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Role of bathymetry in Agulhas Current configuration and behaviour

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Abstract:

The Agulhas Current forms an important link in the global ocean thermohaline circulation by its role in the inter-ocean exchange of water south of Africa. This process of ring shedding at the current's retroflexion is dependent on perturbations to its trajectory that are sensitive to bathymetry. These perturbations may furthermore force the current to intersect shallow regions resulting in substantial changes to its path. A number of other flow characteristics of the system have also been deemed to be influenced by bathymetry. How dependent is Agulhas Current behaviour therefore on the bottom topography? We have used a regional, primitive equation model for initial experimentation. Removing the Agulhas Bank leads to enhanced inter-ocean flux, indicating its importance for inter-ocean exchange. Excising the Agulhas Plateau causes meridional meanders in the Agulhas Return Current to be unlocked from the bathymetry. Smoothing the continental slope weakens the current and substantially increases the direct inter-ocean flux.

25 **1. Introduction**

26 The Agulhas Current (AC in the following) is a major western boundary current
27 along the south-east coast of Africa. South of the continent it retroflects and most of
28 its water subsequently becomes part of the Agulhas Return Current [*Lutjeharms and*
29 *Ansorge, 2001*] that carries out large meridional meanders on its way eastward. The
30 retroflexion of the AC is unstable and creates large Agulhas rings by loop occlusion.
31 This is the prime mechanism by which warm and salty water from the Indian Ocean is
32 transferred to the South Atlantic Ocean [*Gordon, 1986*]. Ring spawning events may
33 be induced by the shedding of a lee eddy from the western side of the Agulhas Bank
34 or by the arrival of a Natal Pulse [*Van Leeuwen et al., 2000*], a singular meander,
35 from far upstream. A well-developed Natal Pulse may even cause an upstream
36 retroflexion [*Lutjeharms and van Ballegooyen, 1988*] that will prevent AC water
37 from reaching the normal retroflexion location thus temporarily interrupting inter-
38 ocean exchange. All these flow features are in some way dependent on the
39 bathymetry.

40 It has been shown that the generation of Natal Pulses is due to an anomalously
41 weak continental slope [*De Ruijter et al., 1999a*] at the Natal Bight. When this
42 meander precipitates an early retroflexion, this is due to the current being forced
43 across shallower topography of the Agulhas Plateau (location: see Figure 1). The
44 disposition of the retroflexion itself may be a function of the shape of the Agulhas
45 Bank [*De Ruijter et al., 1999b*] as is the presence of a lee eddy on its western side.
46 The meridional meanders in the Agulhas Return Current in turn are thought to be
47 forced by the shallow topography of the Agulhas Plateau [*Lutjeharms and van*
48 *Ballegooyen, 1984*] and by the poleward extension of the Mozambique Plateau
49 [*Gründlingh, 1977*]. The sensitivity of the AC to the bathymetry has also been

50 indicated by modelling [e.g. *Lutjeharms and Webb, 1995; Matano, 1996*]. We have
51 therefore experimented by removing certain key components of the bottom
52 topography in a more refined model to see how the current configuration would react
53 and thus to establish the importance of each of these components to the normal
54 current configuration.

55

56 **2. The regional model**

57 Our circulation model is based on the IRD-UCLA version of the Regional Ocean
58 Modelling System (ROMS) [*Shchepetkin and McWilliams, 2003; 2005 ; Penven et al.,*
59 *2005*]. The model domain extends from 5.8°E to 34°E and from 25.4°S to 44°S (Fig.
60 1). The model grid is 168×136 points with a resolution of 1/6° corresponding to a
61 mean grid spacing of 12 km, which resolves the first baroclinic Rossby radius of
62 deformation here (about 30 km, *Chelton et al., 1998*). The grid is isotropic and does
63 not introduce any asymmetry in the horizontal dissipation of turbulence. Therefore, it
64 allows a fair representation of mesoscale dynamics. The bottom topography is derived
65 from a 2' resolution database [*Smith and Sandwell, 1997*]. Although a new pressure
66 gradient scheme associated to a specific equation of state limits errors in the
67 computation of the pressure gradient (*Shchepetkin and McWilliams, 2003a*), the
68 bathymetry has been filtered in order to keep a "slope parameter" (*Beckmann and*
69 *Haidvogel, 1993*) $r = \frac{\Delta h}{2h} = \frac{h^{+1/2} - h^{-1/2}}{H^{+1/2} + H^{-1/2}} \leq 0.3$ for the control run (and smaller for a
70 higher topographic smoothing).

71 The model has 32 vertical levels and the vertical s-coordinate is stretched for
72 boundary layer resolution. All the model external forcing functions are derived from
73 climatologies. At the surface, the model heat and fresh water fluxes are extracted from
74 the COADS climatology [*Da Silva et al., 1994*]. For the wind stress, a monthly mean

75 climatology is computed from QuikSCAT scatterometer data. At the four lateral
76 boundaries an active, implicit, upstream-biased, radiation condition connects the
77 model solution to the surroundings [Marchesiello *et al.*, 2001]. In the case of inflow,
78 the solution at the boundary is nudged toward a climatological velocity field
79 calculated from the OCCAM 1/4° global ocean model that is also used as initial
80 condition. All the simulations discussed in this manuscript were run for 11 years and
81 model outputs were averaged and stored every 5 days of simulation.

82

83 **3. Results**

84 The results of four runs of the model are given here. These are a control run with
85 fully intact bathymetry, a run without the Agulhas Bank, a run without the Agulhas
86 Plateau and a run with a much smoothed shelf. Results from the control run are given
87 in Figure 1.

88 The sea surface temperatures in this figure reproduce the known characteristic
89 flow patterns of the region with a high degree of verisimilitude. The AC appears as a
90 narrow ribbon at the shelf edge of the east coast with surface temperatures exceeding
91 26 °C and an annual mean volume flux to the sea bottom at 30° E of 75 Sv ($Sv = 10^6$
92 m^3/s) [viz. *Beal and Bryden*, 1999]. South-west of the tip of Africa it retroflects.
93 North-west of this retroflexion there is evidence for a newly shed Agulhas ring
94 (named A in Fig. 1a) while between these is an equatorward moving filament of cold,
95 subantarctic water with a temperature of less than 14 °C [viz. *Lutjeharms and Fillis*,
96 2003]. The frequency of ring shedding events in the model is realistic at about 4 per
97 year [viz. *de Ruijter et al.*, 1999b]. The meridional meander of the Agulhas Return
98 Current over the Agulhas Plateau is clearly circumscribed. Eddies shed by this
99 meander move westward [Boebel *et al.*, 2003a]. South of this meander there is a warm

100 Agulhas eddy [Lutjeharms, 1987] that has entered the subantarctic zone (named B in
101 Fig. 1a). Even a number of smaller features are well-represented. These include
102 upwelling inshore of the current at the eastern extremity of the Agulhas Bank
103 [Lutjeharms *et al.*, 2000], a cyclonic lee eddy west of this part of the shelf (named C
104 in Fig. 1a) [Penven *et al.*, 2001a] and an AC filament [Lutjeharms and Cooper, 1996]
105 being drawn equatorward in the South Atlantic. The altimetric results show a number
106 of circulation features even more clearly.

107 In Figure 1b, the anti-cyclonic nature of the southern AC system stands out (warm
108 colours). The meander over the Agulhas Plateau is again well-represented as is the
109 retroflection extending to about 16° E on this occasion. A newly spawned Agulhas
110 ring is evident in the Cape Basin to the west of the subcontinent as are a number of
111 weaker remnants of rings all moving in a north-westward direction [Schouten *et al.*,
112 2000]. Some split, amalgamate with other rings or interact with cyclones [Boebel *et*
113 *al.*, 2003b] that move in a south-westward direction. The lee eddy west of the Agulhas
114 Bank is particularly prominent and is often seen to cut through the retroflection loop
115 [Lutjeharms *et al.*, 2003] thus synchronised with a ring shedding event.

116 To evaluate the Indo-Atlantic inter-ocean exchange we made use of the ARIANE
117 off-line Lagrangian diagnostic [<http://univ-brest.fr/lpo/ariane>; e.g. Blanke *et al.*,
118 1999]. Inter-ocean transport is then computed by releasing 140,000 virtual particles
119 across a zonal section of the AC at 32°S in the Indian Ocean. and by integrating their
120 individual trajectories and related infinitesimal transport forward in time till they
121 reach defined final sections. These vertical sections completely close the modelled
122 area and are located in the Atlantic, Southern Ocean and Indian sectors of the regional
123 domain (Fig. 1c). Each trajectory is computed offline and integrated sequentially on
124 the 5-day mean fields of the simulation. The virtual particles are released starting

125 from year 4 of the simulation. We stop the deployment at year 8 allowing to the last
126 released particles a 3-year delay to exit the domain. At the end of the integration, only
127 a very small percentage of particles are still in the domain (about 2%). The water
128 mass transfer between the AC and the South Atlantic thus derived is, in the control
129 run, 41 ± 2 Sv. The uncertainty on the mass transfers was estimated from the
130 sensitivity of the mass transfer to the particular sampling period adopted for the
131 storage of the model output. This represents 55% of the total AC transport computed
132 at 32°S and it is at the very high end of estimates of such fluxes to date. This is
133 probably due to two different factors. First, the regional modelled domain is relatively
134 small and therefore the final sections for the Lagrangian integration that close the
135 South Atlantic and Southern Ocean sectors are very close to the African continent and
136 still embedded in the very turbulent regime of the Cape Basin. This could induce an
137 overestimate of water transfer to the South Atlantic, while, in reality, as a result of
138 different mesoscale interactions, part of this water recirculates back to the Indian
139 Ocean. Indeed, the Agulhas water flux that crosses the Atlantic section north of 35°S
140 is only 25.4 ± 1.2 Sv. The remaining 15.6 ± 0.8 Sv of the computed leakage leave the
141 Cape Basin with a south-west direction and reach the Atlantic final section south of
142 35°S . Second, the initial and open ocean boundary conditions are a monthly
143 climatology derived from OCCAM, a global ocean model and not an observed
144 climatology. Deviations of the mean thermohaline structure of OCCAM from
145 observations could induce a difference in magnitude for the Indo-Atlantic connection.

146 The strong correspondence between these simulations and the known
147 characteristics of the current system, as reflected in the cited literature, therefore gives
148 us considerable confidence that this model incorporates the appropriate physics and

149 captures the scales and the behaviour of the current adequately to experiment with the
150 bathymetry. In the first experiment (Figure 2) the Agulhas Bank has been removed.

151 The most immediately striking aspect of this simulation is that the AC hugs the
152 now zonal shelf edge south of Africa continuously. An excessive leakage of AC water
153 into the South Atlantic of 56 ± 2.8 Sv takes place, (average for 8 model years) or 69%
154 of the total. This large leakage appears also from the sea surface temperature structure
155 (Figure 2a). A retroflective behaviour is present all the time, but the surface layers of
156 the AC only take part in this about 46% of the time (viz. Figure 2b) usually moving
157 directly west (Figure 2a). A lee eddy is formed on the western side of the land mass
158 where the current overshoots, but is considerably more prominent than when the
159 Agulhas Bank is present. This lee eddy passes south-westward between the ring and
160 the new retroflection loop on 72% of the ring-shedding events (e.g. Figure 2a), more
161 clearly seen in the sea surface height than in surface temperatures. We can only
162 surmise if the movement of this eddy is opportunistic, when a gap appears between
163 ring and retroflection, or is itself the cause of the ring shedding event. The location of
164 the retroflection lies at least 3° of latitude further north than in the control run, but not
165 further west. The latitude of the Subtropical Convergence remains virtually the same,
166 at a mean of 42° S, making the retroflection loop much wider than normal. Meanders
167 in the Agulhas Return Current are realistic and relatively stationary, whereas cold
168 eddies shed from these meanders all move westward. An occasional leakage
169 reminiscent of an upstream retroflection is seen. While the Agulhas Bank is almost
170 completely removed, a small upwelling cell still exist inshore of the AC . Removal of
171 the Agulhas Plateau leads to different current behaviour.

172 First, the behaviour of the Agulhas retroflection is much like that in the control
173 run, including the location of the retroflection (not shown) and the average number of

174 ring shedding events. However, the average volume transport of the AC is reduced to
175 66 Sv. This is due to the enormously reduced recirculation, and therefore of inertia
176 and water entrainment, west of 32° E in the absence of the Agulhas Plateau. The mean
177 Lagrangian flux into the South Atlantic is 34 ± 1.7 Sv. This value is lower than that
178 for the control run, but still represents more than 50% of the total Agulhas transport at
179 32°S. The major change for this experiment is in the meanders in the Agulhas Return
180 Current. When the topography is removed, non-stationary Rossby wave-like meanders
181 forms as they are not anymore constrained to one geographic location as in the control
182 run. They persistently move westward at about one degree of longitude in $11 (\pm 3.6)$
183 days.

184 The effects of reducing the steepness in the continental slope around southern
185 Africa (by decreasing the “slope parameter” r to 0.1) are given in Figure 3. First, the
186 surface speed of the AC is reduced from > 2 m/s in the control run to < 0.8 m/s with
187 this smoothed slope (Figure 3). The current is wider and more diffuse. The intensity
188 of the retroflection is much reduced with a considerable proportion of the current
189 instead following the 1000 m isobath around the tip of Africa into the South Atlantic
190 (Figure 3). The volume flux of the current is reduced to only 65 Sv in this experiment,
191 but the percentage leakage into the South Atlantic is increased to 64%. The weaker
192 the slope gradient, the less inertial the current is and the less will be the tendency to
193 enter the South Atlantic as a free jet in a south-westerly direction and to retroflect.
194 The propensity of the current core to continue to hold close to the shelf edge, well into
195 the South Atlantic, may thus be increased, as is seen in Figure 3.

196

197 **4. Conclusions**

198 These preliminary modelling experiments show that the removal of certain
199 prominent parts of the bottom topography at the AC termination has some important
200 effects on the current's disposition. Removal of the Agulhas Bank leads to a
201 substantial increase in the volume flux of the current into the South Atlantic and a
202 seemingly increased role for a lee eddy off the west coast on the timing of ring
203 shedding events. Excising the Agulhas Plateau leads to meridional meanders in the
204 Agulhas Return Current moving steadily westward while the volume flux of the AC is
205 reduced.

206 The sensitivity of the AC to bathymetry is particularly evident in experiments with
207 the steepness of the shelf slope. Decreasing steepness leads to decreased speeds in the
208 current, a less concentrated current and a greater tendency for it to move directly into
209 the South Atlantic and not to form Agulhas rings.

210 The model we have used has a number of critical limitations. The one concerns
211 the perennial quest for higher spatial resolution in models; the other the inadequacy of
212 the boundary conditions. Both factors result in a lack of perturbations to the flow of
213 the AC itself in the model. Such perturbations, in the form of the Natal Pulse, have
214 been shown [e.g. *van Leeuwen et al.*, 2000] to be crucial to a proper understanding of
215 the mechanisms responsible for inter-ocean exchange in the system. In order therefore
216 to simulate the true situation better, improved model runs that include realistic
217 mesoscale perturbations will doubtless improve these initial results.

218

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224

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311 **Figure 1.** Sea surface temperature distribution simulated for the southern Agulhas
312 Current region on 3 March, model year 11 (**a**), with a temperature scale. Black lines
313 show the bottom topography in km. The sea surface height in cm is given in panel (**b**),
314 for 28 December, model year 1. Note the locations of the Agulhas Bank and the
315 Agulhas Plateau. (**c**) Interocean water mass transfer (with a 5-Sv C.I.) originating in
316 the AC. The four sections of interception are also shown in red: “Agulhas current”
317 (solid line), “Indian ocean” (dashed line), “Southern Ocean” (dotted line), and
318 “Atlantic Ocean” (dash-dotted line).

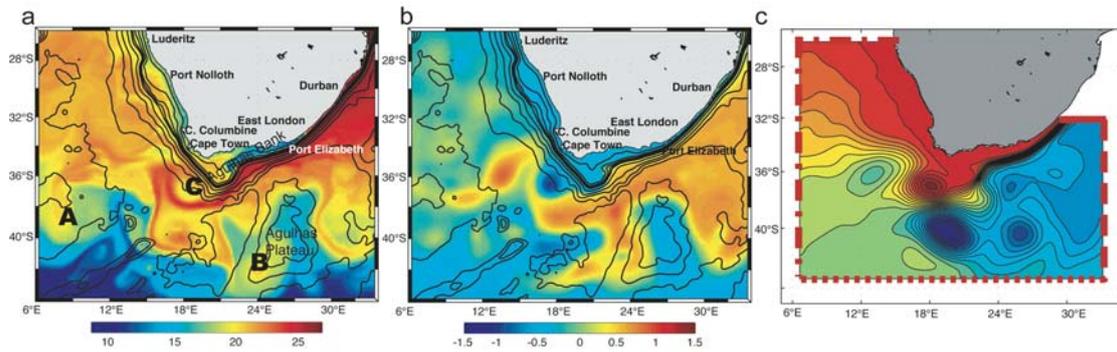
319

320 **Figure 2.** The sea surface temperature distribution for the southern Agulhas Current
321 region on 13 October, model year 4, when the Agulhas Bank has been removed (panel
322 **a**). It shows a reduced retroflexion. For 3 September, model year 4, (panel **b**) the
323 retroflexion is meridionally wide, but much better developed. Otherwise as in Figure
324 1.

325

326 **Figure 3.** The sea surface temperature for the Agulhas Current termination with a
327 smoothed and weakened shelf slope. Otherwise as in Figure 1.

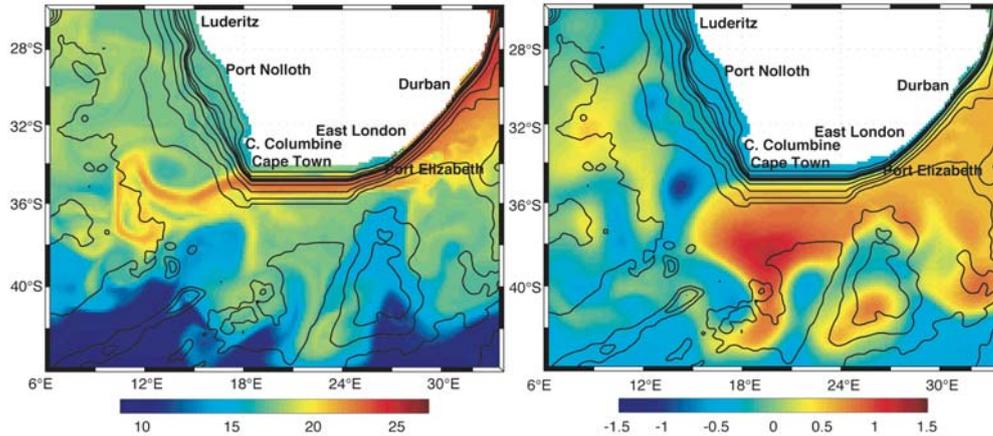
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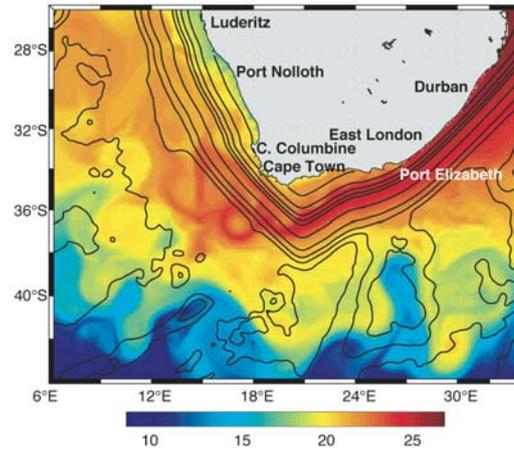
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