

Mesoscale Ocean Processes: The Critical Role of Stratification in the Icelandic Region

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Key Points:

- High-resolution SWOT satellite data, *in situ* data, and numerical models are used to analyze ocean dynamics around Iceland
- North of Iceland, the shallow upper layer creates a "low-energy zone" with reduced energy at the mesoscale
- South of Iceland, the deep upper layer boosts the mesoscale eddy generation through baroclinic instability

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Abstract

The oceanic region around Iceland, a key component of the Atlantic Meridional Overturning Circulation, plays a critical role in global climate through its complex system of surface and subsurface currents. Using high-resolution SWOT satellite data, *in situ* observations, and idealized numerical simulations, this study reveals two distinct dynamical regimes in this region. South of Iceland, the mesoscale eddy field is energetic. In contrast, north of Iceland, the shallow upper layer inhibits baroclinic instability and eddy generation, resulting in a low-energy "eddy desert" — a phenomenon observed synoptically for the first time at SWOT's unprecedented resolution. This new understanding of Iceland's dynamical regimes highlights the role of stratification in mesoscale variability. Analyzing Biogeochemical-Argo float data, it also suggests an impact of mesoscale regimes on the local biogeochemical cycles, with implications for primary production and carbon cycling as stratification patterns shift with climate change.

Plain Language Summary

The waters around Iceland are an important part of the global climate system, influencing ocean currents and climate patterns. In this study, we used cutting-edge satellite data, ocean measurements, and computer models to investigate the behavior of ocean currents in the region. We discovered two distinct ocean zones around Iceland. South of Iceland, strong ocean currents create a turbulent flow. North of Iceland, however, the water is more stable, and the ocean is much calmer, with fewer currents and less mixing. This creates a "low-energy zone" where the usual ocean activity is limited. Our high-resolution data captured this unique, and previously unexplored feature. This study helps us understand how the movement of water around Iceland affects the ocean environment, with important implications for climate change and the future of marine ecosystems, as warmer temperatures may change these ocean patterns.

1 Introduction

The oceanic region surrounding Iceland plays a vital role in the Atlantic Meridional Overturning Circulation (AMOC), a major driver of global climate (Buckley & Marshall, 2016). Surface and subsurface currents in this area are shaped by both, regional ocean dynamics and complex topographic features, influencing the exchange of water masses between the Arctic and Atlantic oceans (Brakstad, Gebbie, et al., 2023). To the north of Iceland, the Iceland Sea is encompassed by the Greenland-Iceland Ridge and the Jan Mayen Ridge, acting as a physical barrier between the Arctic and Atlantic waters (Fig. 1). To the east, the Iceland-Faroe Ridge separates the Norwegian Sea and the Iceland Basin, while in the south of Iceland, the Reykjanes Ridge defines the boundary between the Iceland Basin and the Irminger Sea. The surface circulation is dominated by warm and saline North Atlantic waters flowing from the southwest. It is composed of the North Atlantic Current (NAC) which flows into the Iceland Basin, and the Irminger Current (IC) which flows into the Irminger Sea and detaches into the North Icelandic Irminger Current (NIIC) flowing clockwise reaching the north of Iceland (Hansen & Østerhus, 2000). The other main surface current is the East Greenland Current (EGC), transporting cold and fresh polar waters mainly southward along the east coast of Greenland. A small component of this branch is advected by the NIIC forming the East Icelandic Current (EIC, Logemann et al., 2013; Semper et al., 2022). At mid-depths, northeast of Iceland, the North Icelandic Jet (NIJ) injects a small portion of dense water into the Denmark Strait (Jonsson & Valdimarsson, 2004; Casanova-Masjoan et al., 2020). All of these currents transport water masses that, together with the deep cold and denser waters: Iceland Scotland Overflow Water (ISOW) and the Denmark Strait Overflow Waters (DSOW), are key components of the AMOC.

Over the last two decades, our understanding of regional circulation in this area has steadily improved. However, the dynamics at smaller scales, particularly at the mesoscale and below, remain poorly understood. This is due to a scarcity of *in situ* data (Beaird et al., 2013) and the inability of numerical models and satellite data, until now, to resolve these scales both south and north of Iceland. In the Iceland Basin, several studies have documented the presence of mesoscale eddies, examining their origins, variability, and impact on regional circulation and nutrient distribution (Mahadevan et al., 2012; Godø et al., 2012; Zhao et al., 2018a; Soman et al., 2022; Z. Zhang et al., 2024; Johnson et al., 2024; Voet et al., 2024). These studies rely on AVISO-derived satellite altimetry ($1/4^\circ$ resolution), which is at the limit of what is needed to resolve the mesoscale eddy field in the region (eddies in the Iceland Basin have an average radius of 50 km; see Soman et al., 2022). However, north of Iceland, in the Iceland Sea, pre-SWOT satellite altimetry has been largely ineffective due to the smaller deformation radius and the correspondingly small eddy sizes ($\mathcal{O}(10)$ km; Chelton et al., 1998). Past studies have primarily focused on regions

68 bordering Norway or on large, semi-permanent structures such as the Lofoten Vortex (Bosse et al., 2019)
 69 and intense Norwegian Sea eddies, which are influenced by their Atlantic water component (Bashmach-
 70 nikov et al., 2023). Thus, mesoscale-focused studies are scarce north of Iceland and *absent in the Iceland*
 71 *Sea*.

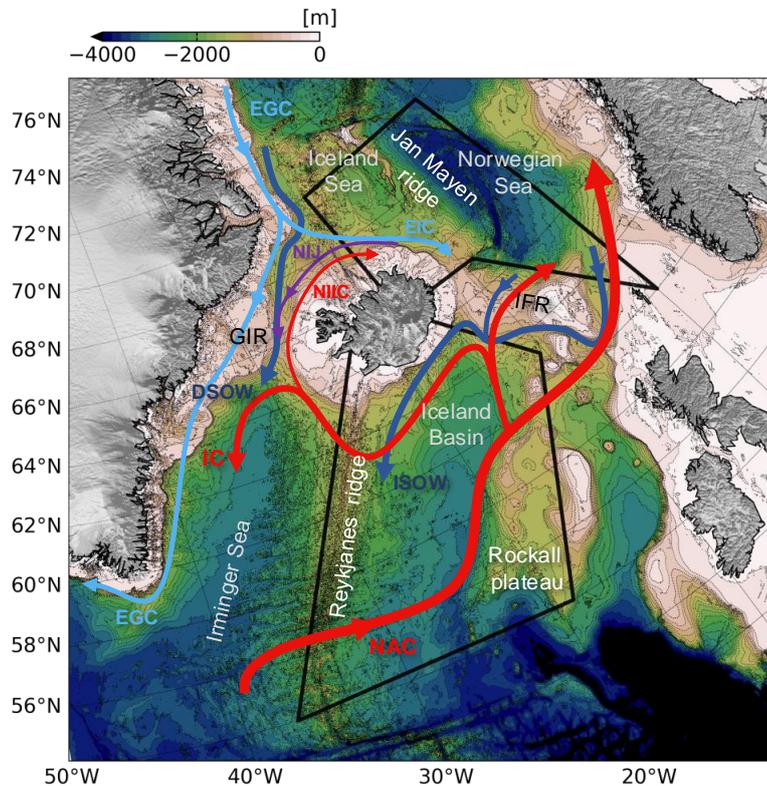


Figure 1. Bathymetry around Iceland highlighting the main topographic features, basins and seas; IFR and GIR stand for Iceland-Faroe Ridge and Greenland-Iceland Ridge, respectively. The areas we denominate as "south of Iceland" and "north of Iceland" in the study are indicated by the black solid lines. Bold arrows show schematic pathways of the main currents in the area: the North Atlantic Current (NAC), the East Greenland Current (EGC), the East Icelandic Current (EIC), the North Icelandic Jet (NIJ) and the North Icelandic Irminger Current (NIIC); dark blue arrows show the pathway of the two main bottom water masses flowing south: the Denmark Strait Overflow Waters (DSOW) and the Iceland Scotland Overflow Water (ISOW).

72 The surface currents around Iceland advect and mix the water masses. They shape the character-
 73 istic thermocline stratification and the mixed layer depth in the Iceland Basin, Iceland Sea and Norwe-
 74 gian Sea. The mixed layer is the uppermost layer of the Ocean, which is in direct contact with the at-

75 mosphere, where exchanges of heat, oxygen, carbon, and other tracers take place (Bindoff et al., 2019;
76 Sallée et al., 2021). The physical processes within this layer influence injection of tracers into the ocean
77 interior, including the oceanic CO₂ uptake and nutrient cycles, which are key components in the regu-
78 lation of the Earth’s climate (Ólafsson, 2003; Ruiz et al., 2019). At smaller scales, particular processes,
79 such as meso and submesoscale eddies, modulate the mixed layer depth and enhance vertical fluxes through
80 eddy-driven upwelling and downwelling (McGillicuddy et al., 2007; Mahadevan et al., 2012). Eddies there-
81 fore act as hotspots for primary production and carbon sink from the surface to the ocean interior (Ruiz
82 et al., 2019). Stratification changes in the upper-ocean layers have been reported in the last 50 years as
83 the result of changes in temperature and salinity (in particular saltier and warmer surface water, see Polyakov
84 et al., 2017; Dai et al., 2019). These changes are particularly faster in the region around Iceland than the
85 global average and are among the critical variables within the IPCC reports (Bindoff et al., 2019; Sallée
86 et al., 2021).

87 As sea surface temperatures continue to increase (Bindoff et al., 2019; Pörtner et al., 2019), it is crit-
88 ical to understand how oceanic currents variability may change in a warming and potentially more strat-
89 ified ocean, ultimately having implications for ocean productivity, carbon uptake and sequestration. In
90 this study, we analyze the characteristics of mesoscale ocean dynamics around Iceland, using newly re-
91 leased SWOT satellite data, ship-based observations, numerical modeling, and Argo float data. We com-
92 pare the regions north and south of Iceland; the northern region encompasses the Iceland Sea and the
93 Norwegian Sea, while the southern region corresponds mainly with the Iceland basin, see definition in Fig. 1.
94 We demonstrate the existence of two distinct dynamical regimes. On the one hand, the southern region
95 consist of highly energetic *balanced geostrophic turbulence* (hereafter called *turbulence*, or *turbulent flow*
96 for conciseness) with intense mesoscale activity. On the other hand, the northern region presents a stark
97 contrast, acting as an "eddy desert" with significantly less mesoscale variability. Our analyses show that
98 these distinct regimes are primarily driven by the difference in water column properties between these
99 two regions: the upper layer is significantly thicker in the south, compared to the north. We finally dis-
100 cuss the potential implications for the biological productivity in these two distinct regimes .

101 **2 Data and methods**

102 **2.1 T/S properties around Iceland**

103 To analyze the thermohaline properties of the ocean around Iceland, we utilize the COriolis Ocean
104 Dataset for Reanalysis product (hereafter CORA, Szekely et al., 2019). It is a global gridded dataset of

105 *in situ* temperature and salinity measurements from autonomous platforms (Argo profilers, fixed moor-
 106 ings, gliders, drifters, sea mammals) and research or opportunity vessels (CTDs, XBTs, ferrybox). The
 107 dataset provides monthly temperature (T) and salinity (S) profiles at a $0.5^\circ \times 0.5^\circ$ resolution over the
 108 period 1960–2023. We restrained our analysis to the 2000–2023 period to maximize the number of au-
 109 tonomous floats representation in the dataset.

110 2.2 Normal vertical mode decomposition

111 To assess the impact of stratification on ocean dynamics, an advanced approach involves decom-
 112 posing the ocean’s vertical structure into normal dynamical modes. This method allows for the separa-
 113 tion of the water column into different depth-dependent modes that capture the full extent of stratifi-
 114 cation influences on the flow (Vallis, 2017). The dynamical regimes discussed here, *i.e.*, the mesoscale ei-
 115 ther north or south of Iceland, are characterized by a relatively small Rossby number (Ro) and timescales
 116 longer than a day. In this case, the ocean can be fairly well described by the continuously stratified Quasi-
 117 Geostrophy (QG) on an f -plane. It is formulated as the material conservation of the QG Potential Vor-
 118 ticity Q_{qg} (Vallis, 2017) as

$$[\partial_t + \mathcal{J}(\Psi, \cdot)] Q_{\text{qg}} = 0, \quad (1)$$

119 where

$$Q_{\text{qg}} = \nabla^2 \Psi + \frac{\partial}{\partial z} \left(\frac{f_0^2}{N^2} \frac{\partial}{\partial z} \Psi \right). \quad (2)$$

120 $\Psi(x, y, z, t)$ is the scalar 3D streamfunction, \mathcal{J} is the Jacobian operator, f_0 the local Coriolis frequency,
 121 N the Brunt–Väisälä frequency, and x, y, z , the zonal, meridional, and vertical coordinates, respectively.
 122 The water-column dynamics are therefore intrinsically linked to the stratification of the water column
 123 N , through the so-called *stretching term* (the last term on the right-hand side in Eq. (2)). This stretch-
 124 ing term can be decomposed into vertical dynamical modes (Gill, 1982) by projecting the streamfunc-
 125 tion over normal modes of vertical structure $h_n(z)$ as

$$\Psi(x, y, z, t) = \sum_{n=0}^{+\infty} \psi_n(x, y, t) h_n(z), \quad (3)$$

126 and solving the eigenvalue problem

$$\frac{\partial}{\partial z} \left(\frac{1}{N^2} \frac{\partial}{\partial z} h_n \right) + \frac{1}{c_n^2} h_n = 0, \quad (4)$$

127 where c_n are the eigenvalues of the n^{th} vertical mode, which is linked to the deformation radius λ_n of each
 128 mode by $\lambda_n = c_n^2/|f_0|$. We hereafter define $\lambda_1 = R_D$, as the first baroclinic deformation radius (Chel-
 129 ton et al., 1998)). The modal decomposition is performed using the Python Dedalus library, with the same
 130 boundary condition as presented in Tedesco et al. (2022). This gives an estimate of h_n and λ_n for $0 \leq$
 131 $n \leq 10$.

132 This method is applied in our area of interest using the CORA dataset (see section 2.1), the T/S
 133 measurements from Argo float #4903532 (see section 2.6), and *in situ* T/S measurements from two par-
 134 ticular repeated stations north of Iceland (blue and red crosses in Fig. 3a) over the 1990-2018 period.

135 2.3 Quasi-Geostrophic idealized simulations of forced turbulence

136 We design idealized simulations that simulate the mesoscale oceanic dynamics north and south of
 137 Iceland. There, the flow mainly consists of a turbulent flow (Villas Bôas et al., 2022), with baroclinic in-
 138 stability (BCI) as its primary source of energy (Callies et al., 2016). It can mainly be represented by the
 139 barotropic and first baroclinic components (*i.e.*, truncate Eq. (4) to $n = 0$ and $n = 1$, a 2-layer ocean,
 140 see Flierl, 1978). In these conditions, BCI can occur if the meridional PV gradients change sign between
 141 the surface and the bottom layer (Pedlosky, 2013), which is always the case on an f -plane if there exists
 142 a baroclinic component. The BCI eventually generates eddies of typical size πR_D (Vallis, 2017) thus defin-
 143 ing the so-called mesoscale. The accuracy of considering only the first two modes was confirmed by use
 144 of the ECCO re-analysis (Fukumori et al., 2021); specifically, the RMSE between the current reconstruc-
 145 tion from the first two modes and the real current is less than 10^{-3} in the open ocean in our area of in-
 146 terest (not shown here).

147 With a mean baroclinic zonal flow $\mathbf{U} = (U, -U)$, Eq. (1), can be re-written in a non-dimensionalized,
 148 2-layer form as

$$\partial_t q_i + \mathcal{J}(\psi_i, q_i) \pm \partial_x q_i \mp F_i \partial_x \psi_i = 0, \quad (5)$$

149 for $i = 1, 2$ the upper and bottom layers, respectively. Where ψ_i is the streamfunction anomaly and q_i
 150 is the potential vorticity anomaly. Let's assume that the surface layer is in the rigid lid approximation
 151 and the bottom layer has no bottom friction, whose respective thicknesses are h_1 and h_2 , thus defining
 152 the aspect ratio $\delta = h_1/h_2$, $F_1 = 1/(1+\delta)$, and $F_2 = \delta/(1+\delta)$. All quantities are non-dimensionalized

153 such that $x, y \sim R_D$, $u, v \sim U$, and $t \sim R_D/U$. In this framework, the only parameters that vary are
 154 U , R_D , and δ .

155 We integrate these equations in a doubly periodic domain, on a 256×256 points grid, with timesteps
 156 adjusted to respect the CFL criterion. Following the formulation of Callies et al. (2016), we add dissipa-
 157 tion terms to the right-hand side of the equations as

$$\partial_t q_i + \mathcal{J}(\psi_i, q_i) \pm \partial_x q_i \mp F_i \partial_x \psi_i = r \nabla^{-2} q_i - \nu (-\nabla^2)^n q_i. \quad (6)$$

158 Small scales are damped through hyperviscosity of order $n = 10$ coefficient ν , and large scales through
 159 hypoviscosity coefficient r . These later parameters are chosen to be the same in all simulations, and to
 160 be the smallest for the simulations to be at equilibrium after the turbulence has set up. They are sim-
 161 ilar to the values used in de Marez & Callies (2025). All quantities are re-dimensionalized after running
 162 the model, and we set $U = 0.04 \text{ m s}^{-1}$ and $R_D = 15 \text{ km}$ in all simulations. Note that we choose to fo-
 163 cus on the open ocean dynamics, thus, the doubly periodic domain allows to eliminate spurious bound-
 164 ary effects and it let eddies evolve freely. As the result, the simulations do not capture shelf-break and
 165 coastal processes. This approach is a well-established method for studying open-ocean mesoscale dynam-
 166 ics (see the literature of vortex studies since McWilliams & Flierl, 1979).

167 2.4 SWOT data and spectral analysis

168 We leverage newly released satellite data from the Surface Water and Ocean Topography (SWOT,
 169 see some background in *e.g.*, Morrow et al., 2019) program, a collaborative effort between NASA and CNES
 170 launched in late 2022, to unveil unprecedented details of surface mesoscale turbulence in the ocean around
 171 Iceland. Specifically, we use the latest release (v2.0) of SWOT_L3_SSH 'Basic' product, derived from the
 172 L3 SWOT KaRIn Low rate ocean data products provided by NASA/JPL and CNES. This dataset is pro-
 173 duced and freely distributed by the AVISO and DUACS teams as part of the DESMOS Science Team
 174 project (AVISO/DUACS, 2023). The "noise-reduced" and "raw" Sea Surface Height anomaly (SSHa),
 175 displayed on a 2-km resolution grid, are used for our analysis (see Dibarboure et al., 2023, for details on
 176 the method). We use data on the 21-day repeat orbit period, from 1st September 2023 to 31st August
 177 2024 to cover a full seasonal cycle. Recent studies (X. Zhang et al., 2024; Z. Zhang et al., 2024; Verger-
 178 Miralles et al., 2024; Du & Jing, 2024; Damerell et al., 2025; Wang et al., 2025; Tchilibou et al., 2025;
 179 Carli et al., 2025) and ongoing personal work (in the Labrador Sea, not shown) show the SWOT's abil-
 180 ity to resolve small eddies previously undetected in gridded products. Therefore, although a complete *in*

181 *situ* validation is not yet available, we are confident that SWOT’s resolution is sufficient to study mesoscale
 182 activity in our regions of interest around Iceland. Note that the noise-reduction process applied to the
 183 SSHa tends to reduce the overall energy, which complicates the interpretation in lower-energy regions due
 184 to the noise levels (see *e.g.*, Callies & Wu, 2019). In our case of study, the conclusions remain the same
 185 when comparing the noise-reduced product with the raw product.

186 For each SWOT pass, whether ascending or descending, we compute wavenumber spectra. The study
 187 of such spectra is usually made in the study of oceanic geostrophic turbulence, to quantify the distribu-
 188 tion of kinetic energy across spatial scales, revealing key dynamical processes such as energy cascades,
 189 dominant eddy sizes, and the relative importance of different forcing and dissipation mechanisms. The
 190 slope of the SSH wavenumber spectrum tends to vary between 4 and 5, depending on the underlying mech-
 191 anisms generating the eddy field (Callies et al., 2015; Lawrence & Callies, 2022; de Marez et al., 2023).
 192 In this study, we compute these spectra in SWOT data following these steps: (1) We select a 600 km along-
 193 track window centered around a position \mathbf{X} (tests confirm that the window size has minimal influence
 194 on the results). (2) The SSHa is re-interpolated onto a grid with a constant along-track spacing of 2 km.
 195 (3) We calculate the along-track wavenumber spectra for each cross-track position on the swath, exclud-
 196 ing data from the central nadir altimeter. (4) These spectra are then averaged across the cross-track po-
 197 sitions. This process produces one wavenumber spectrum for each position \mathbf{X} along the track, with a spa-
 198 tial increment of 200 km and an associated date. To facilitate comparison, all wavenumber spectra are
 199 interpolated onto a fixed wavenumber grid for averaging. Area-averaged spectra are computed by aggre-
 200 gating the spectra corresponding to positions \mathbf{X} within the specified region. It is important to note that
 201 data availability decreases during December, January, and February, as cloud cover and rain cells affect
 202 the altimetric measurements during these months.

203 We fit the wavenumber spectrum estimates to the empirical analytical model

$$\mathcal{S} = \frac{A}{1 + (k/k_0)^s}, \quad (7)$$

204 where A is the spectral amplitude (describing the overall energetic level of the flow), s is the slope, and
 205 k_0 is the transition wavenumber. We refer the reader to de Marez et al. (2023) for details and justifica-
 206 tion about this spectral model.

2.5 Andro dataset

For the estimation of eddy kinetic energy (EKE) at depth, we used the ANDRO dataset (Ollitrault & Rannou, 2013), specifically the 2024 release of the Gridded Velocity Climatology from Deep Andro Velocities. This product provides binned statistics, in particular velocity fields and EKE at a spatial resolution of $3^\circ \times 3^\circ$, based on Argo float displacements at their parking depths of 1000 m.

2.6 Biogeochemical-Argo float data

We used data from five BGC-Argo floats (Roemmich et al., 2009) north and south of Iceland, to study the difference in nutrient input, and the consequent biological production, between the two regions. The Argo program (Argo, 2000) provides a global array of profiling floats that measure temperature, salinity, and pressure across the upper 2000 meters of the ocean. The BGC-Argo array further includes sensors for key biogeochemical parameters such as pH, oxygen, nitrate, particulate backscatter, or fluorescence-derived chlorophyll. They facilitate *in situ* sampling of the water column throughout the year, allowing us to get a comprehensive picture of both surface and subsurface oceanic features. All used floats sampled every 10 days, with a vertical resolution of 5 m (upper 100 m), 10 m (100-360 m), 20 m (360-400 m) and 50 m (400-2000 m). We further re-interpolated all profiles onto a regular 10- m vertical grid, see examples of vertical sections sampled by floats in Fig. S3.

The number of profiles we can use is limited because we only considered floats equipped with fluorescence, particulate backscatter, and nitrate sensors. This resulted in three floats south of Iceland and two floats north of Iceland, spanning part of the period 2013 to 2024. Only adjusted (except for particulate backscatter) and quality controlled data (excluding flags 3 and 4, see Bittig et al., 2019) were considered, from the Coriolis DAC. We use three variables: chlorophyll-a concentration, phytoplankton carbon and nutrient availability to describe the biogeochemical activity in the area.

First, we used the fluorescence-derived chlorophyll-a, CHLA, expressed in mg m^{-3} . Only adjusted chlorophyll-a was used from the floats, where the fluorescence data have been converted to chlorophyll-a following a thorough quality control Schmechtig et al. (2015, 2023). Extended details about the method can be found in Xing et al. (2011); Petit et al. (2022), and about quality check on the BGC-Argo website (<https://biogeochemical-argo.org/>). Chlorophyll-a accuracy is 0.08 mg m^{-3} (Schmechtig et al., 2025), which leads to an uncertainty around 5% in our dataset.

Second, we estimated phytoplankton carbon, C_{phyto} , expressed in mg Cm^{-3} , based on bbp_{700} measured by the floats, following

$$C_{\text{phyto}} = 12,128 \text{ bbp}_{470} + 0.59, \quad (8)$$

237 where

$$\text{bbp}_{470} = \text{bbp}_{700} \left(\frac{470}{700} \right)^{-0.78}; \quad (9)$$

238 A more detailed explanation of the conversions can be found in Vives et al. (2024) and references therein.
 239 Uncertainties for bbp_{700} are $2 \times 10^{-4} \text{ m}^{-1}$ (Bittig et al., 2019) translating to uncertainties of $< 1\%$ for
 240 C_{phyto} in our dataset.

241 Chlorophyll-a and C_{phyto} were vertically integrated over the upper 200 m, and used as a proxy for
 242 phytoplankton growth. We acknowledge that inferring biomass from bio-optical measurements has strong
 243 caveats and limitations, particularly when phytoplankton change their physiology to adapt to environ-
 244 mental stressors. Fluorescence-derived chlorophyll-a, for example, can be affected by iron stress (Schal-
 245 lenberg et al., 2022) and light limitation Vives et al. (2024). Similarly, particulate backscatter data, and
 246 the inferred phytoplankton carbon, are sensitive to seasonal changes and shifts in community composi-
 247 tion (Cetinić et al., 2012; Schallenberg et al., 2019). The conclusions in this study are based on the rel-
 248 ative differences between the floats in the two areas, using the same conversions and estimates. Addition-
 249 ally, the difference between our two areas of interest is larger than the uncertainties from each sensor, min-
 250 imizing the interference in our results.

251 Third, we use nitrate measurements (expressed in $\mu\text{mol kg}^{-1}$) as a proxy for nutrient availability
 252 or uptake by phytoplankton. Nitrate measurements present high measurement errors up to a maximum
 253 of $1.27 \mu\text{mol kg}^{-1}$.

254 **2.7 Chlorophyll-a satellite data**

255 We analysed daily satellite-derived surface chlorophyll-a data using the Global Ocean Colour (Copernicus-
 256 GlobColour), Bio-Geo-Chemical, L4 product (doi.org/10.48670/moi-00281). Timeseries are obtained
 257 by taking the median north and south of Iceland (in areas shown in Fig. 1).

258

3 Results

259

3.1 On the stratification differences around Iceland

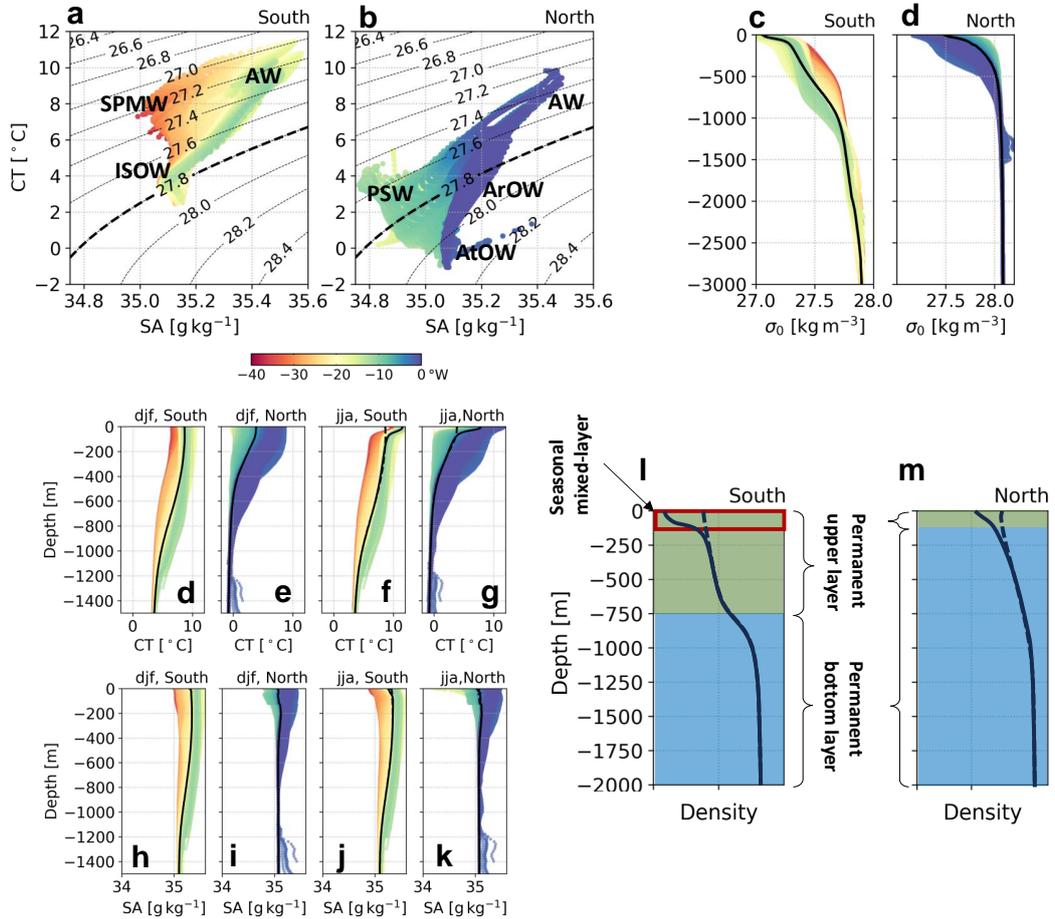


Figure 2. Stratification around Iceland from the *in situ* CORA dataset. (a, resp. b) Yearly averaged T-S diagram south (resp. north) of Iceland (see area definitions in Fig. 1). (c, resp. d) Yearly averaged potential density profiles south (resp. north) of Iceland; black line show the area average. (d-g, resp. h-k) Seasonal averaged profiles of temperature (resp. salinity) south and north of Iceland. (l,m) Schematics of the typical stratification profiles south and north of Iceland; solid line stands for the summer season and dashed line for winter; density gaps have been amplified for presentation purpose. The data is zonally binned to produce color code plots.

260 The main distinction between the regions north and south of Iceland is the contrast of water masses
261 present in each area, see Fig. 2a,b. The southern area is dominated by warm, salty Atlantic-origin wa-
262 ters, and density variations are driven largely by temperature *i.e.*, an α -ocean (Carmack, 2007). The up-
263 per layer, between the surface and ~ 500 m depth, is composed of a combination of Atlantic Water (AW)
264 and Subpolar Mode Water (SPMW), with more of SPMW in the west of the basin; below these, the Ice-
265 land Scotland Overflow Water (ISOW) dominates, see Fig. 2a. In contrast, in the north, water masses
266 are largely of polar origin, with significant freshwater input from sea ice melt. Consequently, density is
267 primarily controlled by salinity, *i.e.*, a β -ocean (Carmack, 2007). The upper layer, between the surface
268 and ~ 100 m depth, contains mainly fresh and cold Polar Surface Water (PSW) in the west of the basin,
269 and warm and salty AW in the east of the basin; Arctic Overflow Water (ArOW) and Atlantic Overflow
270 Water (AtOW) are found just below, see Fig. 2b. Given the strong difference between the northern and
271 southern water-mass properties, the stratification of the water column therefore strongly differ, see Fig. 2c-
272 k. During winter, both areas have (1) a surface layer, (hereafter called *upper layer* and represented in green
273 in Fig. 2l,m), and (2) an homogeneous bottom layer (represented in blue in Fig. 2l,m). During summer,
274 a seasonal mixed-layer appears in the south due to the increased surface temperature, while in the north
275 fresh meltwater enhances the density difference between the top and the base of the upper layer but does
276 not impact its thickness.

277 The northern and southern stratification profiles are primarily distinguished by the difference in the
278 thickness of the upper layer: it is deeper in the south than in the north, see "green layers" in Fig. 2l,m.
279 South of Iceland, classical methods for calculating the mixed layer depth, such as temperature or den-
280 sity threshold approaches fail in the summertime due to the additional thin seasonal mixed-layer (see red
281 box in Fig. 2l), and therefore underestimate the depth of the dynamically relevant surface layer (de Boyer Montégut
282 et al., 2004). A more advanced approach involves decomposing the ocean's vertical structure into nor-
283 mal dynamical modes. This method allows for the separation of the water column into different depth-
284 dependent modes that capture the full extent of stratification influences on the flow, see section 2.2 and
285 *e.g.*, Vallis (2017).

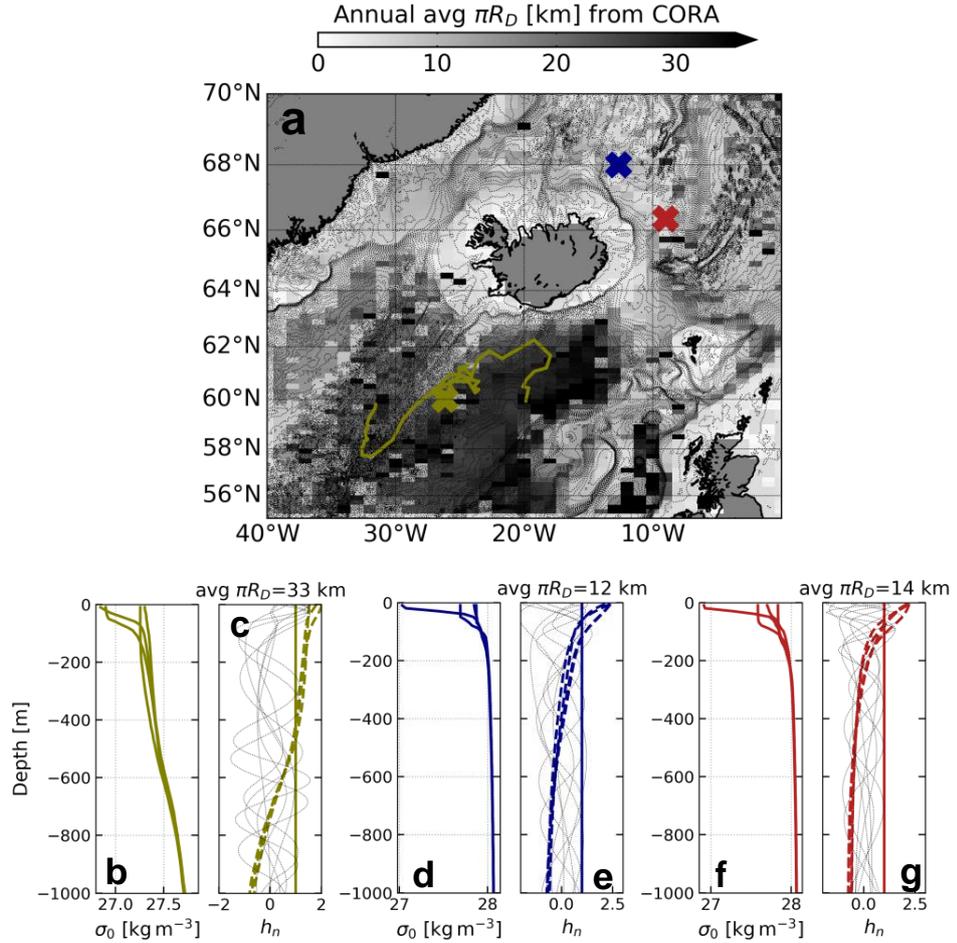


Figure 3. (a) Annual average size of mesoscale eddies πR_D from *in situ* CORA dataset, obtained by solving the normal mode equation (4), it is solely based on vertical profiles and is not affected by the resolution of the CORA product; colored crosses show the average locations where profiles of panels b-g were collected at. Yellow line recalls the trajectory of Argo float # 4903532. (b, d, f) Average potential density south of Iceland in Argo float # 4903532, in the *in situ* station marked with a blue cross, and with a red cross, respectively. (c, e, g) Dynamical modes obtained from profiles in panels (b, d, f); solid color line shows mode 0, dashed color lines shows mode 1, thin lines show modes from 2 to 10. In b-g panels, profiles are averages for each of the four seasons.

286 One of the key quantities obtained from this decomposition is the first baroclinic deformation ra-
 287 dius R_D , which represents the horizontal scale at which rotational effects balance pressure gradient forces
 288 in a stratified fluid. R_D north of Iceland is less than half the value of R_D south of Iceland, with a clear
 289 delimitation made by the Greenland-Scotland ridge, see Fig. 3a. The modal decomposition further em-

phasizes the fact that in the south, the average stratification consists of an homogeneous upper layer about 500 m deep, independently of the season; this leads to a very deep reaching mode 1, and therefore an average value of deformation radius $R_D \sim 10$ km, Fig. 3b,c. The seasonal mixed-layer that appears in summer only impacts higher modes (see similarity of mode 1 profiles in Fig. 3c). Conversely, in the north, the stratification consists of a thin upper layer, about 100 m deep, with increased density gradients during summer, see Fig. 3d,f. The mode 1 therefore reflects the shallow upper layer, and the deformation radius $R_D \sim 5$ km, see Fig. 3e,g. Our results show that the difference in water column properties north and south of Iceland have the effect to divide by ~ 2 the value of R_D north or south of Iceland. We demonstrate in the next section how this reduced upper layer thickness and deformation radius impact the energetics of the flow around Iceland.

3.2 Dynamical regimes around Iceland

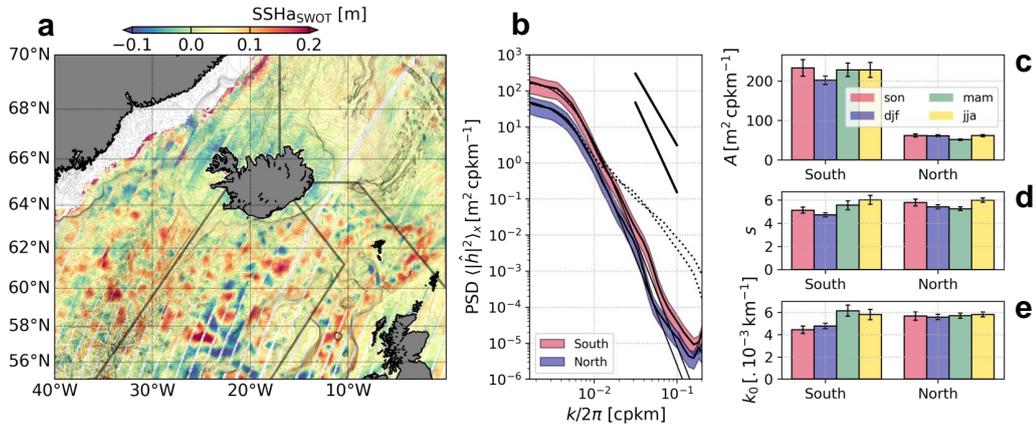


Figure 4. The surface turbulence around Iceland observed from satellite. (a) Noise-reduced SSHa from SWOT's cycle 15 (8th to 28th May 2024) descending passes. (b) Average wavenumber spectra South and North of Iceland for the month of May; solid lines show the average spectra and its fit for the noise-reduced SSHa; color envelopes show the estimated error on the average spectra; dashed line shows the average spectra using the raw SSHa; bold lines show slopes of k^{-4} and k^{-5} . (c,d,e) Amplitude, slope, and transition wavenumber for fits (see Eq. (7)) done over seasonal averaged spectra using noiseless SSHa.

Newly released SWOT altimetric data, spanning a full seasonal cycle, reveal a striking contrast in surface turbulence around Iceland: the ocean north of Iceland is significantly less energetic than the ocean to the south, see Fig. 4a. The 2-km resolution dataset provided by the SWOT (see section 2.4) enables

304 for the first time a detailed study of turbulence in places where classical altimetry used to fail in resolv-
 305 ing mesoscale motions (*e.g.*, north of Iceland, where the deformation radius is small). South of Iceland,
 306 the circulation is dominated by multiple intense—mostly anticyclonic—eddies, with SSHa reaching am-
 307 plitudes of $\mathcal{O}(30)$ cm and sizes of $\mathcal{O}(50)$ km. In contrast, north of Iceland, SSHa rarely exceed $\mathcal{O}(10)$ cm,
 308 signifying a notably quieter ocean compared to the south and much of to the global ocean.

309 To quantify these energetic differences, we computed the wavenumber spectrum of the balanced mo-
 310 tions from SSHa, using the methodology described in section 2.4. As shown in Fig. 4b, while the spec-
 311 tral shapes north and south of Iceland are similar, the amplitude is considerably lower in the north. Fit-
 312 ting the wavenumber spectrum to the model of Eq. (7) yields satisfactory results for the noise-reduced
 313 data (Fig. 4b). The spectral parameters (A , s , and k_0) display no significant seasonal variation (Fig. 4c,d,e).
 314 The spectral amplitude south of Iceland is more than double that of the north (Fig. 4c). The slope of the
 315 spectrum is around $s = 5$ (Fig. 4d), in close agreement with theoretical estimates and prior analyses (Cal-
 316 lies et al., 2015; Lawrence & Callies, 2022; de Marez et al., 2023). Independent estimates from a high-
 317 resolution regional numerical simulation (see Fig. S1) corroborate these findings: the turbulence north
 318 of Iceland is intrinsically less energetic than that in the south. Thus, while the southern region of Ice-
 319 land is characterized by intense turbulent activity, the northern region is identified as an "eddy desert".
 320 To ensure this observation is not biased by seasonal variability, we analyzed SWOT data across all avail-
 321 able cycles from August 2023 to January 2025. The consistency of spectral parameters throughout the
 322 year (Fig. 4c,d,e), as well as the persistent contrast in mesoscale activity between the north and south
 323 of Iceland (not shown), confirm the robustness of our findings.

324 Interestingly, the physiographic characteristics of the two basins are similar: (1) both consist of deep
 325 basins (approximately > 2000 m) enclosed by sharp bathymetric features; (2) the surface forcings, in-
 326 cluding surface heat flux and wind stress, are comparable (see Fig. S2). The only difference between the
 327 north and the south is, therefore, stratification.

328 Idealized simulations of forced turbulence with either a shallow or a deep surface layer (*i.e.*, differ-
 329 ent aspect ratio δ , see section 2.3) allow us to model these two distinct stratification regimes. All sim-
 330 ulations performed follow the same scenario. The energy is injected *via* BCI of the mean current because
 331 the meridional gradients of potential vorticity for the two layers have opposite signs. The growth rate of
 332 the instability is maximum at a scale of ~ 50 km $\sim \pi R_D$ km, with a value of 0.05-0.1 days $^{-1}$ (Fig. 5a).
 333 This leads to an exponential increase of Kinetic Energy (KE) during the first months of the simulation,
 334 see Fig. 5b. Energy subsequently cascades toward larger scales, and it is dissipated at the domain scale

335 (~ 500 km) by hypoviscosity. After a few simulated months, the KE reaches an equilibrium that lasts
 336 throughout the entire simulation (here ten years, see Fig. 5b). This represents a forced turbulence flow,
 337 with a typical wavenumber spectrum slope of 5 (computed during this equilibrium state, see Fig. 5c,e-
 338 j). This value is similar to our SWOT estimates. Strikingly, the energy content of the turbulence increases
 339 along with δ , see Fig. 5d. BCI converts the potential energy of the flow, stocked in the surface layer, into
 340 eddy kinetic energy. Therefore, the smaller δ , the smaller available potential energy, and the weaker tur-
 341 bulance.

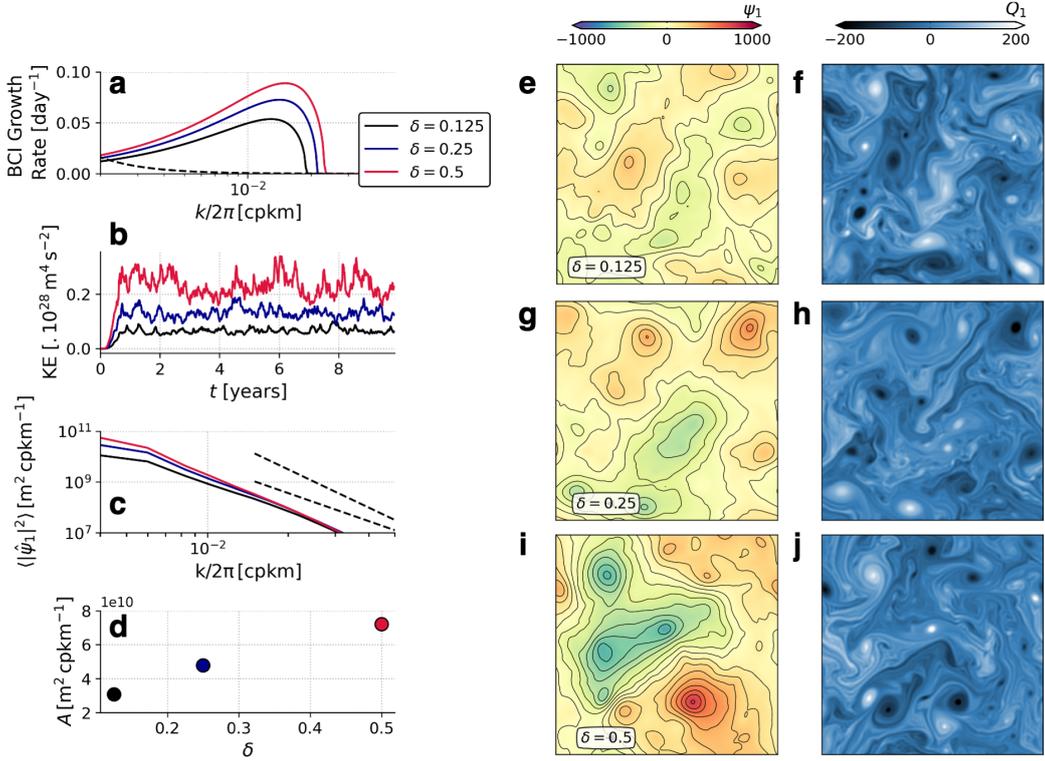


Figure 5. Results from idealized numerical simulations of forced quasi-geostrophic turbulence. (a) Growth rate of the most unstable mode for the primary BCI following linear stability calculation; dashed line shows the damping rate from hypoviscosity. (b) Surface integrated Kinetic Energy evolution in the surface layer of the QG simulation. (c) Surface streamfunction marginal wavenumber spectra; dashed lines show spectrum slopes of 4 and 5. (d) Spectral amplitude (from fits of Eq. (7) over wavenumber spectra of panel c) as a function of δ . The colors in panels a,b,c,d represent the value of δ . (e,g,i) Snapshot of surface streamfunction, equivalent to the SSH (in adimensionalized unit) for the different values of δ . (f,h,j) Snapshot of surface potential vorticity (in adimensionalized unit) for the different values of δ . In panels e-j, the domain size is 500 km.

342 In a nutshell, these simulations explain why the flow is so calm north of Iceland: very little poten-
343 tial energy is available to be converted to kinetic energy through BCI, and therefore turbulence is weak.
344 The SSHa measured by SWOT (Fig. 4a) can be compared to the surface streamfunction in simulations
345 (Fig. 5e,g,i). With small δ (Fig. 5e), eddies are weaker than with larger δ (Fig. 5i), analogously than the
346 north and south of Iceland. These conclusions rely on the open-ocean mesoscale dynamics hypothesis.
347 We cannot exclude the possibility that boundary and shelf-break processes (*e.g.*, Kelvin waves, topographic
348 Rossby waves, upwelling, etc.) play a role in the observed regional variability (Wise et al., 2024). Fur-
349 ther investigation, particularly with dedicated *in situ* observations and numerical model analysis, would
350 be necessary to fully assess their potential role. Also, it is worth mentioning that there exists an inter-
351 play between the turbulent motions and the upper layer formation. As turbulence develops, eddies pro-
352 mote vertical motions within the upper ocean. This intensified mixing further impact the stratification
353 set by surface heat fluxes, solar heating and other air-sea interactions, allowing the upper layer to thicken
354 (Klein & Lapeyre, 2009).

355 **3.3 Phenomenology of oceanic currents north and south of Iceland**

356 In the ocean, except where large western boundary currents are present, mesoscale eddies, are the
357 primary contributors to the surface-intensified turbulence (see *e.g.*, Z. Zhang et al., 2014) as observed by
358 SWOT altimetry (and realistic simulations). Usually, these structures can be observed and statistically
359 analyzed using eddy detection in gridded altimetry products (see *e.g.*, Laxenaire et al., 2018; Ernst et
360 al., 2023). However, the reduced deformation radius north of Iceland (Fig. 3) complicates this approach
361 and makes the interpretation of the eddy detection results very challenging. To better understand the
362 mesoscale dynamics around Iceland, we use *in situ* data from Argo floats. They further confirm the two
363 distinct regimes identified in the precedent section.

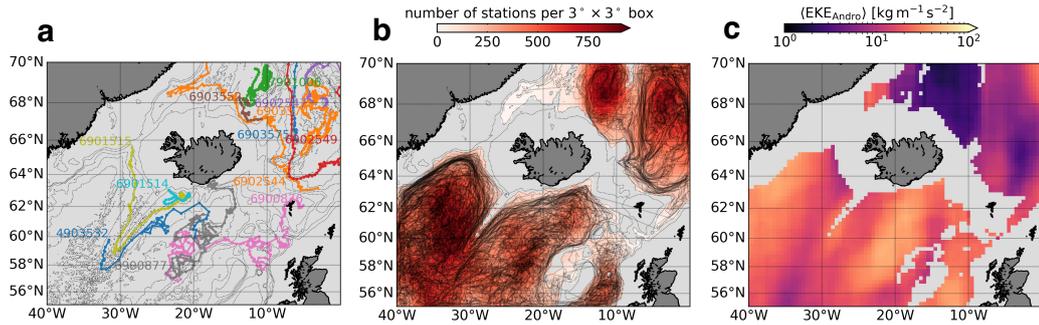


Figure 6. (a) Trajectories of selected Argo floats in the area of interest. (b) Trajectories of Argo floats used in the ANDRO dataset, and the number of stations used to compute the EKE in $3^\circ \times 3^\circ$ bins. (c) Binned average of Eddy Kinetic Energy at 1000 m depth from Argo displacement.

364 South of Iceland, dynamics are primarily driven by deep-reaching mesoscale eddies in the middle
 365 of the Iceland Basin. For instance, floats #6900877 and #6900876, in the southern region, are trapped
 366 by eddies that extend to their 1000-m parking depth (Fig. 6a). More generally, the trajectories of most
 367 floats that enter the interior of the Iceland Basin appear to be advected by vortical structures, resulting
 368 in turbulent trajectories (Fig. 6b). This behavior leads to high values of EKE at 1000 m depth in this
 369 region (Fig. 6c), comparable to those found in high-energy areas such as the western boundaries of oceanic
 370 basins (de Marez et al., 2020). These findings support and align with previous studies using *in situ* data
 371 (see *e.g.*, Zhao et al., 2018b; Johnson et al., 2024). On the edges of the Iceland basin, trajectories of floats
 372 #6901515, #4903532, and #6901514, in the southern regime, highlight the presence of intense bottom
 373 currents flowing along the topography (see *e.g.*, Brakstad, Gebbie, et al., 2023; de Marez et al., 2024)

374 Conversely, north of Iceland, dynamics are mainly driven by weak, deep near-laminar currents. Lim-
 375 ited *in situ* datasets from cruises are available in this area, although more Argo floats have been deployed
 376 there in more recent years (Jayne et al., 2017). The trajectories of some of these floats (Fig. 6a) reveal
 377 that they predominantly follow straight paths, as seen in the trajectories of floats #6903575 and #6902549.
 378 This characteristic is also seen in the trajectories of all Argo floats in the northern regime (Fig. 6b). Also,
 379 the distance between sampling stations—with 10-day intervals—is small, see floats #7901006 (in green),
 380 and #6903552 (in brown). This results in EKE values at 1000 m depth that are approximately one or-
 381 der of magnitude weaker than those observed south of Iceland (Fig. 6c).

4 Discussion

4.1 Unveiling the mesoscale dynamics north of Iceland

Our study provides a new perspective on oceanic circulation around Iceland. Previous studies have focused primarily on the Iceland Basin and their role in heat transport (Zhao et al., 2018a, 2018b), while turbulence in the north remained largely undocumented due to observational limitations (Beaird et al., 2013). Here, high-resolution SWOT altimetry dataset allowed us to provide the first detailed characterization of mesoscale turbulence north of Iceland. We show that the ocean north of Iceland is significantly less energetic than its southern counterpart, in contrast to the south, where strong mesoscale eddies dominate the flow.

These results challenge the assumption that bathymetry or surface forcing alone determine regional differences in ocean dynamics. Instead, we show that the primary driver of this difference is stratification. South of Iceland, a thick upper layer (~ 500 m) supports strong baroclinic instability, allowing eddies to grow and sustain energetic turbulence. In contrast, the north has a much shallower upper layer (~ 100 m), limiting the energy available for eddy formation. These findings contribute to broader research on subpolar ocean variability (Brakstad, Gebbie, et al., 2023) and underscore the importance of high-resolution observations in improving our understanding of mesoscale dynamics in previously unresolved regions. Future work combining SWOT data with additional *in situ* measurements and high-resolution numerical modeling will be essential to further investigate the mesoscale variability at high latitude.

4.2 Implications for ocean biogeochemistry around Iceland

Mesoscale eddies influence ocean biogeochemistry through the water column globally (Su et al., 2021; Cornec et al., 2021). Eddies facilitate upwelling of nutrients (Patel et al., 2020) and carbon flux (Moreau et al., 2017). For instance, cyclonic eddies create optimal conditions for phytoplankton growth in deeper layers when nutrients are all consumed at the surface, creating subsurface or deep chlorophyll-a maxima (Cornec et al., 2021). The North Atlantic is known for being a highly productive area during spring (Westberry et al., 2016). Particularly, in the Iceland Basin, studies using remote sensing (Siegel et al., 2002), and cruise data (Poulton et al., 2010), show highly productive blooms during the growing season. Although cruise data for the Iceland Sea are more limited, they suggest it to be less productive (Jeansson et al., 2015; Richardson & Bendtsen, 2021), thus suggesting a higher primary productivity in the south than in the north of Iceland (Thórdardóttir, 1986; Zhai et al., 2012; Cerfonteyn et al., 2023).

411 We test the hypothesis of the link between mesoscale activity and biological production by analyz-
 412 ing chlorophyll-a and phytoplankton carbon data from BGC-Argo floats (see section 2.6). The timeseries
 413 suggest a clear difference in phytoplankton growth between the northern and the southern regions. We
 414 observe more prolonged, and overall larger, phytoplankton blooms in the south compared to the north
 415 (see Fig. 7).

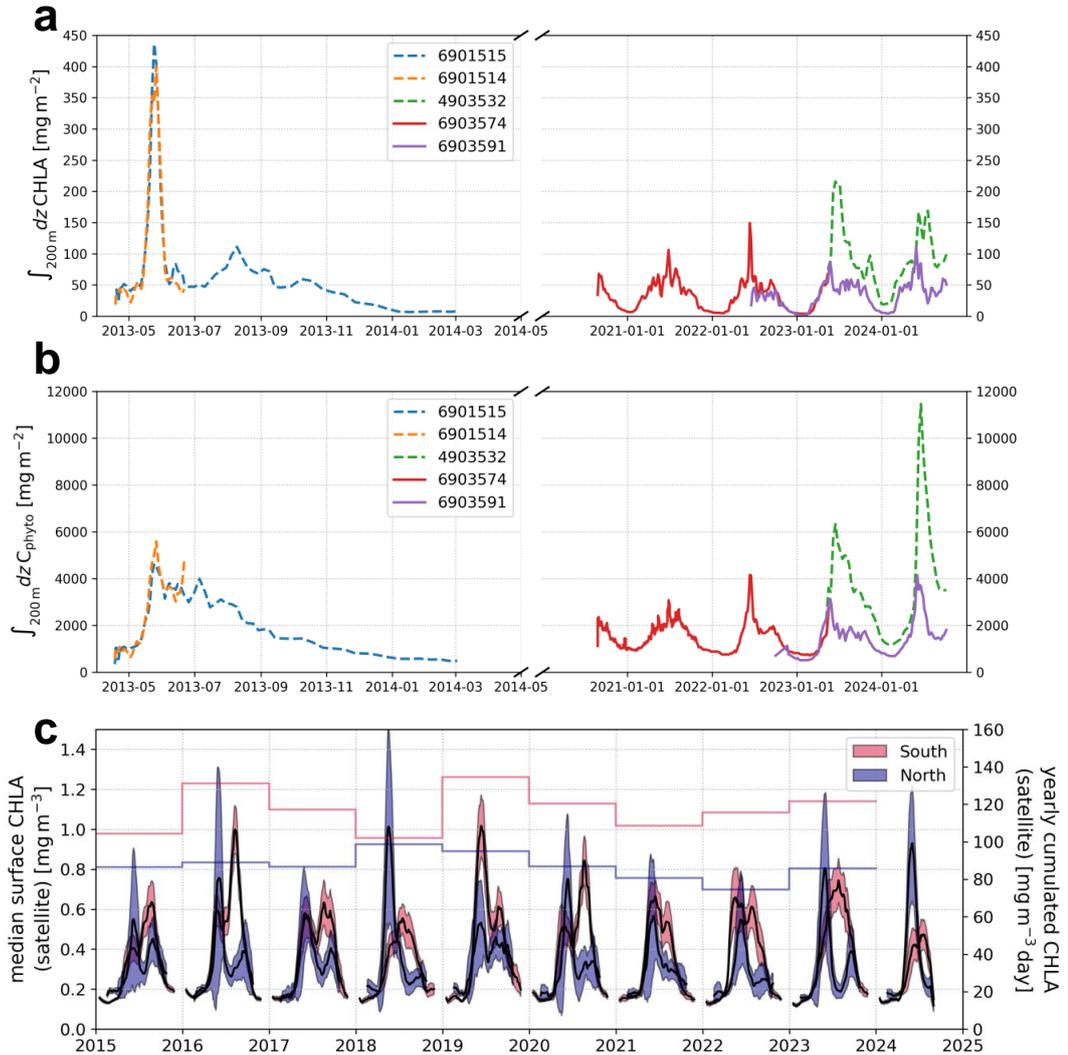


Figure 7. (a, resp. b) Timeseries of chlorophyll-a (resp. phytoplankton carbon) integrated in the upper 200 m from BGC-Argo floats; solid (resp. dashed) lines are for floats located north (resp. south) of Iceland. (c) Median chlorophyll-a north and south of Iceland from satellite measurements in a 10-year period (see section 2.7); envelopes show the standard deviation in the areas; bar plot (right y-axis) shows the cumulated median SSH over each year.

416 The northern region has lower concentrations of chlorophyll-a and phytoplankton carbon compared
417 to the southern region, particularly during the phytoplankton growing season in 2023 and 2024 (*i.e.*, April
418 to June, Fig. 7a,b). In 2024, for example, where there are two floats sampling at the same time, the south-
419 ern regime shows more than double of chlorophyll-a concentrations (200 mg m^{-2}) in the water column
420 compared to the north (50 mg m^{-2}). However, fluorescence-derived chlorophyll-a can sometimes be bi-
421 ased due to iron stress (Schallenberg et al., 2022) and light limitation of phytoplankton growth (West-
422 berry et al., 2008). We therefore use phytoplankton carbon, as an additional and more reliable proxy of
423 phytoplankton production at high latitudes (Vives et al., 2024). The difference in phytoplankton carbon
424 is even greater between the two regimes, where values are up to $10,000 \text{ mg C m}^{-2}$ higher in the south
425 compared to the north.

426 While the temporal frame of these floats is small, a decade of satellite-derived chlorophyll-a con-
427 centrations exhibits the same pattern. Blooms are longer in the southern regime. This leads to overall
428 higher yearly integrated median surface chlorophyll-a in the south compared to the north (Fig. 7c), *i.e.*,
429 higher production over time in the south than in the north.

430 Float data shows slightly higher, although not significantly, nitrate concentrations at the surface
431 and in the upper 200 m in the northern regime compared to the south (Fig. S4). Along with the differ-
432 ences in productivity, this could suggest higher nutrient uptake in the southern regime, where phytoplank-
433 ton are able to utilize more available nutrients to continue growing, leading to prolonged blooms. How-
434 ever, the Iceland Basin is known for being iron limited in the summer (Ryan-Keogh et al., 2013; Nielsdóttir
435 et al., 2009; Moore et al., 2013), meaning that nitrate concentrations may not be representative of nu-
436 trient availability in this area.

437 These findings are consistent with previous long-term observations, using either satellite algorithms
438 to derive primary production or *in situ* data, that point out the high productivity of the Iceland Basin
439 compared to the Iceland Sea (Thórdardóttir, 1986; Zhai et al., 2012; Richardson & Bendtsen, 2021; Cer-
440 fonteyn et al., 2023). This may lead to hypothesize that the high mesoscale turbulent activity in the south
441 ultimately affects the biogeochemical variability south of Iceland, thus explaining the difference in pro-
442 ductivity between the north and the south. We could indeed hypothesize that the larger upper layer thick-
443 nesses and stronger mesoscale current activity at greater depths may be facilitating nutrient upwelling
444 to the surface in the south. There, the higher upwelling of nutrients may facilitate an increase in longevity
445 and vertical extension of blooms. On the contrary, thinner upper layer thicknesses and weaker EKE val-
446 ues in the north, primarily occurring at shallower depths, may be limiting nutrient supply to the surface.

447 As a result, phytoplankton is subjected to the surface in the northern region, where growth may be ter-
448 minated promptly with the diminishment of nutrient availability earlier in the season. While consistent
449 with our observations, these interpretations remain speculative and require further data for confirmation.
450 Our study serves as a first step toward unveiling the coupling between physical and biogeochemical pro-
451 cesses around Iceland, highlighting the need for future targeted observations and modeling.

452 **4.3 Future Changes in a warming climate**

453 Recent studies suggest that global warming is significantly affecting the upper ocean stratification
454 around Iceland, driven by changes in both temperature and salinity, which directly influence the upper
455 layer thickness. The region, particularly in the northeast, has warmed in the last decades (Polyakov et
456 al., 2017; Dai et al., 2019), as the influx of warmer AW initiated a process known as "Atlantification"—where
457 Arctic waters increasingly take on the characteristics of Atlantic conditions due to the penetration of warmer,
458 saltier waters. The effects of Atlantification are twofold. In northern regions, driven by the increasing in-
459 trusion of Atlantic Water, Atlantification is expected to weaken the typical upper-ocean stratification.
460 This would promote deeper mixing, resulting in a deeper mixed layer. In contrast, south of Iceland, sur-
461 face warming will strengthen upper-ocean stratification, limiting vertical mixing to shallower depths and
462 likely leading to a shallower mixed layer depth. These changes may have significant impacts on dynam-
463 ical biogeochemical activity in the region. A thicker upper layer in the northern areas would allow for
464 a more turbulent ocean with greater vertical mixing, which could enhance nutrient availability, support-
465 ing biological productivity at the surface if other conditions are optimal for phytoplankton growth. How-
466 ever, in the southern region, where the upper layer is expected to become shallower, eddy generation *via*
467 BCI may decrease, thus limiting nutrient upwelling from deeper waters. This could reduce primary pro-
468 ductivity, particularly during the growing season, as essential nutrients for phytoplankton growth become
469 less accessible. As the southern region is observed to be more productive, warming-derived changes might
470 have consequences for the biological carbon pump. Overall, this might shift northward the primary pro-
471 duction patterns, and therefore perturb the entire marine food web around Iceland.

472 **Open Research Section**

473 CORA data, can be downloaded from the Copernicus Marine Service at [https://data.marine.copernicus](https://data.marine.copernicus.eu/product/INSITU_GLO_PHY_TS_OA_MY_013_052/)
474 [.eu/product/INSITU_GLO_PHY_TS_OA_MY_013_052/](https://data.marine.copernicus.eu/product/INSITU_GLO_PHY_TS_OA_MY_013_052/). BGC-Argo float data are available on the Ifremer Argo
475 Data Assembly Center at <https://data-argo.ifremer.fr/dac>. *in situ* data from the two stations north
476 of Iceland is provided by the SeaDataNet Pan-European infrastructure for ocean and marine data man-
477 agement (<https://www.seadatanet.org>), and can be downloaded as part of the SDC_ARC_DATA_TS_V2
478 dataset and the Norwegian Marine Data Center (Brakstad, Våge, et al., 2023). The measurement were
479 made during the observational program of the Icelandic Marine and Freshwater Research Institute. QG
480 numerical simulation code is available at <https://github.com/joernc/QGModel>. SWOT data can be
481 downloaded on AVISO website <https://www.aviso.altimetry.fr/en/my-aviso-plus.html>. ANDRO
482 dataset is available through the SEANOE platform and is part of the comprehensive ANDRO Argo-based
483 velocity product managed by Coriolis, at <https://www.seanoe.org/data/00360/47077/>. Satellite mea-
484 surements of chlorophyll-a can be downloaded using E.U. Copernicus Marine Service Information at [https://](https://data.marine.copernicus.eu/product/OCEANCOLOUR_GLO_BGC_L4_MY_009_104/)
485 data.marine.copernicus.eu/product/OCEANCOLOUR_GLO_BGC_L4_MY_009_104/. Scripts to reproduce re-
486 sults of this paper can be obtained online (de Marez, 2024).

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