeled by Bonjean and Lagerloef (2002). Further research can explore the seasonal and regional sensitivities of Ekman parameters such as those computed in Rio et al. (2014). The estimation of the vertical stress divergence may also benefit from more complete Ekman models that use knowledge on the mixed layer depth (Mulet et al., 2021) and/or the wind time history (Lilly & Elipot, 2021). At present, we cannot directly quantify the role of ERA* product inaccuracies on momentum closure. In this analysis, we relied on ERA* being the best product available to estimate the wind term. A sensitivity analysis of different wind reanalyses would be an interesting undertaking. In the preliminary stages of the analysis, we compared ERA5 with ERA* with no significant differences in global momentum closure.

The spatial colocation mismatch criterion set to 25 km resulted from a compromise between the mitigation of colocation errors and statistical reliability. It is possible to adjust this criterion regionally to match the geographical and temporal (seasonal) distributions of ocean spatial scales. For example, higher mismatch values at the equator may be able to help account for larger spatial scales (Chelton et al., 2011). Adjusting this cri-

643 terion however requires building prior knowledge on these distributions. An a posteri-
644 ori consideration of the small sensitivity of residuals on temporal mismatch can allow
645 increasing the temporal tolerance when assembling colocations. Relaxing the temporal
646 tolerance should be done cautiously, because it can affect closure in areas where ener-
647 getic high-frequency motions are present and may also introduce correlated colocations.
648 The largest benefit of these refinements is expected to be for geographically bin-averaged
649 diagnoses whose statistical reliability is currently limited by the amount of available colo-
650 cations. Another benefit would be to improve the closure in global diagnoses.

651 For the global averages, we did not attempt to account for biases introduced by non-
652 uniform geographical sampling (driven by joint drifter and altimeter availability), due
653 to limited data availability. Our methodology cannot identify three term balances (e.g.
654 cyclogeostrophy, internal tide dynamics); clustering approaches need to be developed to
655 do so. Statistical diagnoses may have been conditioned on temporal (seasonal) criteria
656 to highlight associated variability. Testing our momentum reconstruction methodology
657 with numerical simulations of ocean circulation can help assess the above-cited limits and
658 actual usefulness of the proposed improvements. In particular, the signature of specific
659 dynamical processes (e.g. internal tides, near inertial waves) in simplified simulated cases
660 can help identify such processes in reconstructions based on observational data.

661 With its unprecedented resolution and accuracy (Fu et al., 2024), the use of SWOT
662 wide-swath sea level is expected to substantially reduce instrumental and resolution mis-
663 match errors, and to considerably improve horizontal surface momentum balance recon-
664 structions. The two-dimensional nature of SWOT sea level observations will make it pos-
665 sible to reconstruct momentum balance in both horizontal dimensions simultaneously.
666 The wide-swath sea level data will also lead to a larger number of colocations and smaller,
667 if not negligible, spatial colocation mismatch errors. The reconstructions of horizontal
668 momentum balances and the methodology developed here should be instrumental to 1/
669 gain an observation-based understanding of the upper ocean dynamics, and 2/ assess the
670 need to go beyond the framework employed to date to estimate the ocean circulation,
671 namely geostrophy plus Ekman. This research effort can also contribute to validating
672 SWOT sea level observations and to improving our ability to estimate MDT from these
673 observations.

674 Appendix A Reconstruction sensitivity

675 Colocation datasets constructed as the one in Section 2.2, but considering a dif-
 676 ferent drifter type, drogue status, and low-pass filter are built. Here, the sensitivity of
 677 closure to the low-pass filter cutoff frequency applied on drifter trajectories, the use of
 678 GPS or Argos positioning and drogued or undrogued drifters are studied using meth-
 679 ods described in Section 2.4.

680 A1 Applying low-pass filters on drifter trajectories

681 We low-filtered drifter trajectories at different cutoff frequencies with a finite im-
 682 pulse response filter and applied forwards and backwards on drifter velocity. The com-
 683 bined filter thus has zero phase and a filter order twice that of the original.

684 Applying a 2.5 cpd low-pass filter on drifter trajectories reduces the residual con-
 685 tribution of the Lagrangian acceleration term by $0.84 \gamma^2$ (-44% of its MS), whereas the
 686 balanced signal contribution remains approximately constant (Figure A1a). This 2.5 cpd
 687 low-pass filter thus removed a substantial part of acceleration errors from drifter obser-
 688 vations. GDP drifter instrumental velocity errors were estimated using the LOWESS smooth-
 689 ing method and directly available in the GDP dataset (Elipot et al., 2016). The mean
 690 velocity error is about $\epsilon_v = 2 \times 10^{-2} \text{m.s}^{-1}$, yielding approximate instrumental error
 691 variances of $(f\epsilon_v)^2 = 0.04 \gamma^2$ for Coriolis acceleration and $\frac{\epsilon_v^2}{2dt^2} = 0.40 \gamma^2$ for Lagrangian
 692 acceleration after centered finite differentiation. These errors are comparable to the am-
 693 plitudes of the signal filtered out with the 2.5 cpd low-pass filter described above and
 694 corroborate the hypothesis that this information mostly represents instrumental posi-
 695 tion error. Differences between these two quantities can be explained by an imperfect
 696 velocity error estimation or by unbalanced resolution mismatch errors that are filtered
 697 out with instrumental errors. In any case, these results justify the use of low-pass filter-
 698 ing on drifter data for the global momentum balance reconstruction (Section 3.1).

699 Interestingly, decreasing the cutoff frequencies from 1.5 to 0.5 cpd progressively can-
 700 cels the balanced signal contribution of Lagrangian acceleration and reduces that of Cori-
 701 olis acceleration (Figure A1 a) and b)). This leads to a decrease in the balanced signal
 702 contributions of the other terms (Figure A1 c) and d)). These two filters thus hinder clo-
 703 sure, but these experiments demonstrate that high-frequency processes, such as near in-
 704 ertial and tidal motions, are captured and do contribute to momentum balance closure.

705 A2 GPS vs. ARGOS

706 Despite their original difference in terms of positioning error (median of 60 m for
 707 GPS vs. 400 m for Argos), reconstruction with GPS or Argos drifters give similar results
 708 in terms of the balanced signal and residual contributions of the equation once filtered
 709 with 2.5 cpd low-pass filter (Figure A1e). This similarity justifies the use of both GPS
 710 and Argos drifters, making it possible to nearly double the number of available coloca-
 711 tions.

712 A3 Drogued vs. undrogued

713 Undrogued drifters follows currents at the very surface and are more subject to wind
 714 drift than drogued drifters (Poulain et al., 2009). Therefore, the undrogued Coriolis ac-
 715 celeration MS is higher than the drogued one ($+1.26 \gamma^2$). This additional variability re-
 716 mains unbalanced in the reconstruction, and is thus added in the residual contribution
 717 of the Coriolis term (Figure A1f). This additional variability therefore justifies using only
 718 drogued drifters for our reconstruction analysis.

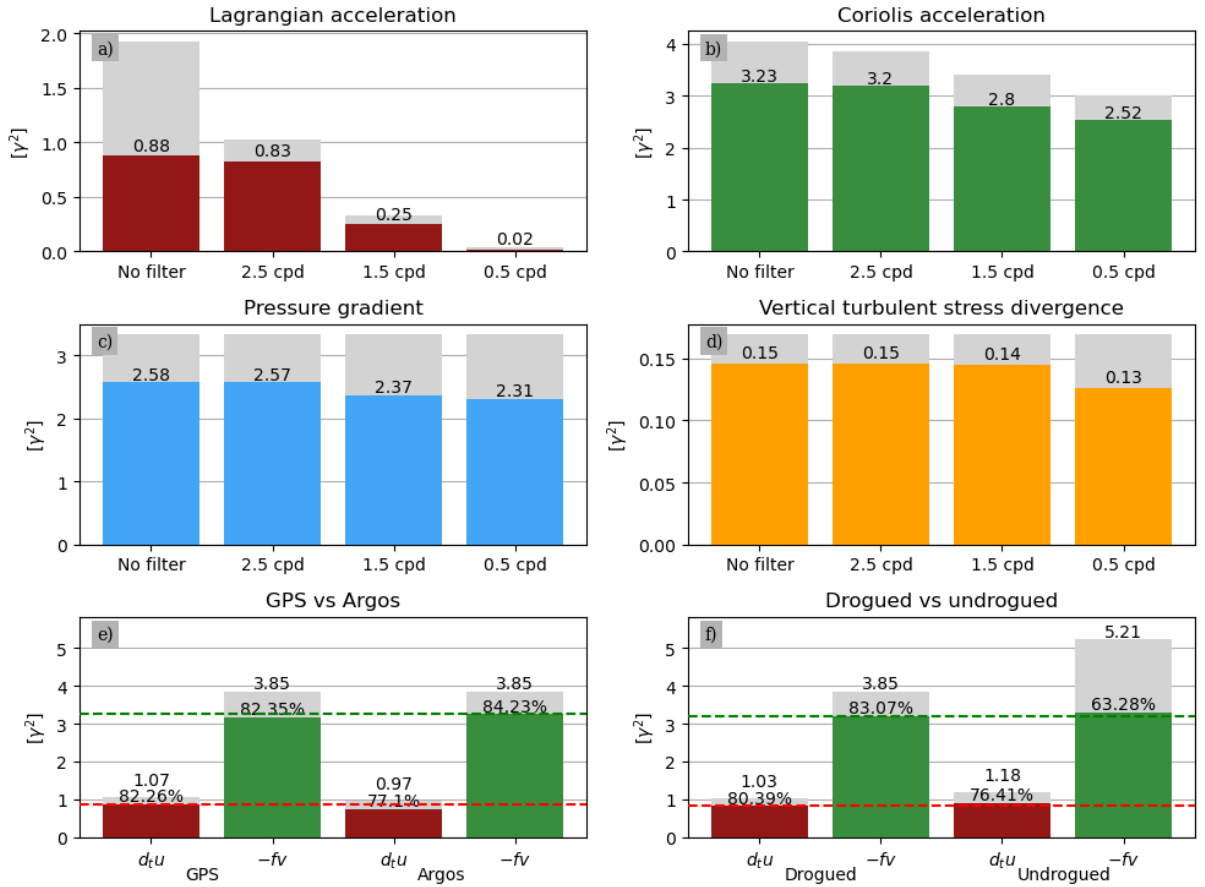


Figure A1. Decomposition into balanced physical signal (in color) and error (gray) components: (a to d) for all terms applying no filter or 2.5, 1.5 and 0.5 cpd-low pass filter on drifter trajectories; comparing the Lagrangian and Coriolis accelerations for drogued GPS and Argos drifters (e) and for drogued and undrogued drifters (f) with 2.5 cpd low-pass filtering

Open Research Section

Archived datasets used are available at the following references:

- for ERA* data : Trindade et al. (2022)
- for L3-along-track altimetry: European Union-Copernicus Marine Service (2021a)
- for L4-AVISO-gridded altimetry: European Union-Copernicus Marine Service (2021b)

All codes and notebooks used are available in the dedicated repository: <https://doi.org/10.5281/zenodo.14513262> (Demol, 2024).

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